Hindawi Publishing Corporation Abstract and Applied Analysis Volume 2011, Article ID 901631, 14 pages doi:10.1155/2011/901631

Research Article

Oscillation Properties for Second-Order Partial Differential Equations with Damping and Functional Arguments

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Received 19 September 2011; Accepted 21 October 2011

Academic Editor: Norio Yoshida

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Using an integral averaging method and generalized Riccati technique, by introducing a parameter $\beta \ge 1$, we derive new oscillation criteria for second-order partial differential equations with damping. The results are of high degree of generality and sharper than most known ones.

1. Introduction

Consider the second-order partial delay differential equation

$$\frac{\partial}{\partial t} \left(r(t) \frac{\partial}{\partial t} u(x, t) \right) + p(t) \frac{\partial u(x, t)}{\partial t} = a(t) \Delta u(x, t) + \sum_{k=1}^{s} a_k(t) \Delta u(x, t - \rho_k(t)) - q(x, t) f(u(x, t))$$

$$- \sum_{j=1}^{m} q_j(x, t) f_j(u(x, t - \sigma_j)), \quad (x, t) \in \Omega \times R_+ \equiv G,$$
(1.1)

where Δ is the Laplacian in \mathbb{R}^N , $\mathbb{R}_+ = [0, \infty)$ and Ω is a bounded domain in \mathbb{R}^N with a piecewise smooth boundary $\partial\Omega$.

Throughout this paper, we assume that

$$(H_1)$$
 $r(t) \in C^1(R_+, (0, \infty)), p(t) \in C(R_+, R);$

$$(H_2) \ q(x,t), q_j(x,t) \in C(\overline{G}, R_+), \ q(t) = \min_{x \in \overline{G}} q(x,t), \ q_j(t) = \min_{x \in \overline{G}} q_j(x,t), \ j \in I_m = \{1, 2, \ldots, m\};$$

- (H₃) a(t), $a_k(t)$, $\rho_k(t) \in C(R_+, R_+)$, $\lim_{t\to\infty} (t \rho_k(t)) = \infty$, and σ_j are nonnegative constants, $j \in I_m$, $k \in I_s = \{1, 2, ..., s\}$;
- (H_4) $f(u) \in C^1(R, R)$, $f_j(u) \in C(R, R)$ are convex in R_+ with $uf_j(u) > 0$, uf(u) > 0, and $f'(u) \ge u > 0$, $(u \ne 0)$.

We say that a continuous function H(t,s) belongs to the function class ω , denoted by $H \in \omega$, if $H \in C(D,R_+)$, where $D = \{(t,s) : -\infty < s \le t < +\infty\}$, satisfy

$$H(t,t) = 0, \quad H(t,s) > 0, \quad -\infty < s < t < +\infty.$$
 (1.2)

Furthermore, the continuous partial derivative $\partial H/\partial S$ exists on D, and there is $h \in L_{loc}(D, R)$, such that

$$\frac{\partial H}{\partial s} = -h(t, s)\sqrt{H(t, s)}. \tag{1.3}$$

Various results on the oscillation for the partial functional differential equation have been obtained recently. We refer the reader to [1–3] for parabolic equations and to [4–11] for hyperbolic equations.

Recently, Li and Cui [12] studied the equation of the form

$$\frac{\partial}{\partial t} \left[p(t) \frac{\partial}{\partial t} \left(u(x,t) + \sum_{i=1}^{l} \lambda_i(t) u(x,t - \tau_i) \right) \right] = a(t) \Delta u(x,t) + \sum_{k=1}^{s} a_k(t) \Delta u(x,t - \rho_k(t)) - q(t) u(x,t) \\
- \sum_{j=1}^{m} q_j(x,t) u(x,t - \sigma_j), \quad (x,t) \in \Omega \times R_+ \equiv G \tag{1.4}$$

with Robin boundary condition

$$\frac{\partial u(x,t)}{\partial y} + g(x,t)u(x,t) = 0, \quad (x,t) \in \partial \Omega \times R_+, \tag{1.5}$$

where γ is the unit exterior vector to $\partial\Omega$ and g(x,t) is a nonnegative continuous function on $\partial\Omega \times R_+$ and obtained the following result.

