# Metrizability of Affine Connections

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#### Abstract

An affine connection  $\Gamma$  on a vector bundle  $\eta = (E, \pi, M, V)$  of a rank r is called Riemann metrizable if there exists on M a Riemann metric which preserves the scalar product of vector fields parallel displaced according to  $\Gamma$ .  $\Gamma$  determines a connection G in a bundle, where M is fibered by the manifold of the ellipsoids of  $R^r = \pi^{-1}, x \in M$ . We prove that  $\Gamma$  is Riemann metrizable iff G is integrable.

An analogous result is deduced in the case, where  $\eta$  is replaced by a Finsler vector bundle,  $\Gamma$  means a Finsler connection, and the metric is a Finsler metric.

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## 1 Introduction

We consider a vector bundle  $\eta = (E, \pi, M, V)$  over the *n*-dimensional base manifold M with an *r*-dimensional real vector space V as typical fiber, where E is the total space and  $\pi : E \to M$  is the projection operator. An affine connection  $H_{\eta}$  in  $\eta$  is given by a special splitting  $T_{z}E = V_{z}E \oplus H_{z}E$ ,  $z \in E$  and it is determined locally by the connection coefficients  $\Gamma_{\beta}{}^{\alpha}{}_{i}(x)$ ;  $\alpha, \beta, \ldots = 1, \ldots, r$ ;  $i, j, \ldots = 1, \ldots, n$ , where  $x \in M$  has the local coordinates  $x^{i}$ .  $H_{\eta}$  or  $\Gamma$  is called *Riemann metrizable* if there exists a Euclidean scalar product  $\langle , \rangle$  in each fiber  $\pi^{-1}(x)$ , i.e. a symmetrical bilinear form g(x), in local coordinates  $\langle \xi, \zeta \rangle = g_{\alpha\beta}(x)\xi^{\alpha}(x)\zeta^{\beta}(x)$ , such that the length of the parallel translated  $\|_{\Gamma}P_{C}\xi_{0}\|_{g}$  of a vector  $\xi_{0} \in \pi^{-1}(x_{0})$  along any curve  $C(t) \subset M$ ,  $C(t_{0}) = x_{0}$  is constant, i.e. if the connection  $\Gamma$  is compatible with the Riemannian metric  $g. g(x_{0})$  is equivalent with an ellipsoid  $\mathcal{E}(\S_{1}) : \}_{\alpha\beta}(\S_{1})\xi^{\alpha}\xi^{\beta} = \infty$  in  $\pi^{-1}(x_{0})$  called *indicatrix*.  $_{\Gamma}P_{C}$  establishes a linear mapping  $\pi^{-1}(x(t_{0})) \to \pi^{-1}(x(t))$ .  $\Gamma$  is *metrizable* if there exists a field  $\mathcal{E}(\S)$  such that from  $\xi_{0} \in \mathcal{E}_{1}$  follows  $_{\Gamma}P_{C(t)}\xi_{0} \in \mathcal{E}(\S(\sqcup))$ ,  $\forall \mathcal{C}(\sqcup) \subset \mathcal{M}$ . Indicatrices play the role of the unit sphere.

The most simple case is r = n. If  $\Gamma_j{}^i{}_h(x)$  is symmetrical and metrizable by a g(x), then  $\Gamma$  is the Levi-Civita connection  $\overset{g}{\Gamma}$  of the Riemannian manifold  $V_n = (M, g)$ .

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Denoting the set of the Levi-Civita connections for the different g by  $\{\stackrel{g}{\Gamma}\}$  and supposing the symmetry  $\Gamma_j{}^i{}_h(x) = \Gamma_h{}^i{}_j(x)$  the question is whether  $\Gamma \in \{\stackrel{g}{\Gamma}\}$ . — Riemann metrizability of affine connections has been investigated by many authors from different points of view. I mention here only [1], [4], [5], [6], [9], [12].

