Electronic Journal of Differential Equations, Vol. 2023 (2023), No. 13, pp. 1–15. ISSN: 1072-6691. URL: https://ejde.math.txstate.edu or https://ejde.math.unt.edu

EXISTENCE OF NONTRIVIAL SOLUTIONS FOR SCHRÖDINGER-KIRCHHOFF EQUATIONS WITH INDEFINITE POTENTIALS

SHUAI JIANG, LI-FENG YIN

ABSTRACT. We consider a class of Schrödinger-Kirchhoff equations in \mathbb{R}^3 with a general nonlinearity g and coercive sign-changing potential V so that the Schrödinger operator $-a\Delta + V$ is indefinite. The nonlinearity considered here satisfies the Ambrosetti-Rabinowitz type condition $g(t)t \ge \mu G(t) > 0$ with $\mu > 3$. We obtain the existence of nontrivial solutions for this problem via Morse theory.

1. INTRODUCTION

In this article, we consider the Schrödinger-Kirchhoff type problem

$$-\left(a+b\int_{\mathbb{R}^3}|\nabla u|^2\mathrm{d}x\right)\Delta u+V(x)u=g(u),\quad u\in H^1(\mathbb{R}^3),\tag{1.1}$$

where a and b are positive constants, $V \in C(\mathbb{R}^3)$ is the potential and $g \in C(\mathbb{R})$ is the nonlinearity.

The above problem is nonlocal as the term $\int_{\mathbb{R}^3} |\nabla u|^2 dx$ implies that (1.1) is not a pointwise identity. This feature causes some mathematical difficulties, which make the investigation of (1.1) particularly interesting. Problem (1.1) arises in an interesting physical context. Indeed, if we set V(x) = 0 and replace \mathbb{R}^3 by a bounded domain $\Omega \subset \mathbb{R}^3$, problem (1.1) reduces to the Dirichlet problem

$$-\left(a+b\int_{\Omega}|\nabla u|^{2}\mathrm{d}x\right)\Delta u = g(u) \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \partial\Omega,$$

(1.2)

which is related to the stationary analogue of the equation

$$\rho \frac{\partial^2 u}{\partial t^2} - \left(\frac{P_0}{h} + \frac{E}{2L} \int_0^L |\frac{\partial u}{\partial x}|^2 \mathrm{d}x\right) \frac{\partial^2 u}{\partial x^2} = 0, \tag{1.3}$$

proposed by Kirchhoff [10] as an extension of the classical D'Alembert's wave equation for free vibrations of elastic strings, where ρ , P_0 , h, E and L are positive

²⁰²⁰ Mathematics Subject Classification. 35B38, 35J20, 35J60.

Key words and phrases. Schrödinger-Kirchhoff equations; Palais-Smale condition; Morse theory.

^{©2023.} This work is licensed under a CC BY 4.0 license.

Submitted March 28, 2022. Published February 10, 2023.

constants. Lions [15] introduced an abstract functional analysis framework for the equation

$$u_{tt} - \left(a + b \int_{\Omega} |\nabla u|^2 \mathrm{d}x\right) \Delta u = g(x, u).$$
(1.4)

Since then, problem (1.4) have received much attention.

Problem (1.1) has been studied extensively in recent years (in bounded or unbounded domain), see, for example, [1, 7, 14, 22, 23, 26, 28] and references therein. We emphasize that in all these papers, the potential V is assumed to be nonnegative. In this case, the quadratic part of the variational functional Φ given in (2.1) is positively definite, the zero function u = 0 is a local minimizer of Φ and the mountain pass theorem [2] can be applied. However, when the potential V is negative somewhere so that the quadratic part of Φ is indefinite, the zero function u = 0is no longer a local minimizer of Φ , the mountain pass theorem is not applicable anymore. For the semilinear problem (b = 0)

$$\Delta u + V(x)u = g(u)$$

with indefinite Schrödinger operator $-\Delta + V$, one usually applies the linking theorem to obtain a solution, see e.g. [11, 19]. For problem (1.1), it seems hard to verify the linking geometry because of the nonlocal term $\int_{\mathbb{R}^3} |\nabla u|^2 dx$, which prevents the functional Φ to be nonpositive on the negative space of the Schrödinger operator, see Remark 1.2 for detail. Hence, the classical linking theorem [25, Lemma 2.12] is also not applicable.

For this reason, there are very few results about (1.1) with indefinite potential. To the best of our knowledge, the first work on this situation is due to Chen and Liu [6]. To overcome the above difficulty, a crucial observation in [6] is that Φ has a local linking at u = 0, thus for certain $i \in \mathbb{N}$, the *i*-th critical group of Φ at u = 0 is nontrivial. Assuming that g is 3-superlinear

$$\lim_{t \to \infty} \frac{g(t)}{t^3} = +\infty, \tag{1.5}$$

with suitable technical conditions it was shown in [6] that all critical groups of Φ at infinity are trivial, then a nonzero critical point of Φ can be found via Morse theory. See also [13] for a recent result for 3-superlinear nonlinearity.

For other results on the indefinite problem (1.1), we mention [27] and [9], where the nonlinearity g is sublinear and subquadratic, respectively. In these two papers the variational functional is coercive. The three critical points theorem of Liu and Su [16] and the classical Clark theorem are applied to obtain multiple solutions.

Having the above results in mind, it is natural to ask what will happen when g is 2-superlinear (g satisfies (1.5) with 3 replaced by 2)? This is the motivation of the current paper.

The main difficulty under this assumption is that it is not know whether Palais-Smale (or even Cerami) sequences are bounded or not. Motivated by Liu and Mosconi [20] on the study of nonlinear Schrödinger-Poisson systems (see also [8]), we add a dummy variable and consider an augmented functional $\tilde{\Phi} : \mathbb{R} \times X \to \mathbb{R}$, see (2.4). It turns out that $\tilde{\Phi}$ satisfies the (*PS*) condition (see Lemma 2.6), and if (\bar{s}, \bar{u}) is a critical point of $\tilde{\Phi}$, then \bar{u} is a critical point of Φ (see Lemma 2.5). Moreover, $\tilde{\Phi}$ has a local linking at zero, and all critical groups of $\tilde{\Phi}$ at infinity are trivial. Eventually, a nonzero critical point of $\tilde{\Phi}$.

Without loss of generality, we assume that a = b = 1. Then the problem (1.1) can be rewritten as follows

$$-\left(1+\int_{\mathbb{R}^3} |\nabla u|^2 \mathrm{d}x\right) \Delta u + V(x)u = g(u), \quad u \in H^1(\mathbb{R}^3).$$
(1.6)

Let $H^1(\mathbb{R}^3)$ be the standard Sobolev space. If $V \in C(\mathbb{R}^3)$ is bounded from below, we can choose a constant m > 2 such that

$$\tilde{V}(x) := V(x) + m \ge \frac{m}{2} \ge 1, \quad \text{for all} \quad x \in \mathbb{R}^3.$$
(1.7)

On the linear subspace

$$X = \{ u \in H^1(\mathbb{R}^3) : \int V(x) u^2 < \infty \},$$
(1.8)

where from now on all integrals are taken over \mathbb{R}^3 except stated explicitly, we equip with the inner product

$$(u,v) = \int \left(\nabla u \cdot \nabla v + \tilde{V}(x) u v \right)$$

and the corresponding norm $\|\cdot\|$. Note that if the assumptions (A1) below holds, then X is a Hilbert space and by Bartsch and Wang [4] we have a compact embedding $X \hookrightarrow L^s(\mathbb{R}^3)$ for $s \in [2, 6)$.