Theorem A (see [12, Theorem 2.2]). *Suppose that* $H \in \omega$, *let*

- (C₁) $0 < \inf_{s \ge t_0} \{ \liminf_{t \to \infty} (H(t,s)/H(t,t_0)) \} \le \infty$, suppose that there exists some $j_0 \in I_m$ and there exist two functions $\phi \in C^1[t_0,\infty)$, $A \in C[t_0,\infty)$ satisfying,
- (C₂) $\limsup_{t\to\infty} (1/H(t,t_0)) \int_{t_0}^t p(s-\sigma_{j_0})\phi(s)h^2(t,s)ds < \infty$,
- (C_3) $\int_{t_0}^{\infty} (A_+^2(s)/p(s-\sigma_{j_0})\phi(s))ds = \infty$, and for every $t_1 \ge t_0$,
- $(C_4) \limsup_{t\to\infty} (1/H(t,t_1)) \int_{t_1}^t [H(t,s)\psi(s) (1/4)\phi(s)p(s-\sigma_{i_0})h^2(t,s)]ds \ge A(t_1),$

where $\phi(s) = \exp\{-2\int^s \phi(\xi)d\xi\}$, $A_+(s) = \max\{A(s), 0\}$, and $\psi(s) = \phi(s)\{\alpha_{j_0}q_{j_0}(s) [1-\sum_{i=1}^l \lambda_i(s-\sigma_{j_0})] + p(s-\sigma_{j_0})\phi^2(s) - [p(s-\sigma_{j_0})\phi(s)]'\}$. Then every solution u(x,t) of the problem (1.4), (1.5) is oscillatory in G.

In 2008, Rogovchenko and Tuncay [13] established new oscillation criteria for secondorder nonlinear differential equations with damping term

$$(r(t)x'(t))' + p(t)x'(t) + q(t)f(x(t)) = 0,$$
(1.6)

without an assumption that has been required in related results reported in the literature over the last two decades. Motivated by the ideas in [12, 13], by introducing a Parameter $\beta \geq 1$, we will further improve Theorems A and derive new interval criteria for oscillation of (1.1). We suggest two different approaches which allow one to remove condition (C_2) in Theorem A. A modified integral averaging technique enables one to simplify essentially the proofs of oscillation criteria.

2. Main Results

Theorem 2.1. Suppose that there exists a function $y \in C^1[t_0, \infty)$ such that for some $\beta \ge 1$ and for some $H \in \omega$,

$$\lim_{t\to\infty}\sup\frac{1}{H(t,t_0)}\int_{t_0}^t \left(H(t,s)\psi(s)-\frac{\beta}{4\mu}v(s)r(s)h^2(t,s)\right)ds=\infty,\quad t\geq 0, \tag{2.1}$$

where

$$v(t) = \exp\left(-2\int_{-\infty}^{t} \left(\mu y(s) - \frac{p(s)}{2r(s)}\right) ds\right), \tag{2.2}$$

$$\psi(t) = v(t) \left[q(t) + \mu r(t) y^{2}(t) - p(t) y(t) - (r(t) y(t))' \right], \tag{2.3}$$

then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Proof. Suppose to the contrary that there is a nonoscillatory solution u(x,t) of the problem (1.1), (1.5) which has no zero on $\Omega \times [t_0,\infty)$ for some $t_0>0$. Without loss of generality, we assume that u(x,t)>0, $u(x,t-\rho_k(t))>0$ and $u(x,t-\sigma_j)>0$ in $\Omega \times [t_1,\infty)$, $t_1\geq t_0$, $k\in I_s$, $j\in I_m$. Integrating (1.1) with respect to x over the domain Ω , we have

$$\frac{d}{dt}\left(r(t)\frac{d}{dt}\int_{\Omega}u(x,t)dx\right) + p(t)\frac{d}{dt}\int_{\Omega}u(x,t)dx$$

$$= a(t)\int_{\Omega}\Delta u(x,t)dx + \sum_{k=1}^{s}a_{k}(t)\int_{\Omega}\Delta u(x,t-\rho_{k}(t))dx$$

$$-\int_{\Omega}q(x,t)f(u(x,t))dx - \sum_{j=1}^{m}\int_{\Omega}q_{j}(x,t)f_{j}(u(x,t-\sigma_{j}))dx, \quad t \ge t_{1}.$$
(2.4)