A Finsler space  $F_n = (M, \mathcal{L})$  on the manifold M is given by the smooth fundamental function  $\mathcal{L}: \mathcal{TM} \to \mathcal{R}^+$ ;  $(x, y) \mapsto \mathcal{L}(\S, \dagger), y \in T_x M$  which is supposed to be first order positively homogeneous:  $\mathcal{L}(\S, \lambda^{\dagger}) = |\lambda| \mathcal{L}(\S, \dagger), \lambda \in \mathbb{R}$ . Its indicatrix is given by  $I(x_0) = \{y \mid \mathcal{L}(\S_I, \dagger) = \infty\} \subset \mathcal{T}_{\S_I} \mathcal{M}$  (the convexity of I is mostly also supposed). Giving of  $F_n$  is equivalent to giving of  $\{I(x)\}$ . Then an affine metrical connection should satisfy that from  $y_0 \in I(x_0)$  follows  $\Gamma P_C y_0 \in I(x_1), x_1 \in C(t_1)$ (this could be denoted by  $_{\Gamma}P_{C}I(x_{0}) = I(x_{1})$ ), while  $_{\Gamma}P_{C}$  is an affine mapping. However, this is impossible in general, e.g. if  $I(x_0)$  is an ellipsoid and  $I(x_1)$  is not so. This necessitates the introduction of the so called Finsler vector fields which are sections of a vector bundle  $\zeta = (E, \pi, TM, V^n)$ , in components  $\xi^i(x, y)$  with the property  $\xi^i(x,\lambda y) = \xi^i(x,y), \ \lambda \in \mathbb{R}, \ \lambda y \neq 0.$  The set  $\{(x_0,\lambda y_0) \mid \lambda \in \mathbb{R}, \ \lambda y_0 \neq 0\}$  is geometrically a point  $x_0$  and the direction of  $y_0$  in  $T_{x_0}M$ ; this is called a *line-element*. So Finsler vectors are defined in line-elements. The length (the norm) of such a vector is defined by  $g_{ij}(x,y)\xi^i(x,y)\xi^j(x,y) := \|\xi(x,y)\|^2$ , where  $g_{ij} := \frac{1}{2} \frac{\partial^2 \mathcal{L}^{\epsilon}}{\partial y^i \partial y^j}$  and hence  $g_{ij}(x,\lambda y) = g_{ij}(x,y)$ . In an  $F_n = (M, \mathcal{L}), g_{ij}$  is derived from  $\mathcal{L}$ . A more general structure is  $F_n = (M, z)$ . ture is  $F_n = (M, g)$ , called generalized Finsler space, where we start directly with the metric tensor  $g_{ii}(x, y)$ .

An affine connection  $\Gamma$  in the Finsler vector bundle  $\zeta$  can be given locally by the connection coefficients  $F_{j}{}^{i}{}_{k}(x,y), V_{j}{}^{i}{}_{h}(x,y)$  in the form  $\Gamma \xi = \xi - d_{\Gamma}\xi$ , where

(1) 
$$d_{\Gamma}\xi^{i}(x,y) = F_{j}^{i}{}_{k}(x,y)\xi^{j}(x,y)dx^{k} + V_{j}^{i}{}_{k}(x,y)\xi^{j}(x,y)dy^{k}.$$

 $\Gamma$  is metrizable if there exists a scalar product  $g_{ij}(x, y)$  in each  $\pi^{-1}(x, y)$  such that  $\|_{\Gamma} P_C \xi_0\| = \text{constant for any curve } C(t) \subset M.$ 

## **2** Connection in $\mu$

We want to find a new, geometric condition for the Riemann metrizability of a vector bundle  $\eta = (E, \pi, M, V^r)$  endowed with the affine connection  $H_{\eta}$  given by  $\Gamma_{\beta}{}^{\alpha}{}_i(x)$ . First we derive from  $H_{\eta}$  an affine connection  $H_{\mu}$  in  $\mu = (E_{\mu}, \pi_{\mu}, M, V^{r^2})$ , and then from  $H_{\mu}$  a connection  $H_{\nu}$  in the bundle  $\nu = (E_{\nu}, \pi_{\nu}, \pi_{\nu})$ .

 $M, \mathbf{E}$ ), where  $\mathbf{E}$  is the manifold of the ellipsoids in  $\pi^{-1}(x) \cong V^r$  centered at the origin O of  $V^r$ .