Now we present our assumptions on the potential V(x) and the nonlinearity g(u).

- (A1) $V \in C(\mathbb{R}^3)$ is bounded from below and $|\{V \leq k\}| < \infty$ for all $k \in \mathbb{R}$, where $|\cdot|$ is the Lebesgue measure on \mathbb{R}^3 .
- (A2) $V \in C^1(\mathbb{R}^3)$ and there exists R > 0 such that $V(x) \nabla V(x) \cdot x \ge 0$ for $|x| \ge R$.
- (A3) There exists $\kappa > 0$ such that

$$|\nabla V(x) \cdot x| \le \kappa \left(V(x) + m \right) =: \kappa \tilde{V}(x)$$

for all $x \in \mathbb{R}^3$, here *m* is the constant from (1.7).

(A4) $g \in C(\mathbb{R})$ and there exist C > 0 and $p \in (2, 6)$ such that

$$|g(t)| \le C \left(|t| + |t|^{p-1} \right)$$

(A5) There exists $\mu > 3$ such that $g(t)t \ge \mu G(t) > 0$ for all $t \in \mathbb{R} \setminus \{0\}$.

Consider the bilinear form

$$Q(u,v) = \frac{1}{2} \int \left(\nabla u \cdot \nabla v + V(x) u v \right) \quad u, v \in X,$$
(1.9)

then $X = X^+ \oplus X^- \oplus X^0$, where X^+ , X^- and X^0 are positive, negative, and null eigenspaces of the Schrödinger operator. It is well known that there exist constants $\eta^{\pm} > 0$ such that

$$\pm Q(u,u) \ge \eta^{\pm} ||u^{\pm}||^2, \quad \text{for } u \in X^{\pm} \oplus X^0, \text{ respectively.}$$
(1.10)

We are now ready to state our main result.

Theorem 1.1. Suppose that (A1)–(A5) hold. If either

(1) $\dim X^- > 0$, $\dim X^0 = 0$, or

(2) dim $X^0 > 0$ and $G(t) = \int_0^t g(\tau) d\tau \ge c|t|^{\nu}$ for some $\nu < 4$,

then problem (1.6) has at least one nontrivial solution.

Remark 1.2. As we have mentioned, because of the nonlocal term $\int_{\mathbb{R}^3} |\nabla u|^2 dx$, our functional Φ (see (2.1)) does not satisfy the geometric assumption of the linking theorem. For R > r > 0 set

$$N = \{ u \in X^+ : ||u|| = r \}, \quad M = \{ u \in X^- \oplus X^0 \oplus \phi \mathbb{R}^+ : ||u|| \le R \},\$$

where $\phi \in X^+ \setminus \{0\}$. Then we have

$$b = \underset{N}{\inf} \Phi > 0, \quad \underset{u \in \partial M, \|u\| = R}{\sup} \Phi < 0$$

provided R is large enough and r is small enough. However, because the integral $(\int |\nabla u|^2)^2$ in our functional, Φ may be very large, it is possible that $\Phi(u) > b$ for some $u \in \partial M \cap (X^- \oplus X^0)$. Therefore, the following geometric assumption of the linking theorem

$$\inf_{N} \Phi > \sup_{\partial M} \Phi$$

can not be satisfied. Thus, the linking theorem is not applicable for problem (1.6).

2. PALAIS-SMALE CONDITION

Now, we investigate the functional Φ . Under the assumptions (A1) and (A4) we can show that the functional $\Phi: X \to \mathbb{R}$,

$$\Phi(u) = \frac{1}{2} \int \left(|\nabla u|^2 + V(x)u^2 \right) + \frac{1}{4} \left(\int |\nabla u|^2 \right)^2 - \int G(u)$$
(2.1)

is well defined and of class C^1 . The derivative of Φ is given by

$$\langle D\Phi(u), v \rangle = \left(1 + \int |\nabla u|^2\right) \int \nabla u \cdot \nabla v + \int V(x)uv - \int g(u)v.$$

for $u, v \in X$. Consequently, critical points of Φ are weak solutions of problem (1.6).

Proposition 2.1 ([20, Lemma 3.1]). Assume (A3) holds and let $\tilde{V} := V + m$. Then for any t > 0 and $x \in \mathbb{R}^3$, we have

$$\tilde{V}(tx) \le \max\{t^{\kappa}, t^{-\kappa}\}\tilde{V}(x).$$
(2.2)

For t > 0 and $u \in X$ we define

$$u_t(x) = tu(\frac{x}{t}),\tag{2.3}$$

and define on the Hilbert space $\mathbb{R}\times X$ (with natural norm $\|(s,u)\|^2=s^2+\|u\|^2)$ the augmented functional

$$\tilde{\Phi}(s,u) := \frac{s^2}{2} + \Phi(u_{e^s}).$$
(2.4)

Remark 2.2. Obviously, for s, t > 0, from (2.3) we have $(u_t)_s = u_{ts} = (u_s)_t$.

Lemma 2.3. Suppose that (A3) and (A4) hold. Then the functional $\tilde{\Phi}$ is well defined on $\mathbb{R} \times X$, of class C^1 , and

$$\tilde{\Phi}(s,u) = \frac{s^2}{2} + \frac{e^{3s}}{2} \int |\nabla u|^2 + \frac{e^{6s}}{4} \left(\int |\nabla u|^2 \right)^2 + \frac{e^{5s}}{2} \int V(xe^s) u^2 - e^{3s} \int G(e^s u),$$
(2.5)

$$\begin{split} \langle \partial_u \tilde{\Phi}(s,u), \varphi \rangle &= e^{3s} \int \nabla u \nabla \varphi + e^{6s} \left(\int |\nabla u|^2 \right) \int \nabla u \nabla \varphi \\ &+ e^{5s} \int V(xe^s) u\varphi - e^{4s} \int g(e^s u) \varphi, \\ \partial_s \tilde{\Phi}(s,u) &= s + \frac{3e^{3s}}{2} \int |\nabla u|^2 + \frac{3e^{6s}}{2} \left(\int |\nabla u|^2 \right)^2 \\ &+ \frac{e^{5s}}{2} \int \left[5V(xe^s) + \nabla V(xe^s) \cdot xe^s \right] u^2 \\ &- e^{3s} \int \left[3G(e^s u) + g(e^s u)e^s u \right]. \end{split}$$
(2.6)