From Green's formula and the boundary condition (1.5), we have

$$\int_{\Omega} \Delta u(x,t) dx = \int_{\partial\Omega} \frac{\partial u(x,t)}{\partial \gamma} ds = -\int_{\partial\Omega} g(x,t) u(x,t) ds \le 0,$$

$$\int_{\Omega} \Delta u(x,t-\rho_{k}(t)) dx = \int_{\partial\Omega} \frac{\partial u(x,t-\rho_{k}(t))}{\partial \gamma} ds$$

$$= -\int_{\partial\Omega} g(x,t-\rho_{k}(t)) u(x,t-\rho_{k}(t)) ds \le 0, \quad t \ge t_{1}, \ k \in I_{s},$$
(2.5)

where ds denotes the surface element on $\partial\Omega$. Moreover, from (H_2) , (H_4) and Jensen's inequality, we have

$$\int_{\Omega} q(x,t)f(u(x,t))dx \ge q(t) \int_{\Omega} f(u(x,t))dx \ge |\Omega|q(t)f\left(\frac{1}{|\Omega|} \int_{\Omega} u(x,t)dx\right),$$

$$\int_{\Omega} q_{j}(x,t)f_{j}\left(u(x,t-\sigma_{j})\right)dx \ge q_{j}(t) \int_{\Omega} f_{j}\left(u(x,t-\sigma_{j})\right)dx$$

$$\ge |\Omega|q_{j}(t)f_{j}\left(\frac{1}{|\Omega|} \int_{\Omega} u(x,t-\sigma_{j})\right)dx,$$
(2.6)

where $|\Omega| = \int_{\Omega} dx$. Set

$$U(t) = \frac{1}{|\Omega|} \int_{\Omega} u(x, t) dx, \quad t \ge t_1.$$
 (2.7)

In view of (2.5)–(2.7), (2.4) yields that

$$(r(t)U'(t))' + p(t)U'(t) + q(t)f(U(t)) + \sum_{j=1}^{m} q_j(t)f_j(U(t-\sigma_j)) \le 0, \quad t \ge t_1.$$
 (2.8)

Note that (H_4) , (2.8) yields that

$$(r(t)U'(t))' + p(t)U'(t) + q(t)f(U(t)) \le 0, \quad t \ge t_1.$$
(2.9)

Put

$$w(t) = v(t)r(t) \left[\frac{U'(t)}{f(U(t))} + y(t) \right], \quad t \ge t_1,$$
 (2.10)

where v(t) is given by (2.2), then

$$\begin{split} w'(t) &= \left[-2\mu y(t) + \frac{p(t)}{r(t)} \right] w(t) + v(t) \left[\frac{(r(t)U'(t))'}{f(U(t))} - \frac{r(t)f'(U(t))(U'(t))^2}{f^2(U(t))} + (r(t)y(t))' \right] \\ &\leq \left[-2\mu y(t) + \frac{p(t)}{r(t)} \right] w(t) + v(t) \left[\frac{(r(t)U'(t))'}{f(U(t))} - \frac{\mu r(t)(U'(t))^2}{f^2(U(t))} + (r(t)y(t))' \right] \\ &\leq \left[-2\mu y(t) + \frac{p(t)}{r(t)} \right] w(t) - v(t) \left[\frac{p(t)U'(t)}{f(U(t))} + q(t) + \mu r(t) \left(\frac{U'(t)}{f(U(t))} \right)^2 - (r(t)y(t))' \right] \\ &= \left[-2\mu y(t) + \frac{p(t)}{r(t)} \right] w(t) - v(t) \left[p(t) \left(\frac{w(t)}{v(t)r(t)} - y(t) \right) + q(t) \right. \\ &\qquad \qquad + \mu r(t) \left[\frac{w(t)}{v(t)r(t)} - y(t) \right]^2 - (r(t)y(t))' \right] \\ &= \left[-2\mu y(t) + \frac{p(t)}{r(t)} \right] w(t) \\ &- v(t) \left[\frac{p(t)}{v(t)r(t)} w(t) - p(t)y(t) + q(t) \right. \\ &\qquad \qquad + \mu r(t) \frac{w^2(t)}{v^2(t)r^2(t)} - 2\mu \frac{w(t)y(t)}{v(t)} + \mu r(t)y^2(t) - (r(t)y(t))' \right] \\ &= -2\mu y(t)w(t) + \frac{p(t)}{r(t)} w(t) - \frac{p(t)}{r(t)} w(t) - v(t) \left[-p(t)y(t) + q(t) + \mu r(t)y^2(t) - (r(t)y(t))' \right] \\ &+ 2\mu y(t)w(t) - \mu \frac{w^2(t)}{v(t)r(t)} \\ &= -v(t) \left[-p(t)y(t) + q(t) + \mu r(t)y^2(t) - (r(t)y(t))' \right] - \mu \frac{w^2(t)}{v(t)r(t)} \\ &= -\psi(t) - \mu \frac{w^2(t)}{v(t)r(t)}, \end{split}$$