Let us consider a canonical coordinate system  $(x^i, v^{\alpha})$  in  $\pi^{-1}(U) \subset E$ , where  $U \subset M$  is a coordinate neighbourhood of  $x \in M$  and  $v^{\alpha}$  are components of  $v \in \pi^{-1} \cong V^r$ . Similarly we have local coordinates  $(x^i, y^a)$  in  $\pi_{\mu}^{-1}(U) \subset E_{\mu}$ , where  $y^a, a = 1, \ldots, r^2$  are components of  $y \in \pi_{\mu}^{-1}(x) \cong V^{r^2}$ . Let  $\overset{\alpha}{v} \in \pi^{-1}(x) \cong V^r$ ,  $\alpha, \beta = 1, \ldots, r$  be r vectors with components  $(\overset{\alpha}{v})^{\beta}$ . Since any integer a  $(1 \leq a \leq r^2)$  can uniquely be represented in the form  $a = (\alpha - 1)r + \beta$ , and conversely, any pair  $\alpha, \beta$  uniquely determines such an a and thus

(2) 
$$y^a = (\ddot{v})^{\beta}, \quad a = (\alpha - 1)r + \beta$$

determines a 1:1 mapping between  $\pi_{\mu}^{-1}(x)$  and the vector *r*-tuples  $(v, \ldots, v)$  which can be considered as elements of  $\stackrel{r}{\oplus} \pi^{-1}(x) \cong \stackrel{r}{\oplus} V^r$ .

Having an affine connection  $H_{\eta}$  in  $\eta$  with local connection coefficients  $\Gamma_{\beta}{}^{\alpha}{}_{i}(x)$ , we obtain for the parallel translated of v from x to x + dx

$$_{\Gamma}P_{x,x+dx}v(x) = v(x) - d_{\Gamma}v(x), \quad d_{\Gamma}v^{\beta}(x) = \Gamma_{\sigma}{}^{\beta}{}_{i}(x)v^{\sigma}dx^{i}.$$

Then we define an affine connection  $H_{\mu}$  in  $\mu$  with local coefficients  $G_{b}{}^{a}{}_{i}(x)$  by

(3) 
$$d_G y := (d_{\Gamma} \stackrel{1}{v}, \dots, d_{\Gamma} \stackrel{r}{v}), \qquad \begin{array}{l} y = (\stackrel{1}{v}, \dots, \stackrel{r}{v}) \\ d_{\Gamma}(\stackrel{\alpha}{v})^{\beta} = \Gamma_{\sigma}{}^{\beta}{}_i(x)(\stackrel{\alpha}{v})^{\sigma} dx^i \end{array}$$

 $G_b{}^a{}_i$  can be expressed explicitly by  $\Gamma_\beta{}^\alpha{}_i$  as follows:

(4) 
$$\begin{aligned} d_G y^a &= G_b{}^a{}_i(x) y^b dx^i \\ &= d_G y^{(\alpha-1)r+\beta} = G_{(\kappa-1)r+\lambda}{}^{(\alpha-1)r+\beta}{}_i(x) y^{(\kappa-1)r+\lambda} dx^i, \end{aligned}$$

since  $a = (\alpha - 1)r + \beta$ ,  $b = (\kappa - 1)r + \lambda$ . By (3) and (2) we get

(5) 
$$d_G y^{(\alpha-1)r+\beta} = d_{\Gamma} (\overset{\alpha}{v})^{\beta} = \Gamma_{\sigma}{}^{\beta}{}_i(x) (\overset{\alpha}{v})^{\sigma} dx^i = \Gamma_{\sigma}{}^{\beta}{}_i(x) y^{(\alpha-1)r+\sigma} dx^i.$$

From (4) and (5) we obtain

$$\begin{split} G_{(\kappa-1)r+\lambda}{}^{(\alpha-1)r+\beta}{}_i(x) \, y^{(\kappa-1)r+\lambda} &= \Gamma_{\sigma}{}^{\beta}{}_i(x) \delta^{\alpha}_{\kappa} \delta^{\sigma}_{\lambda} y^{(\kappa-1)r+\lambda} \\ &= \Gamma_{\lambda}{}^{\beta}{}_i(x) \delta^{\alpha}_{\kappa} y^{(\kappa-1)r+\lambda} \end{split}$$

and hence

$$G_{(\kappa-1)r+\lambda}{}^{(\alpha-1)r+\beta}{}_i(x) = \delta^{\alpha}_{\kappa} \Gamma_{\lambda}{}^{\beta}{}_i(x).$$