Proof. By changing variables, we obtain

$$\left(\int |\nabla u_t|^2\right)^2 = t^6 \left(\int |\nabla u|^2\right)^2 \le t^6 ||u||^4.$$

Similarly,

$$\int |\nabla u_t|^2 = t^3 \int |\nabla u|^2 \le t^3 ||u||^2,$$
$$\left| \int G(u_t) \right| \le C |t|^5 \int u^2 + C |t|^{p+3} \int |u|^p.$$

By the continuity of V, changing variables x = yt on any fixed ball B_R and applying Proposition 2.1, we obtain

$$\int_{B_R} \tilde{V}(x) u_t^2 = t^5 \int_{B_{R/t}} \tilde{V}(tx) u^2 \le t^5 \max\{t^{\kappa}, t^{-\kappa}\} \int_{B_{R/t}} \tilde{V}(x) u^2 \le c_t \|u\|^2.$$

Letting $R \to +\infty$ we deduce that

$$\int \tilde{V}(x)u_t^2 \le c_t \|u\|^2.$$

From the above estimates and continuous embeddings,

$$\begin{split} \Phi(u_t) &= \frac{1}{2} \int \left(|\nabla u_t|^2 + V(x)u_t^2 \right) + \frac{1}{4} \left(\int |\nabla u_t|^2 \right)^2 - \int G(u_t) \\ &\leq \frac{t^3}{2} \|u\|^2 + \frac{c_t}{2} \|u\|^2 + \frac{t^6}{4} \|u\|^4 + C|t|^5 \|u\|^2 + C|t|^{p+3} \|u\|^p. \end{split}$$

Together with (2.4), it is easy to see that $\tilde{\Phi}$ is well defined. Formula (2.5) directly follows from changing variables.

Formula (2.6) can be computed in a standard way, while (2.7) is obtained by differentiating under the integral sign in (2.5). By (A4), we deduce that

$$|3G(e^{s}u) + g(e^{s}u)e^{s}u| \le C(e^{2s}|u|^{2} + e^{ps}|u|^{6}).$$
(2.8)

It follows from $|V| \leq \tilde{V} + m$, (A3) and Proposition 2.1 that

$$\begin{aligned} |5V(xe^s) + \nabla V(xe^s) \cdot xe^s| &\leq 5|V(xe^s)| + |\nabla V(xe^s) \cdot xe^s| \\ &\leq 5m + (5+\kappa)\tilde{V}(xe^s) \\ &\leq 5m + (5+\kappa)e^{\kappa|s|}\tilde{V}(x). \end{aligned}$$
(2.9)

Noting that, by the dominated convergence theorem, both

$$s \mapsto \int |5V(xe^s) + \nabla V(xe^s) \cdot xe^s|, \quad s \mapsto \int |3G(e^su) + g(e^su)e^su|$$

are continuous, standard arguments yields the differentiation formula (2.7). Finally, estimates (2.8) and (2.9) ensure the continuity of the corresponding Nemitskii operators appearing in (2.6) and (2.7), so that $\tilde{\Phi}$ is of class C^1 .

Lemma 2.4 (Pohožaev identity). Suppose that (A3) and (A4) hold. Let $u \in X$ be a weak solution to problem (1.6), then we have the following Pohožaev identity,

$$\frac{1}{2}\int |\nabla u|^2 + \frac{3}{2}\int V(x)u^2 + \frac{1}{2}\int (\nabla V(x), x)u^2 + \frac{1}{2}\left(\int |\nabla u|^2\right)^2 - 3\int G(u) = 0.$$
(2.10)

Moreover,

$$\frac{\mathrm{d}}{\mathrm{d}t}\big|_{t=1}\Phi(u_t) = 0.$$

Proof. The proof of Pohožaev identity is standard, so we omit it and refer the reader to [12, Lemma 2.1]. Now, it suffices to show that

$$\frac{\mathrm{d}}{\mathrm{d}t}\big|_{t=1}\Phi(u_t) = 0.$$

By changing variables and (2.1), we have

$$\Phi(u_t) = \frac{t^3}{2} \int |\nabla u|^2 + \frac{t^5}{2} \int V(xt)u^2 + \frac{t^6}{4} \left(\int |\nabla u|^2\right)^2 - t^3 \int G(tu).$$
(2.11)

Then, by (2.10), (2.11), and $D\Phi(u) = 0$, we obtain

$$\begin{aligned} \frac{\mathrm{d}\Phi(u_t)}{\mathrm{d}t}\Big|_{t=1} &= \frac{3}{2} \int |\nabla u|^2 + \frac{5}{2} \int V(x)u^2 + \frac{1}{2} \int (\nabla V(x), x)u^2 \\ &\quad + \frac{3}{2} \Big(\int |\nabla u|^2\Big)^2 - \int [3G(u) + g(u)u] \\ &= \langle D\Phi(u), u \rangle + \frac{1}{2} \int |\nabla u|^2 + \frac{3}{2} \int V(x)u^2 \\ &\quad + \frac{1}{2} \int (\nabla V(x), x)u^2 + \frac{1}{2} \Big(\int |\nabla u|^2\Big)^2 - 3 \int G(u) = 0. \end{aligned}$$

In the following, we denote by $\tilde{D}\tilde{\Phi}$ the total differential of $\tilde{\Phi}$ with respect to both variables s and u.

Lemma 2.5. Suppose that (A3) and (A4) hold, then

$$D\bar{\Phi}(\bar{s},\bar{u}) = 0 \quad \Leftrightarrow \quad \bar{s} = 0 \text{ and } D\Phi(\bar{u}) = 0.$$

Proof. The proof is similar to the proof of [20, Lemma 3.5], we include it here for the reader's convenience.

 (\Leftarrow) From the definition of $D\Phi$, we have

$$\tilde{D}\tilde{\Phi}(\bar{s},\bar{u})[l,\varphi] = \langle \partial_u \tilde{\Phi}(\bar{s},\bar{u}),\varphi \rangle + \partial_s \tilde{\Phi}(\bar{s},\bar{u})l, \quad \text{for all } [l,\varphi] \in \mathbb{R} \times X.$$
(2.12)

Note that $D\Phi(\bar{u}) = 0$ implies

$$\langle \partial_u \Phi(0, \bar{u}), \varphi \rangle = \langle D \Phi(\bar{u}), \varphi \rangle = 0.$$

Therefore, it suffices to prove that $\partial_s \tilde{\Phi}(0, \bar{u}) = 0$. From $D\Phi(\bar{u}) = 0$, Lemma 2.4 gives

$$\frac{\mathrm{d}}{\mathrm{d}t}\big|_{t=1}\Phi(\bar{u}_t) = 0. \tag{2.13}$$

The map $t \mapsto \Phi(\bar{u}_t)$ is C^1 by Lemma 2.3 and

$$\Phi(\bar{u}_t) = \tilde{\Phi}(\log t, \bar{u}) - \frac{\log^2 t}{2}.$$

Now (2.13) gives

$$0 = \frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=1} \Big(\tilde{\Phi}(\log t, \bar{u}) - \frac{\log^2 t}{2}\Big) = \Big(\partial_s \tilde{\Phi}(\log t, \bar{u}) \frac{1}{t} - \frac{\log t}{t}\Big)\Big|_{t=1} = \partial_s \tilde{\Phi}(0, \bar{u}).$$