that is,

$$\psi(t) \le -w'(t) - \mu \frac{w^2(t)}{v(t)r(t)},$$
(2.12)

where $\psi(t)$ is defined by (2.3). Multiplying (2.12) by H(t,s) and integrating from T to t, we have, for some $\beta \ge 1$ and for all $t \ge T \ge t_1$,

$$\int_{T}^{t} H(t,s)\psi(s)ds \leq -\int_{T}^{t} H(t,s)w'(s)ds - \int_{T}^{t} H(t,s)\frac{\mu}{v(s)r(s)}w^{2}(s)ds
= H(t,T)w(T) - \int_{T}^{t} \left(h(t,s)\sqrt{H(t,s)}w(s) + H(t,s)\frac{\mu w^{2}(s)}{v(s)r(s)}\right)ds
= H(t,T)w(T) - \int_{T}^{t} \left[\sqrt{\frac{\mu H(t,s)}{\beta v(s)r(s)}}w(s) + \sqrt{\frac{\beta v(s)r(s)}{4\mu}}h(t,s)\right]^{2}ds
+ \frac{\beta}{4\mu}\int_{T}^{t} v(s)r(s)h^{2}(t,s)ds - \int_{T}^{t} \frac{(\beta-1)\mu}{\beta v(s)r(s)}H(t,s)w^{2}(s)ds.$$
(2.13)

Writing the latter inequality in the form

$$\int_{T}^{t} \left[H(t,s)\psi(s) - \frac{\beta}{4\mu}v(s)r(s)h^{2}(t,s) \right] ds$$

$$\leq H(t,T)w(T) - \int_{T}^{t} \left[\sqrt{\frac{\mu H(t,s)}{\beta v(s)r(s)}}w(s) + \sqrt{\frac{\beta v(s)r(s)}{4\mu}}h(t,s) \right]^{2} ds \qquad (2.14)$$

$$- \int_{T}^{t} \frac{(\beta-1)\mu}{\beta v(s)r(s)} H(t,s)w^{2}(s)ds.$$

Using the properties of H(t, s), we have

$$\int_{t_1}^{t} \left[H(t,s)\psi(s) - \frac{\beta}{4\mu} v(s)r(s)h^2(t,s) \right] ds \le H(t,t_1)|w(t_1)| \le H(t,t_0)|w(t_1)|, \quad t \ge t_1, \quad (2.15)$$

and for all $t \ge t_1 \ge t_0$,

$$\int_{t_0}^{t} \left[H(t,s)\psi(s) - \frac{\beta}{4\mu} v(s)r(s)h^2(t,s) \right] ds \le H(t,t_0) \left[\int_{t_0}^{t_1} |\psi(s)| ds + |w(t_1)| \right]. \tag{2.16}$$

By (2.16),

$$\lim_{t \to \infty} \sup \frac{1}{H(t, t_0)} \int_{t_0}^t \left(H(t, s) \psi(s) - \frac{\beta}{4\mu} v(s) r(s) h^2(t, s) \right) ds \le \int_{t_0}^{t_1} |\psi(s)| ds + |w(t_1)| < \infty, \tag{2.17}$$

which contradicts (2.1). This proves Theorem 2.1.

Consider a Kamenev-type function H(t,s) defined by $H(t,s) = (t-s)^{n-1}$, $(t,s) \in D$, where n > 2 is an integer. Obviously, H belongs to the class ω , and $h(t,s) = (n-1)(t-s)^{(n-3)/2}$, $(t,s) \in D$. Then, we can get the following results.