#### 3 Connection in $\nu$

An ellipsoid  $\mathcal{E}$  in  $\pi^{-1}(x) \cong V^r$  centered at the origin O of  $V^r$  has the equation  $a_{\alpha\beta}v^{\alpha}v^{\beta} = 1$ ,  $a_{\alpha\beta} = a_{\beta\alpha}$ ,  $\operatorname{Det}|a_{\alpha\beta}| > 0$ . The set  $\{\mathcal{E}\} = \mathbf{E}$  can be given a natural manifold structure, namely each  $\mathcal{E}$  can be identified with the coefficients  $a_{\alpha\beta}$  which correspond to a point of  $R^{r^2}$ . Hence  $\mathbf{E}$  can be identified with a variety of the Euclidean space  $R^{r^2}$ . Thus  $\nu = (E_{\nu}, \pi_{\nu}, B, \mathbf{E})$  is a fiber bundle.

Now we want to derive from the  $H_{\mu}$  determined by the affine connection  $H_{\eta}$  a connection  $H_{\nu}$  in  $\nu : H_{\eta} \Rightarrow H_{\mu} \Rightarrow H_{\nu}$ . — Let  $y = (\overset{1}{v}, \ldots, \overset{r}{v}) \in \pi_{\mu}^{-1}(x) \subset E_{\mu}$  be such that  $\overset{1}{v}, \ldots, \overset{r}{v}$  are linearly independent vectors in  $\pi^{-1}(x)$ . From now on, in this section y denotes elements of  $E_{\mu}$  with this independence property. The set of these (x, y)-s will be denoted by  $E_{\mu}^{*}$  and the corresponding bundle by  $\overset{*}{\mu} = (E_{\mu}^{*}, \pi_{\mu}^{*}, M, V_{*}^{r^{2}})$ . We remark that  $V_{*}^{r^{2}}$  is no vector space, and  $\pi_{*}^{*}$  is a restriction of  $\pi_{\mu}$  to  $E_{\mu}^{*} \subset E_{\mu}$ .  $H_{\mu}$  is equivalent with the splitting  $T_{u}E_{\mu} = V_{u}E_{\mu} \oplus H_{u}E_{\mu}, u \in E_{\mu}$ . The restriction

of an affine connection  $H_{\mu}$  to  $E_{\mu}^* \subset E_{\mu}$  is also a connection in  $E_{\mu}^*$ , i.e.  $H_{\mu} \subset E_{\mu}^*$  if  $u \in E_{\mu}^* \subset E_{\mu}$ . This is so, because  $H_{\eta}$  takes by parallel translation linearly independent vectors of  $\pi^{-1}(x)$  into linearly independent vectors again. Also,  $H_{\mu}^*$  can be extended by continuity to a  $H_{\mu}$ , and if  $H_{\mu}^*$  is a restriction of an affine connection  $H_{\mu}$ , then its extension yields this  $H_{\mu}$ .

The vectors  $\vec{v}$  of a y can be considered as a system of conjugate axes of an ellipsoid  $\mathcal{E} \in \pi_{\nu}^{-\infty}(\S)$  centered at the origin O, and we order this  $\mathcal{E}$  to y. Doing this with every (x, y) we obtain a strong bundle mapping

$$\rho: E_{\stackrel{*}{\mu}} \to E_{\nu}, \qquad \pi_{\stackrel{*}{\mu}}^{-1}(x) \to \pi_{\nu}^{-1}(x), \quad y \mapsto \mathcal{E}.$$

The inverse  $\rho^{-1}(\mathcal{E}) = \{\dagger_{l}, \dagger_{\infty}, \dots, \dagger_{l}, \dots\}$  is an infinite set consisting of  $y_{0} = (\overset{1}{v}_{0}, \dots, \overset{r}{v}_{0})$  $), y_{1} = (\overset{1}{v}_{1}, \dots, \overset{r}{v}_{1}), \dots, y = (\overset{1}{v}, \dots, \overset{r}{v}), \dots$  such that every system  $\overset{1}{v}_{0}, \dots, \overset{r}{v}_{0}; \overset{1}{v}_{1}, \dots, \overset{r}{v}_{1}; \dots, \overset{r}{v}; \dots, \overset{r}{v}_{1}; \dots, \overset{r}{v}; \dots, \overset{r}{v}_{1}; \dots, \overset{r}{v}; \dots, \overset{r}{v$ 