 (\Rightarrow) If $\tilde{D}\tilde{\Phi}(\bar{s},\bar{u})=0$, combining (2.4), we immediately infer that

$$0 = \partial_u \tilde{\Phi}(\bar{s}, \bar{u}) = \partial_u \Phi(\bar{u}_{e^{\bar{s}}}) = D\Phi(\bar{u}_{e^{\bar{s}}}).$$

Thus, we only have to prove that $\bar{s} = 0$. Applying Lemma 2.4 to $v := \bar{u}_{e^{\bar{s}}}$ gives

$$\frac{\mathrm{d}}{\mathrm{d}t}\big|_{t=1}\Phi(v_t) = 0.$$

The map $t \mapsto \Phi(v_t)$ is C^1 by Lemma 2.3 and

$$\Phi(v_t) = \Phi(\bar{u}_{te^{\bar{s}}}) = \tilde{\Phi}(\bar{s} + \log t, \bar{u}) - \frac{(\bar{s} + \log t)^2}{2}.$$

By the chain rule and $\partial_s \tilde{\Phi}(\bar{s},\bar{u})=0$ we have

$$0 = \frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=1} \Phi(v_t)$$

= $\left(\partial_s \tilde{\Phi}(\bar{s} + \log t, \bar{u}) \frac{1}{t} - \frac{(\bar{s} + \log t)}{t}\right)\Big|_{t=1}$
= $\partial_s \tilde{\Phi}(\bar{s}, \bar{u}) - \bar{s} = -\bar{s}.$

Lemma 2.6. Suppose that (A1)–(A5) hold. Then $\tilde{\Phi}$ satisfies the (PS) condition.

Proof. Let $\{(s_n, u_n)\}$ be a (PS) sequence for $\tilde{\Phi}$ in $\mathbb{R} \times X$, that is $|\tilde{\Phi}(s_n, u_n)| \leq s_n = \tilde{D}\tilde{\Phi}(s_n, u_n) > 0$

$$\sup_{n} |\tilde{\Phi}(s_n, u_n)| < \infty, \quad \tilde{D}\tilde{\Phi}(s_n, u_n) \to 0.$$

Choose $\lambda \in (3, \mu)$, where $\mu > 3$ is given by (A5). Then

$$\begin{aligned} &(\lambda+3)\Phi(s,u) - \partial_s \Phi(s,u) \\ &= \frac{\lambda+3}{2}s^2 - s + \frac{\lambda}{2}e^{3s} \int |\nabla u|^2 + \frac{\lambda-3}{4}e^{6s} \Big(\int |\nabla u|^2\Big)^2 \\ &+ \frac{e^{5s}}{2} \int \left[(\lambda-2)V(xe^s) - \nabla V(xe^s) \cdot xe^s \right] u^2 + e^{3s} \int \left[g(e^su)e^su - \lambda G(e^su) \right] \\ &\geq \frac{\lambda+3}{2}s^2 - s + \frac{\lambda}{2}e^{3s} \int |\nabla u|^2 + \frac{\lambda-3}{4}e^{6s} \Big(\int |\nabla u|^2\Big)^2 \\ &+ \frac{e^{5s}}{2} \int \left[(\lambda-2)V(xe^s) - \nabla V(xe^s) \cdot xe^s \right] u^2 + (\mu-\lambda)e^{3s} \int G(e^su). \end{aligned}$$

$$(2.14)$$

The third integral is bounded from below through (A1)–(A3), and Hölder's inequality. Indeed, we set

$$v(x) = e^s u\left(\frac{x}{e^s}\right), \quad W_\lambda(x) = (\lambda - 2)V(x) - \nabla V(x) \cdot x.$$

Then, by changing variables,

$$e^{5s} \int [(\lambda - 2)V(xe^s) - \nabla V(xe^s) \cdot xe^s] u^2 = \int W_{\lambda} v^2.$$
 (2.15)

As W_{λ} is bounded on bounded sets, we let $C_{\lambda} \in \mathbb{R}$ be such that $W_{\lambda} \geq -C_{\lambda}$ in B_R , where R is given in (A2), then

$$\int_{B_R} W_{\lambda} v^2 \ge -C_{\lambda} \int_{B_R} v^2 \ge -C_{\lambda} \Big(\int |v|^{\mu}\Big)^{2/\mu}.$$
(2.16)

We write the integral over $\mathbb{R}^3 \setminus B_R$ as the sum of the integrals over the intersections of $\mathbb{R}^3 \setminus B_R$ with $\{V \ge 0\}$ and $\{V < 0\}$. Obviously, assumption (A1) implies that $\{V < 0\}$ has finite measure. Because $\lambda - 2 > 1$, assumption (A2) implies that

$$W_{\lambda}(x) = (\lambda - 2)V(x) - \nabla V(x) \cdot x \ge V(x) - \nabla V(x) \cdot x \ge 0$$

for $x \in \{V \ge 0\} \setminus B_R$.

On the other hand, by (A3) and $V \ge -m$ we have

$$W_{\lambda}(x) \ge (\lambda - 2)V - \kappa(V + m) \ge -((\lambda - 2) + \kappa)m \quad \text{on } \{V < 0\}.$$

For some possibly larger C_{λ} , we have

$$\int_{\mathbb{R}^3 \setminus B_R} W_{\lambda} v^2 \ge \int_{\{V < 0\} \setminus B_R} W_{\lambda} v^2 \ge -C_{\lambda} \int_{\{V < 0\}} v^2$$

$$\ge -C_{\lambda} |\{V < 0\}|^{(\mu-2)/\mu} \left(\int |v|^{\mu}\right)^{2/\mu}.$$
(2.17)

Combining (2.15), (2.16), and (2.17), we obtain

$$e^{5s} \int \left[(\lambda - 2)V(xe^s) - \nabla V(xe^s) \cdot xe^s \right] u^2 \ge -C_\lambda \left(\int |v|^\mu \right)^{2/\mu}.$$
 (2.18)

Condition (A5) implies that for some C > 0, $G(t) \ge C|t|^{\mu}$. We deduce

$$(\mu - \lambda)e^{3s} \int G(e^s u) = (\mu - \lambda) \int G(v) \ge (\mu - \lambda)C \int |v|^{\mu}.$$
 (2.19)