Corollary 2.2. Suppose that there exists a function $y(t) \in C^1([t_0, \infty); R)$ such that for some integer n > 2 and some $\beta \ge 1$,

$$\lim_{t \to \infty} \sup \frac{1}{t^{n-1}} \int_{t_0}^t (t-s)^{n-3} \left[(t-s)^2 \psi(s) - \frac{\beta(n-1)^2}{4\mu} v(s) r(s) \right] ds = \infty, \tag{2.18}$$

where v(t) and $\psi(t)$ are as defined in Theorem 2.1. Then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Theorem 2.3. Suppose that

$$0 < \inf_{s \ge t_0} \left(\lim_{t \to \infty} \inf \frac{H(t, s)}{H(t, t_0)} \right) \le \infty.$$
 (2.19)

Assume that there exist functions $f \in C^1([t_0, \infty); R)$ and $\phi \in C([t_0, \infty); R)$ such that, for all $t \ge T \ge t_0$ and for some $\beta > 1$,

$$\lim_{t \to \infty} \sup \frac{1}{H(t,T)} \int_{T}^{t} \left(H(t,s)\psi(s) - \frac{\beta}{4\mu} v(s)r(s)h^{2}(t,s) \right) ds \ge \phi(T), \tag{2.20}$$

where v(t), $\psi(t)$ are as defined in Theorem 2.1 and suppose further that

$$\lim_{t \to \infty} \sup \int_{t_0}^t \frac{\phi_+^2(s)}{v(s)r(s)} = \infty, \tag{2.21}$$

where $\phi_+(t) = \max(\phi(t), 0)$. Then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Proof. Suppose to the contrary that there is a nonoscillatory solution u(x,t) of the problem (1.1), (1.5) which has no zero on $\Omega \times [t_1, \infty)$ for some $t_1 > t_0$, without loss of generality, we assume that u(x,t) > 0, $u(x,t-\rho_k(t)) > 0$ and $u(x,t-\sigma_j) > 0$ in $\Omega \times [t_1,\infty)$, $t \ge t_1 \ge t_0$, $k \in I_s$, $j \in I_m$.

As in the proof of Theorem 2.1, (2.14) holds for all $t \ge T \ge t_1$, we have

$$\frac{1}{H(t,T)} \int_{T}^{t} \left[H(t,s)\psi(s) - \frac{\beta}{4\mu} v(s)r(s)h^{2}(t,s) \right] ds$$

$$\leq w(T) - \frac{1}{H(t,T)} \int_{T}^{t} \left[\sqrt{\frac{\mu H(t,s)}{\beta v(s)r(s)}} w(s) + \sqrt{\frac{\beta v(s)r(s)}{4\mu}} h(t,s) \right]^{2} ds$$

$$- \frac{1}{H(t,T)} \int_{T}^{t} \frac{(\beta - 1)\mu}{\beta v(s)r(s)} H(t,s)w^{2}(s) ds$$

$$\leq w(T) - \frac{1}{H(t,T)} \int_{T}^{t} \frac{(\beta - 1)\mu}{\beta v(s)r(s)} H(t,s)w^{2}(s) ds.$$
(2.22)

Therefore, for $t > T \ge T_0$,

$$\lim_{t \to \infty} \sup \frac{1}{H(t,T)} \int_{T}^{t} \left[H(t,s)\psi(s) - \frac{\beta}{4\mu} v(s)r(s)h^{2}(t,s) \right] ds$$

$$\leq w(T) - \lim_{t \to \infty} \inf \frac{1}{H(t,T)} \int_{T}^{t} \frac{(\beta - 1)\mu}{\beta v(s)r(s)} H(t,s)w^{2}(s) ds.$$
(2.23)

It follows from (2.20) that

$$w(T) \ge \phi(T) + \lim_{t \to \infty} \inf \frac{1}{H(t, T)} \int_{T}^{t} \frac{(\beta - 1)\mu}{\beta v(s)r(s)} H(t, s) w^{2}(s) ds$$
 (2.24)

for all $T \ge t_1$ and for any $\beta > 1$. Then, for all $T \ge t_1$,

$$w(T) \ge \phi(T),\tag{2.25}$$

$$\lim_{t \to \infty} \inf \frac{1}{H(t, t_1)} \int_{t_1}^t \frac{H(t, s)}{v(s)r(s)} w^2(s) ds \le \frac{\beta}{\mu(\beta - 1)} (w(t_1) - \phi(t_1)). \tag{2.26}$$