 $H_{\mu}$  takes  $\pi_{\mu}^{-1}(x)$  into  $\pi_{\mu}^{-1}(x+dx)$  and so it takes  $y \in \pi_{\mu}^{-1}(x)$  into  $\hat{y} \in \pi_{\mu}^{-1}(x+dx)$ . However, according to (3),  $H_{\mu}$  is defined via  $H_{\eta}$ , and in such a way that the images  $\hat{y}_0, \hat{y}_1, \ldots, \hat{y}, \ldots$  by  $H_{\mu}$  of the elements of an equivalence class  $\{y_0, y_1, \ldots, y, \ldots\}$  (i.e. of conjugate axes systems of an ellipsoid  $\mathcal{E}$ ) form again an equivalence class in  $\pi_{\mu}^{-1}(x+dx)$  (i.e.  $\hat{y}_0, \hat{y}_1, \ldots, \hat{y}, \ldots$  are conjugate axes systems of an ellipsoid again). This is shown on the diagram

(6)  

$$\rho(x)\{y_0, y_1, \dots, y, \dots\} = \mathcal{E}(\S) \in \pi_{\nu}^{-\infty}(\S)$$

$$\downarrow H_{\mu} \qquad \downarrow H_{\nu}$$

$$\rho(x+dx)\{\hat{y}_0, \hat{y}_1, \dots, \hat{y}, \dots\} = \hat{\mathcal{E}}(x+dx) \in \pi_{\nu}^{-1}(x+dx).$$

It means that  $H_{\mu}: \pi_{\mu}^{-1}(x) \to \pi_{\mu}^{-1}(x+dx)$  preserves equivalence classes. Thus

$$\rho \circ H_{\mu} \circ \rho^{-1} : \pi_{\nu}^{-1}(x) \to \pi_{\nu}^{-1}(x+dx)$$

yields a connection  $H_{\nu}$  in  $\nu$  (This fact is discussed in more detail in [10], [11]).

If  $H_{\nu}$  is integrable at least for one  $\mathcal{E}_{\ell} \in \pi_{\nu}^{-\infty}(\S_{\ell})$  and  $\mathcal{E}(\S), \mathcal{E}(\S_{\ell}) = \mathcal{E}_{\ell}$  is the integral manifold, then  $\mathcal{E}(\S)$  can be considered as indicatrix I(x) and  $g_{\alpha\beta}(x)$  in the equation  $g_{\alpha\beta}(x)v^{\alpha}v^{\beta} = 1$  of  $\mathcal{E}(\S)$  as metric tensor. Any  $v_0$  leading to a point of  $\mathcal{E}_{\ell} : \sqsubseteq_{\ell} \in \mathcal{E}_{\ell}$  can be an axe of a conjugate axes system of  $\mathcal{E}_{\ell}$ . Then, according to our construction, the parallel translated v of  $v_0$  according to  $H_{\eta}$  along a curve  $C \subset M$  from  $x_0$  to x is an element of  $\mathcal{E}(\S)$ :

$${}_{H_{\eta}}P_{C;x_{0},x}v_{0} = v \in {}_{H_{\nu}}P_{C;x_{0},x}\mathcal{E}_{\prime} = \mathcal{E}(\S)$$

and hence

$$||v_0||_{g(x_0)} = ||v||_{g(x)}.$$

We remark that v depends on the path C joining  $x_0$  and x, but  $\mathcal{E}(\S)$  does not. — This means: if  $H_{\nu}$  is integrable, then  $H_{\eta}$  is metrizable.

The converse is obvious. If  $H_{\eta}$  is metrical with respect to g(x), then  $\mathcal{E}(\S) := \mathcal{I}(\S)$  is an integral manifold of  $H_{\nu}$ .