From (2.14), (2.18), and (2.19), for our (PS) sequence $\{(s_n, u_n)\}$, we have

$$\begin{split} O(1) &\geq (\lambda+3)\tilde{\Phi}(s_n, u_n) - \partial_s \tilde{\Phi}(s_n, u_n) \\ &\geq \frac{\lambda+3}{2} s_n^2 - s_n + \frac{\lambda}{2} e^{3s_n} \int |\nabla u_n|^2 + \frac{\lambda-3}{4} e^{6s_n} \Big(\int |\nabla u_n|^2 \Big)^2 \\ &+ (\mu - \lambda) e^{3s_n} \int G(e^{s_n} u_n) + \frac{e^{5s_n}}{2} \int [(\lambda - 2)V(xe^{s_n}) - \nabla V(xe^{s_n}) \cdot xe^{s_n}] u_n^2 \\ &\geq \frac{\lambda+3}{2} s_n^2 - s_n + \frac{\lambda}{2} e^{3s_n} \int |\nabla u_n|^2 + \frac{\lambda-3}{4} e^{6s_n} \Big(\int |\nabla u_n|^2 \Big)^2 \\ &+ (\mu - \lambda)C \int |v_n|^\mu - C_\lambda \Big(\int |v_n|^\mu \Big)^{2/\mu}, \end{split}$$

where $v_n(x) = e^s u_n\left(\frac{x}{e^s}\right)$. From $\mu > \lambda$, $2/\mu < 1$ we infer $(\mu - \lambda)C\xi_n - C_\lambda {\xi_n}^{2/\mu} \to +\infty$ if $\xi_n = |v_n|^2_{\mu} \to +\infty$, we deduce from the previous estimate that

$$|s_n|, \quad \int |\nabla u_n|^2, \quad \int |v_n|^{\mu}, \quad \int G(e^{s_n}u_n) \quad \text{are bounded};$$
 (2.20)

recalling (2.5) and $\Phi(s_n, u_n) = O(1)$ we obtain

$$\int V(xe^{s_n})u_n^2 \le O(1). \tag{2.21}$$

To complete the proof of the boundedness of $||u_n||$, let S > 0 be such that $|s_n| \leq S$. Since $V \geq -m$ on $\{V \leq k\}$, we deduce

$$\begin{split} \int V(xe^{s_n})u_n^2 &= e^{-2s_n} \int V(xe^{s_n})v_n^2(xe^{s_n}) = e^{-5s_n} \int V(x)v_n^2 \\ &\geq e^{-5S} \Big(\int_{\{V > k\}} V(x)v_n^2 - m \int_{\{V \le k\}} v_n^2 \Big) \\ &\geq e^{-5S} \Big(\int_{\{V > k\}} V(x)v_n^2 - m |\{V \le k\}|^{(\mu-2)/\mu} \Big(\int |v_n|^{\mu} \Big)^{2/\mu} \Big) \\ &\geq e^{-5S} \int_{\{V > k\}} V(x)v_n^2 - O(1), \end{split}$$

where we used (A1) and (2.20) in the last inequality. From (2.21) we thus infer

$$\int_{\{V > k\}} V(x) v_n^2 \le O(1).$$

Next, k > m implies $V + m \le 2V$ on $\{V > k\}$ and $V + m \le 2k$ on $\{V \le k\}$. Using Proposition 2.1 we obtain

$$\begin{split} \|u_n\|^2 &= \int \left(|\nabla u_n|^2 + \tilde{V}(x)u_n^2 \right) \\ &= \int |\nabla u_n|^2 + e^{-5s_n} \int \tilde{V}(xe^{-s_n})v_n^2 \\ &\leq \int |\nabla u_n|^2 + e^{(|\kappa| - 5)S} \int \tilde{V}(x)v_n^2 \\ &\leq \int |\nabla u_n|^2 + 2e^{(|\kappa| - 5)S} \Big(\int_{\{V > k\}} V(x)v_n^2 + k \int_{\{V \le k\}} v_n^2 \Big) \\ &\leq \int |\nabla u_n|^2 + 2e^{(|\kappa| - 5)S} \Big(\int_{\{V > k\}} V(x)v_n^2 \\ &+ k |\{V \le k\}|^{(\mu - 2)/\mu} \Big(\int |v_n|^{\mu} \Big)^{2/\mu} \Big) \le O(1) \end{split}$$

by (2.20), proving the boundedness of $\{u_n\}$ in X. Using the compact embedding $X \hookrightarrow L^s$ for $s \in [2, 6)$, by standard argument as in [9, Lemma 3.2], we see that $\{u_n\}$ has a convergent subsequence. Hence $\{(s_n, u_n)\}$ has a convergent subsequence. \Box

3. Critical groups and proof of Theorem 1.1

Having established the (PS) condition for $\tilde{\Phi}$, we are now ready to present the proof of Theorem 1.1. We start by recalling some concepts and results from infinitedimensional Morse theory (see e.g., Chang [5] and Mawhin and Willem [21, Chapter 8]).

Let X be a Banach space, $\varphi : X \to \mathbb{R}$ be a C^1 functional, u be an isolated critical point of φ and $\varphi(u) = c$. Then

$$C_i(\varphi, u) := H_i(\varphi_c, \varphi_c \setminus \{0\}), \quad i \in \mathbb{N} = \{0, 1, 2, \ldots\},\$$

is called the *i*-th critical group of φ at u, where $\varphi_c := \varphi^{-1}(-\infty, c]$ and H_* stands for the singular homology with coefficients in \mathbb{Z} .

If φ satisfies the (PS) condition and the critical values of φ are bounded from below by α , then following Bartsch and Li [3], we define the *i*-th critical group of φ at infinity by

$$C_i(\varphi, \infty) := H_i(X, \varphi_\alpha), \quad i \in \mathbb{N}.$$

By the deformation lemma, it is well known that the homology on the right-hand side does not depend on the choice of α .

Proposition 3.1 ([3, Proposition 3.6]). If $\varphi \in C^1(X, \mathbb{R})$ satisfies the (PS) condition and $C_{\ell}(\varphi, 0) \neq C_{\ell}(\varphi, \infty)$ for some $\ell \in \mathbb{N}$, then φ has a nonzero critical point.

Proposition 3.2 ([17, Theorem 2.1]). Suppose $\varphi \in C^1(X, \mathbb{R})$ satisfies the (PS) condition and has a local linking at 0 with respect to the decomposition $X = Y \oplus Z$, *i.e.*, for some $\varepsilon > 0$,

$$\begin{aligned} \varphi(u) &\leq 0 \quad for \ u \in Y \cap B_{\varepsilon}, \\ \varphi(u) &> 0 \quad for \ u \in (Z \setminus \{0\}) \cap B_{\varepsilon} \end{aligned}$$

where $B_{\varepsilon} = \{u \in X : ||u|| \le \varepsilon\}$. If $\ell = \dim Y < \infty$, then $C_{\ell}(\varphi, 0) \ne 0$

Since $\tilde{\Phi}$ satisfies the (PS) condition, the critical group $C_*(\tilde{\Phi}, \infty)$ of $\tilde{\Phi}$ at infinity makes sense. To study $C_*(\tilde{\Phi}, \infty)$ we need the following lemma.