Thus we obtain the

**Theorem**. The affine connection  $H_{\eta}$  of a vector bundle  $\eta$  is Riemann metrizable iff the constructed connection  $H_{\nu}$  in a bundle  $\nu$  fibered with ellipsoids is integrable.

## 4 Coefficients of $H_{\nu}$

We want to determine the connection coefficients of  $H_{\nu}$ .  $H_{\nu}$  orders to the ellipsoid  $\mathcal{E}(\S)$ 

(7) 
$$a_{\alpha\beta}(x)v^{\alpha}v^{\beta} = 1 \in \pi_{\nu}^{-1}(x)$$

the ellipsoid  $\hat{\mathcal{E}}(x+dx)$ 

(8) 
$$a_{\alpha\beta}(x+dx)v^{\alpha}(x+dx)v^{\beta}(x+dx) = 1 \in \pi_{\nu}^{-1}(x+dx).$$

According to the definition (construction) of  $H_{\nu}$  this last equation is satisfied by the parallel translated with respect to  $H_{\eta}$  of  $v^{\alpha}(x)$ , i.e. by  $v^{\alpha}(x + dx) = v^{\alpha}(x) - \Gamma_{\sigma}{}^{\alpha}{}_{i}(x)v^{\sigma}(x)dx^{i} + o(dx^{i})$ . (Since we work with linear connections,  $o(dx^{i})$ , i.e. higher order terms in  $dx^{i}$ , can be omitted.) Then the parallel translated of  $a_{\alpha\beta}(x)$  according to  $H_{\nu}$  are the  $a_{\alpha\beta}(x + dx)$  appearing in (8). Denoting the connection coefficients of  $H_{\nu}$  by  $M_{\alpha\beta i}(x, a_{\kappa\lambda})$  we obtain from (8)

$$(a_{\alpha\beta} + M_{\alpha\beta i}(x, a_{\kappa\lambda})dx^{i})(v^{\alpha} - \Gamma_{\sigma}{}^{\alpha}{}_{i}v^{\sigma}dx^{i})(v^{\beta} - \Gamma_{\sigma}{}^{\beta}{}_{i}v^{\sigma}dx^{i}) = 1$$

or

$$a_{\alpha\beta}v^{\alpha}v^{\beta} + \left[M_{\alpha\beta i} - a_{\kappa\lambda}(\Gamma_{\beta}{}^{\lambda}{}_{i}\delta^{\kappa}_{\alpha} + \Gamma_{\alpha}{}^{\kappa}{}_{i}\delta^{\lambda}_{\beta})\right]v^{\alpha}v^{\beta}dx^{i} + o(dx^{i}) = 1$$

By (7) the right hand side drops out with  $a_{\alpha\beta}v^{\alpha}v^{\beta}$ . The remaining expression must vanish for every  $v \in \mathcal{E}(\S)$  and for every  $dx^i$ . Thus, omitting  $o(dx^i)$ , we get

$$M_{\alpha\beta i}(x, a_{\kappa\lambda}) = (\Gamma_{\beta}{}^{\lambda}{}_{i}\delta^{\kappa}_{\alpha} + \Gamma_{\alpha}{}^{\kappa}{}_{i}\delta^{\lambda}_{\beta})a_{\kappa\lambda}$$

This means that  $M_{\alpha\beta i}(x, a_{\kappa\lambda})$  is linear in  $a_{\kappa\lambda}$ , i.e.  $H_{\nu}$  is an affine connection and its connection coefficients are

(9) 
$$M_{\alpha\beta}{}^{\kappa\lambda}{}_{i}(x) = \Gamma_{\alpha}{}^{\kappa}{}_{i}(x)\delta^{\lambda}_{\beta} + \Gamma_{\beta}{}^{\lambda}{}_{i}(x)\delta^{\kappa}_{\alpha}$$

We remark that these coefficients are symmetric in the sense that  $M_{\alpha\beta}{}^{\kappa\lambda}{}_i = M_{\beta\alpha}{}^{\lambda\kappa}{}_i$ . Thus the symmetry of  $a_{\alpha\beta}(x)$  implies the symmetry of  $a_{\alpha\beta}(x + dx) = a_{\alpha\beta}(x) + M_{\alpha\beta}{}^{\kappa\lambda}{}_i(x)a_{\kappa\lambda}dx^i$  too, which are the coefficients of  $\hat{\mathcal{E}}(x + dx)$ .

The condition of the absolute parallelism of  $a_{\alpha\beta}(x)$  with respect to  $H_{\nu}$  is

$$\frac{\partial a_{\alpha\beta}}{\partial x^i} = -M_{\alpha\beta}{}^{\kappa\lambda}{}_i(x)a_{\kappa\lambda}(x).$$

This is integrable iff

$$T_{\alpha\beta}{}^{\kappa\lambda}{}_{ij}(x)a_{\kappa\lambda}(x) = 0$$
$$T_{\alpha\beta}{}^{\kappa\lambda}{}_{ij} \equiv \left(\frac{\partial M_{\alpha\beta}{}^{\kappa\lambda}{}_i}{\partial x^j} - M_{\alpha\beta}{}^{\mu\nu}{}_i M_{\mu\nu}{}^{\kappa\lambda}{}_j\right)_{[i,j]}$$

has a solution for  $a_{\kappa\lambda}$  with positive determinant. We find that

$$T_{\alpha\beta}{}^{\kappa\lambda}{}_{ij} = R_{\alpha}{}^{\kappa}{}_{ij}\delta^{\lambda}_{\beta} + R_{\beta}{}^{\lambda}{}_{ij}\delta^{\kappa}_{\alpha},$$

where R is the curvature tensor of  $\Gamma_{\beta}{}^{\alpha}{}_{i}(x)$ .

## 5 Finsler vector bundles

Considering a Finsler vector bundle  $\zeta = (E, \pi, TM, V^n)$  and a connection  $\Gamma$  with connection coefficients  $F_{j}{}^{i}{}_{h}(x, y), V_{j}{}^{i}{}_{h}(x, y)$  we have (1). In this case the base manifold TM has dimension 2n. Its coordinates can be denoted by  $u^A$ ,  $A = 1, \ldots, 2n$ ;  $u^i = x^i$ ,  $u^{n+i} = y^i$ .  $\mathcal{E}(\S, \dagger)$  has the equation  $a_{ij}(x, y)\xi^i\xi^j = 1$ , and the equation of  $\hat{\mathcal{E}}(x + dx)$  is

$$a_{ij}(x + dx, y + dy)\xi^{i}(x + dx, y + dy)\xi^{j}(x + dx, y + dy) = 1.$$

Here

$$a_{ij}(x + dx, y + dy) = a_{ij}(x) + M_{ij}{}^{rs}{}_{h}(x, y)a_{rs}(x, y)dx^{h} + M_{ij}{}^{rs}{}_{n+k}(x, y)a_{rs}dy^{h}.$$

Contrasting with (9), here the last index of M runs from 1 to 2n the other indices from 1 to n. Considerations and calculations similar to those done above yield

$$M_{ij}{}^{rs}{}_{h} = F_{j}{}^{s}{}_{h}\delta^{r}_{i} + F_{i}{}^{r}{}_{h}\delta^{s}_{j}$$
$$M_{ij}{}^{rs}{}_{n+k} = V_{j}{}^{s}{}_{k}\delta^{r}_{i} + V_{i}{}^{r}{}_{k}\delta^{s}_{j},$$

and furthermore

$$T_{ij}{}^{rs}{}_{kh} = {}^{F}R_{i}{}^{r}{}_{kh}\delta_{j}^{s} + {}^{F}R_{j}{}^{s}{}_{kh}\delta_{i}^{r}$$
$$T_{ij}{}^{rs}{}_{n+k\ n+h} = {}^{V}R_{i}{}^{r}{}_{kh}\delta_{j}^{s} + {}^{V}R_{j}{}^{s}{}_{kh}\delta_{i}^{r},$$

where  ${}^{F}R$  and  ${}^{V}R$  are formed from  $F_{j}{}^{s}{}_{i}$  and  $V_{j}{}^{s}{}_{i}$  resp. like common curvature tensors. Finally

$$T_{ij}{}^{rs}{}_{n+h\ k} = \frac{\partial M_{ij}{}^{rs}{}_{n+h}}{\partial x^k} - \frac{\partial M_{ij}{}^{rs}{}_k}{\partial y^h} + (V_j{}^s{}_kF_s{}^c{}_h - F_j{}^s{}_kV_s{}^c{}_h)\delta^b_i + V_j{}^c{}_kF_i{}^b{}_h - F_j{}^c{}_kV_i{}^b{}_h + V_i{}^b{}_kF_j{}^c{}_h - F_i{}^b{}_kV_j{}^c{}_h + (V_i{}^r{}_kF_r{}^b{}_h - F_i{}^r{}_kV_r{}^b{}_h)\delta^c_j.$$