Lemma 3.3. Suppose that (A1)–(A5) hold. Then there exists a constant $M_{\lambda} \ge 0$ such that

$$\frac{\mathrm{d}}{\mathrm{d}t}\tilde{\Phi}(\tau, u_t) \le \frac{\lambda+3}{t} \left(\tilde{\Phi}(\tau, u_t) + M_\lambda\right) \tag{3.1}$$

for all t > 0, $u \in X$ and $\tau \in \mathbb{R}$.

Proof. Let $v = u_{e^{\tau}}$. By (2.4) and Remark 2.2 we observe that

$$\tilde{\Phi}(\tau, u_t) = \frac{\tau^2}{2} + \Phi(u_{te^{\tau}}) = \frac{\tau^2}{2} - \frac{\log^2 t}{2} + \tilde{\Phi}(\log t, v) \ge \tilde{\Phi}(\log t, v) - \frac{\log^2 t}{2} \quad (3.2)$$

$$\frac{\mathrm{d}}{\mathrm{d}t}\tilde{\Phi}(\tau, u_t) = -\frac{\log t}{t} + \frac{1}{t}\partial_s\tilde{\Phi}(\log t, v).$$
(3.3)

We claim that for given $\lambda \in (3, \mu)$, there exists $M_{\lambda} > 0$ such that

$$\partial_s \tilde{\Phi}(s,v) - s \le (\lambda+3) \left(\tilde{\Phi}(s,v) - \frac{s^2}{2} + M_\lambda \right)$$
(3.4)

for all $s \in \mathbb{R}$ and $v \in X$. Indeed, by (2.14), (2.18), and (2.19) we have

$$(\lambda+3)\tilde{\Phi}(s,v) - \partial_s\tilde{\Phi}(s,v) \ge \frac{\lambda+3}{2}s^2 - s + (\mu-\lambda)C\int |v|^{\mu} - C_{\lambda}\left(\int |v|^{\mu}\right)^{2/\mu}.$$

The last two terms are bounded from below because $\mu > \lambda$ and $2/\mu < 1$, thus (3.4) holds. Then, using (3.2), (3.3), and (3.4), with $s = \log t$ and $v = u_{e^{\tau}}$, we deduce that

$$\begin{aligned} \frac{\mathrm{d}}{\mathrm{d}t}\tilde{\Phi}(\tau, u_t) &= -\frac{\log t}{t} + \frac{1}{t}\partial_s\tilde{\Phi}(\log t, v) \\ &\leq \frac{\lambda+3}{t}\Big(\tilde{\Phi}(\log t, u_{e^{\tau}}) - \frac{\log^2 t}{2} + M_\lambda\Big) = \frac{\lambda+3}{t}\left(\Phi(u_{te^{\tau}} + M_\lambda)\right) \\ &\leq \frac{\lambda+3}{t}\Big(\tilde{\Phi}((\tau, u_t) + M_\lambda)\Big). \end{aligned}$$

Lemma 3.4. Suppose (A2), (A4), (A5) hold. Then

$$\lim_{t \to +\infty} \tilde{\Phi}(s, u_t) = -\infty$$

for all $(s, u) \in \mathbb{R} \times X \setminus \{0\}$.

Proof. Because $u_{e^s} \neq 0$, and $(u_t)_{e^s} = (u_{e^s})_t$, it suffices to prove that if $v \neq 0$ then $\Phi(v_t) \to -\infty$ as $t \to +\infty$, see (2.4). From the proof of (2.5) we obtain

$$\Phi(v_t) = \frac{t^3}{2} \int |\nabla v|^2 + \frac{1}{2} \int V(x) t^2 v^2 \left(\frac{x}{t}\right) + \frac{t^6}{4} \left(\int |\nabla v|^2\right)^2 - t^3 \int G(tv)$$

Assumption (A5) implies

$$G(t) \le C|t|^{\mu}$$
 for $|t| < 1$, $G(t) \ge C|t|^{\mu}$ for $|t| \ge 1$. (3.5)

Since $v \neq 0$, we can assume that for some $\varepsilon > 0$, $|\{|v| \ge \varepsilon\}|$ is finite and positive and by (3.5) we have

$$\int G(tv) \ge C \int_{|v| \ge \varepsilon} t^{\mu} |v|^{\mu} \ge C \varepsilon^{\mu} |\{|v| \ge \varepsilon\} |t^{\mu} =: C_{\nu} t^{\mu},$$

for some $C_{\nu} > 0$ and $t \ge 1/\varepsilon$.

As V is bounded on B_R , it follows that

$$\int_{B_R} V(x) t^2 v^2\left(\frac{x}{t}\right) \le \|V\|_{L^{\infty}(B_R)} \int t^2 v^2\left(\frac{x}{t}\right) = t^5 \|V\|_{L^{\infty}(B_R)} \int v^2.$$

Assumption (A2) implies that for $|\omega| = R$ and $r \ge 1$,

$$\frac{\mathrm{d}}{\mathrm{d}r} \left(\frac{\tilde{V}(r\omega)}{r} \right) = \frac{\nabla V(r\omega) \cdot r\omega - \tilde{V}(r\omega)}{r^2} \le 0$$

where $\tilde{V}(x) = V(x) + m$, so that

$$H(x) = \frac{\tilde{V}(x)}{|x|} \chi_{\mathbb{R}^3 \setminus B_R}(x) \text{ is radially non-increasing.}$$

Letting $w(x) = v(x)|x|^{1/2}$, we have

$$\int_{\mathbb{R}^3 \setminus B_R} \tilde{V} t^2 v^2 \left(\frac{x}{t}\right) = |t|^3 \int H(x) w^2 \left(\frac{x}{t}\right) = |t|^6 \int H(xt) w^2,$$

and by the monotonicity of H, $H(xt)w^2 \searrow 0$ as $t \to +\infty$. Therefore, by the monotone convergence theorem,

$$\int V(x)t^2v^2\left(\frac{x}{t}\right) \leq \int_{B_R} V(x)t^2v^2\left(\frac{x}{t}\right) + \int_{\mathbb{R}^3 \setminus B_R} \tilde{V}t^2v^2\left(\frac{x}{t}\right)$$
$$\leq t^5 \|V\|_{L^{\infty}(B_R)} \int v^2 + o(t^6).$$

In summary,

$$\begin{split} \Phi(v_t) &\leq \frac{t^3}{2} \int |\nabla v|^2 + \frac{t^6}{4} \Big(\int |\nabla v|^2 \Big)^2 + \frac{t^5 ||V||_{L^{\infty}(B_R)}}{2} \int v^2 + o(t^6) - C_{\nu} t^{\mu+3} \to -\infty, \\ \text{as } t \to +\infty, \text{ because } \mu+3 > 6. \end{split}$$

Lemma 3.5. Suppose (A1)–(A5) hold. Then for any sufficiently negative $a \in \mathbb{R}$, $C_i(\tilde{\Phi}, \infty) = 0$ for all $i \in \mathbb{N}$.