One can use other connections, e.g. a pre-Finsler connection  $F\Gamma(F_j{}^i{}_k, N^i{}_j, V_j{}^i{}_h)$ and *h*- and *v*-covariant derivatives

$$\xi^{i}{}_{|k} = \frac{\partial \xi^{i}}{\partial x^{k}} - \frac{\partial \xi^{i}}{\partial y^{r}} N^{r}{}_{k} + F_{j}{}^{i}{}_{k}\xi^{j}$$
$$\xi^{i}{}_{|k} = \frac{\partial \xi^{i}}{\partial y^{k}} + V_{j}{}^{i}{}_{k}\xi^{j}.$$

In this case (1) becomes

$$d_{\Gamma}\xi^{i} = (F_{j}{}^{i}{}_{k} - V_{j}{}^{i}{}_{r}N^{r}{}_{k})\xi^{j}dx^{k} + V_{j}{}^{i}{}_{k}\xi^{j}dy^{k},$$

or

$$d_{\Gamma}\xi^{i} = \left[ (F_{j}{}^{i}{}_{k} - V_{j}{}^{i}{}_{r}F_{s}{}^{r}{}_{k}y^{s})dx^{k} + V_{j}{}^{i}{}_{k}dy^{k} \right]\xi^{j}$$

if  $F\Gamma$  is without deflection. These lead to other formulae for  $M_{ij}{}^{rs}{}_A$  and  $T_{ij}{}^{rs}{}_{AB}$ . If  $F_j{}^i{}_k$  and  $V_j{}^i{}_k$  are symmetric,  $F\Gamma$  is without deflection and metrizable, then  $F\Gamma$  is the Cartan connection.

Finally we mention still another affine connection introduced by M. Matsumoto [7], [8] (see also [2], [3]) which is an ordinary affine connection derived from a Finsler connection  $F\Gamma(F_j{}^i{}_k, N^i{}_j, V_j{}^i{}_k)$ . Starting with an  $F\Gamma$  and a nonvanishing vector field Y(x) which depends on the point x only

(10) 
$$\underline{\Gamma}_{j\,k}^{i\,k}(x) := F_{j\,k}^{i\,k}(x,Y(x)) + V_{j\,r}^{i\,r}(x,Y(x)) \left(\frac{\partial Y^{r}}{\partial x^{k}} + Y^{s}(x)F_{s\,r\,k}^{r\,r}(x,Y(x))\right)$$

turn out to be connection coefficients of an ordinary affine connection. Using the vector field Y(x) one can associate to any Finsler vector field  $\xi^i(x, y)$  an ordinary vector field  $\underline{\xi}^i(x) := \xi^i(x, Y(x))$ . Then there exists a nice relation among the covariant derivative  $\underline{\xi}^i_{;k}$  constructed with  $\underline{\Gamma}$ , and the *h*- and *v*-covariant derivatives with respect to  $F\Gamma$ , namely

$$\underline{\xi}^{i}_{;k} = \left[\xi^{i}_{|k} + \xi^{i}_{|k} \left(\frac{\partial Y^{r}}{\partial x^{k}} + Y^{s} F_{s}^{r}_{k}\right)\right]_{|_{y=Y(x)}}$$

Given a  $\underline{\Gamma}$  and a Y(x), there are many  $F\Gamma$  which satisfy (10). Then we can use our method to search for metrizable ones among these  $F\Gamma$ , e.g. for such, where  $F\Gamma$ satisfies (10) with the given  $\underline{\Gamma}$  and Y(x) and  $g_{ij|k} = g_{ij}|_{k} = 0$  with respect to this  $F\Gamma$ .

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