Proof. Let $\dot{X} = X \setminus \{0\}$ and consider the continuous map $\mathbb{R} \times \dot{X} \times \mathbb{R}^+ \to \mathbb{R} \times \dot{X}$ given by

$$(s, u, t) \mapsto (s, u_t).$$

Choose $\lambda \in (3, \mu)$ and $a < -M_{\lambda}$, where $M_{\lambda} \ge 0$ is given in Lemma 3.3. Then by Lemma 3.4, we have

$$\lim_{t \to +\infty} \tilde{\Phi}(s, u_t) = -\infty, \quad \forall (s, u) \in \mathbb{R} \times \dot{X}.$$

So there is $t_{s,u} > 0$ such that $\tilde{\Phi}(s, u_{t_{s,u}}) = a$. A direct computation and Lemma 3.3 gives

$$\frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=t_{s,u}}\tilde{\Phi}(s,u_t) \leq \frac{\lambda+3}{t_{s,u}}\left(\tilde{\Phi}(s,u_{t_{s,u}}) + M_{\lambda}\right) = \frac{\lambda+3}{t_{s,u}}(a+M_{\lambda}) < 0.$$

By the implicit function theorem, given $(s, u) \in \mathbb{R} \times \dot{X}$, there is a unique solution t = T(s, u) for the equation $\tilde{\Phi}(s, u_t) = a$, and the map

$$T: \{(s,u) \in \mathbb{R} \times \dot{X} : \tilde{\Phi}(s,u) > a\} \to \mathbb{R}^+$$

is continuous. Using the continuous function T, it is standard (see [24, 18]) to construct a deformation from $\mathbb{R} \times \dot{X}$ to the level set $\tilde{\Phi}_a = \tilde{\Phi}^{-1}(-\infty, a]$, and deduce

$$C_i(\tilde{\Phi},\infty) = H_i(\mathbb{R} \times X, \tilde{\Phi}_a) \cong H_i(\mathbb{R} \times X, \mathbb{R} \times \dot{X}) \cong H_i(\mathbb{R} \times X, \{0\} \times S^\infty) = 0,$$

for all $i \in \mathbb{N}$, where $\mathbb{R} \times \dot{X}$ deformation retracts to $\{0\} \times S^{\infty}$.

Recalling that X^+ , X^- and X^0 are the negative, positive and null eigenspaces of the bilinear form defined in (1.9). To study the critical group $C_*(\tilde{\Phi}, 0)$ of $\tilde{\Phi}$ at origin, we consider the decomposition

$$\mathbb{R} \times X = \tilde{X}^- \oplus \tilde{X}^+$$

where $\tilde{X}^- = X^- \oplus X^0, \ \tilde{X}^+ = \mathbb{R} \oplus X^+.$

Lemma 3.6. Assume (A1)-(A5) hold. If either

- (1) $\dim X_{-} > 0$, $\dim X_{0} = 0$,
- (2) dim $X_0 > 0$ and $G(t) \ge c|t|^{\nu}$ for some $\nu < 4$,

then the functional $\tilde{\Phi}$ has a local linking at 0 with respect to the decomposition $\mathbb{R} \times X = \tilde{X}^- \oplus \tilde{X}^+$.

Proof. It suffices to show that for r > 0 small enough,

$$\tilde{\Phi}(s,u) > 0 \quad \text{for } (s,u) \in \tilde{X}^+, \ 0 < ||(s,u)|| < r,$$
(3.6)

$$\Phi(s, u) \le 0 \quad \text{for } (s, u) \in X^{-}, \ \|(s, u)\| < r.$$
(3.7)

We have

$$\int V(xe^s)u^2 = \int V(x)u^2 + \int \left(\int_0^1 \frac{\mathrm{d}}{\mathrm{d}\tau} V(e^{s\tau}x)\mathrm{d}\tau\right)u^2$$
$$= \int V(x)u^2 + s \int \left(\int_0^1 \nabla V(e^{s\tau}x) \cdot (e^{s\tau}x)\mathrm{d}\tau\right)u^2.$$

12

Condition (A3) and (2.2) yield

$$\begin{split} |\int V(e^{s}x)u^{2} - \int V(x)u^{2}| &\leq |s| \int \Big(\int_{0}^{1} |\nabla V(e^{s\tau}x) \cdot (e^{s\tau}x)| \mathrm{d}\tau\Big)u^{2} \\ &\leq \kappa |s| \int \Big(\int_{0}^{1} \tilde{V}(e^{s\tau}x) \mathrm{d}\tau\Big)u^{2} \\ &\leq \kappa |s| \int \Big(\int_{0}^{1} e^{\kappa |s|\tau} \tilde{V}(x) \mathrm{d}\tau\Big)u^{2} \\ &\leq (e^{\kappa |s|} - 1) \int \tilde{V}(x)u^{2} \\ &\leq O(s) \|u\|^{2} = o(\|(s,u)\|^{2}). \end{split}$$
(3.8)

Moreover, (A4) and (A5) imply that

$$|G(t)| \le C(|t|^{\mu} + |t|^{p+1}).$$

Consequently, for $||(s, u)|| \to 0$, we have

$$e^{3s} \left| \int G(e^{s}u) \right| \le Ce^{3s} \int \left(|e^{s}u|^{\mu} + |e^{s}u|^{p+1} \right) \le o(||u||^{2}), \tag{3.9}$$

while by (1.7) we easily have $|V(x)| \leq 3\tilde{V}(x)$, hence

$$|s| \int |V(x)| u^2 \le 3|s| \int \tilde{V}(x) u^2 = o(||(s,u)||^2).$$
(3.10)

Combining (3.8)-(3.10), we obtain

$$\begin{split} \tilde{\Phi}(s,u) &= \frac{s^2}{2} + \frac{e^{3s}}{2} \int |\nabla u|^2 + \frac{e^{6s}}{4} \left(\int |\nabla u|^2 \right)^2 + \frac{e^{5s}}{2} \int V(xe^s) u^2 - e^{3s} \int G(e^s u) \\ &= \frac{s^2}{2} + \frac{1 + O(s)}{2} \int |\nabla u|^2 + \frac{1 + O(s)}{2} \int V u^2 \\ &+ O(||u||^4) + o(||(s,u)||^2) + o(||u||^2) \\ &= \frac{1}{2} \int \left(|\nabla u|^2 + V(x) u^2 \right) + \frac{s^2}{2} + O(||u||^4) + o(||(s,u)||^2). \end{split}$$

Then, it is easy to see that (3.6) is true.

Now we prove (3.7) for r > 0 small. In item (1) we have s = 0, it is easy to see that

$$\Phi(0, u) < 0 \text{ for } u \in X^{-}, ||u|| < r.$$

In item (2), by simple computations,

$$\begin{split} \tilde{\Phi}(0,u) &= \frac{1}{2} \int \left(|\nabla u|^2 + V(x)u^2 \right) + \frac{1}{4} \left(\int |\nabla u|^2 \right)^2 - \int G(u) \\ &\leq \frac{1}{2} \int \left(|\nabla u| + V(x)u^2 \right) + O(||u||^4) - c \int |u|^\nu. \end{split}$$

Since all norms on $X^- \oplus X^0$ are equivalent we deduce that

$$\tilde{\Phi}(0,u) \le -\eta^{-} ||u^{-}||^{2} + O(||u||^{4}) - c||u||^{\nu} \quad \text{for } u \in X^{-} \oplus X^{0}.$$

Because $\nu < 4$, we can choose r > 0 small enough such that

$$\Phi(0, u) < 0 \text{ for } u \in X^- \oplus X^0, \ ||u|| < r.$$

Hence, (3.7) has also been verified item (2).

Proof of Theorem 1.1. We have shown that $\tilde{\Phi}$ satisfies (PS) condition and has a local linking at 0 with respect to the decomposition $\mathbb{R} \times X = \tilde{X}^- \oplus \tilde{X}^+$. By Proposition 3.2 we have

$$C_i(\Phi, 0) \neq 0$$

where $i = \dim \tilde{X}^-$. By Lemma 3.5, $C_i(\tilde{\Phi}, \infty) = 0$, therefore

$$C_i(\Phi, 0) \neq C_i(\Phi, \infty).$$

Applying Proposition 3.1, we know that Φ has a nonzero critical point (\bar{s}, \bar{u}) . Furthermore, Lemma 2.5 implies that $\bar{s} = 0$, $\bar{u} \neq 0$ and $D\Phi(\bar{u}) = 0$. Thus Φ has a nonzero critical point \bar{u} , which is a nontrivial solution of the problem (1.6).

Acknowledgments. This research was partially supported by NSFC (12071387).

References

- C. O. Alves, G. M. Figueiredo; Nonlinear perturbations of a periodic Kirchhoff equation in
 ^N, Nonlinear Anal., 75 (2012), 2750–2759.
- [2] A. Ambrosetti, P. H. Rabinowitz; Dual variational methods in critical point theory and applications, J. Functional Analysis, 14 (1973), 349–381.
- [3] T. Bartsch, S. Li; Critical point theory for asymptotically quadratic functionals and applications to problems with resonance, *Nonlinear Anal.*, 28 (1997), 419–441.
- [4] T. Bartsch, Z. Q. Wang; Existence and multiplicity results for some superlinear elliptic problems on R^N, Comm. Partial Differential Equations, 20 (1995), 1725–1741.
- K.-c. Chang; Infinite-dimensional Morse theory and multiple solution problems, Progress in Nonlinear Differential Equations and their Applications, 6, Birkhäuser Boston, Inc., Boston, MA, 1993.
- [6] S. Chen, S. Liu; Standing waves for 4-superlinear Schrödinger-Kirchhoff equations, Math. Methods Appl. Sci., 38 (2015), 2185–2193.
- [7] X.-M. He, W.-M. Zou; Multiplicity of solutions for a class of Kirchhoff type problems, Acta Math. Appl. Sin. Engl. Ser., 26 (2010), 387–394.
- [8] L. Jeanjean; Existence of solutions with prescribed norm for semilinear elliptic equations, Nonlinear Anal., 28 (1997), 1633–1659.
- [9] S. Jiang, S. Liu; Multiple solutions for Schrödinger-Kirchhoff equations with indefinite potential, Appl. Math. Lett., 124 (2022), Paper No. 107672, 9.
- [10] G. Kirchhoff; Mechanik, Teubner, Leipzig, 1883.
- [11] W. Kryszewski, A. Szulkin; Generalized linking theorem with an application to a semilinear Schrödinger equation, Adv. Differential Equations, 3 (1998), 441–472.
- [12] G. Li, H. Ye; Existence of positive ground state solutions for the nonlinear Kirchhoff type equations in \mathbb{R}^3 , J. Differential Equations, **257** (2014), 566–600.
- [13] L. Li, J. Xu; Kirchhoff equations with indefinite potentials, Appl. Anal., 101 (2022), 6081– 6089.
- [14] Y. Li, F. Li, J. Shi; Existence of a positive solution to Kirchhoff type problems without compactness conditions, J. Differential Equations, 253 (2012), 2285–2294.
- [15] J.-L. Lions; On some questions in boundary value problems of mathematical physics, in Contemporary developments in continuum mechanics and partial differential equations (Proc. Internat. Sympos., Inst. Mat., Univ. Fed. Rio de Janeiro, Rio de Janeiro, 1977), North-Holland Math. Stud., vol. 30, North-Holland, Amsterdam-New York, 1978, 284–346.
- [16] J. Liu, J. Su; Remarks on multiple nontrivial solutions for quasi-linear resonant problems, J. Math. Anal. Appl., 258 (2001), 209–222.
- [17] J. Q. Liu; The Morse index of a saddle point, Systems Sci. Math. Sci., 2 (1989), 32–39.
- [18] S. Liu; Existence of solutions to a superlinear p-Laplacian equation, Electron. J. Differential Equations, 2001 (2001), No. 66, 1–6.
- [19] S. Liu; On superlinear Schrödinger equations with periodic potential, Calc. Var. Partial Differential Equations, 45 (2012), 1-9.
- [20] S. Liu, S. Mosconi; On the Schrödinger-Poisson system with indefinite potential and 3sublinear nonlinearity, J. Differential Equations, 269 (2020), 689–712.

- [21] J. Mawhin, M. Willem; Critical point theory and Hamiltonian systems, Applied Mathematical Sciences, vol. 74, Springer-Verlag, New York, 1989.
- [22] K. Perera, Z. Zhang; Nontrivial solutions of Kirchhoff-type problems via the Yang index, J. Differential Equations, 221 (2006), 246-255.
- [23] J. Sun, S. Liu; Nontrivial solutions of Kirchhoff type problems, Appl. Math. Lett., 25 (2012), 500–504.
- [24] Z. Q. Wang; On a superlinear elliptic equation, Ann. Inst. H. Poincaré Anal. Non Linéaire, 8 (1991), 43–57.
- [25] M. Willem; *Minimax theorems*, Progress in Nonlinear Differential Equations and their Applications, 24, Birkhäuser Boston Inc., Boston, MA, 1996.
- [26] X. Wu; Existence of nontrivial solutions and high energy solutions for Schrödinger-Kirchhofftype equations in \mathbb{R}^N , Nonlinear Anal. Real World Appl., **12** (2011), 1278–1287.
- [27] Y. Wu, S. Liu; Existence and multiplicity of solutions for asymptotically linear Schrödinger-Kirchhoff equations, Nonlinear Anal. Real World Appl., 26 (2015), 191–198.
- [28] Y. Zhang, X. Tang, D. Qin; Infinitely many solutions for Kirchhoff problems with lack of compactness, *Nonlinear Anal.*, **197** (2020), 111856, 31.

Shuai Jiang

SCHOOL OF MATHEMATICAL SCIENCES, XIAMEN UNIVERSITY, XIAMEN 361005, CHINA Email address: jiangshuai0915@163.com

LI-FENG YIN

SCHOOL OF MATHEMATICAL SCIENCES, XIAMEN UNIVERSITY, XIAMEN 361005, CHINA Email address: yin136823@163.com