Now, we claim that

$$\int_{t_1}^{\infty} \frac{w^2(s)}{v(s)r(s)} ds < \infty. \tag{2.27}$$

Suppose the contrary, that is,

$$\int_{t_1}^{\infty} \frac{w^2(s)}{v(s)r(s)} ds = \infty. \tag{2.28}$$

By (2.19), there is a positive constant M_1 , satisfying

$$\inf_{s \ge t_0} \left(\lim_{t \to \infty} \inf \frac{H(t, s)}{H(t, t_0)} \right) > M_1 > 0.$$
 (2.29)

Let M be any arbitrary positive number, then from (2.28) we get that there exists a $T_1 > t_1$ such that, for all $t \ge T_1$,

$$\int_{t_1}^{t} \frac{w^2(s)}{v(s)r(s)} ds \ge \frac{M}{M_1}.$$
(2.30)

Using integration by parts, for all $t \ge T_1$, we get

$$\frac{1}{H(t,t_1)} \int_{t_1}^{t} H(t,s) \frac{w^2(s)}{v(s)r(s)} ds = \frac{1}{H(t,t_1)} \int_{t_1}^{t} \left(-\frac{\partial H(t,s)}{\partial s} \right) \left(\int_{t_1}^{s} \frac{w^2(\tau)}{v(\tau)r(\tau)} d\tau \right) ds$$

$$\geq \frac{1}{H(t,t_1)} \int_{T_1}^{t} \left(-\frac{\partial H(t,s)}{\partial s} \right) \left(\int_{t_1}^{s} \frac{w^2(\tau)}{v(\tau)r(\tau)} d\tau \right) ds$$

$$\geq \frac{M}{M_1} \frac{1}{H(t,t_1)} \int_{T_1}^{t} \left(-\frac{\partial H(t,s)}{\partial s} \right) ds$$

$$= \frac{M}{M_1} \frac{H(t,T_1)}{H(t,t_1)}.$$
(2.31)

By (2.29), there exists a $T_2 > T_1$ such that, for all $t \ge T_2$,

$$\frac{H(t, T_1)}{H(t, t_1)} \ge M_1. \tag{2.32}$$

It follows from (2.31) that for all $t \ge T_2$,

$$\frac{1}{H(t,t_1)} \int_{t_1}^t H(t,s) \frac{w^2(s)}{v(s)r(s)} ds \ge M. \tag{2.33}$$

Since M is an arbitrary positive constant,

$$\lim_{t \to \infty} \inf \frac{1}{H(t, t_1)} \int_{t_1}^t H(t, s) \frac{w^2(s)}{v(s)r(s)} ds = \infty, \tag{2.34}$$

which contradicts (2.26). Consequently, (2.27) holds. And from (2.25), we obtain

$$\int_{t_1}^{\infty} \frac{\phi_+^2(s)}{v(s)r(s)} ds \le \int_{t_1}^{\infty} \frac{w^2(s)}{v(s)r(s)} ds < \infty, \tag{2.35}$$

which contradicts (2.21). This completes the proof of Theorem 2.3.

Choosing *H* as in Corollary 2.2, by Theorem 2.3, we can obtain the following corollary.

Corollary 2.4. Let v(t) and $\psi(t)$ be as in Theorem 2.1, assume further that there exist functions $f \in C^1([t_0,\infty);R)$ and $\phi \in C([t_0,\infty);R)$ such that, for all $T \ge t_0$, for some integer n > 2, and for some $\beta > 1$,

$$\lim_{t \to \infty} \sup \frac{1}{t^{n-1}} \int_{T}^{t} (t-s)^{n-3} \left[(t-s)^{2} \psi(s) - \frac{\beta(n-1)^{2}}{4\mu} v(s) r(s) \right] ds \ge \phi(T)$$
 (2.36)

and (2.21) hold. Then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Theorem 2.5. Suppose that there exists a function $f \in C^1([t_0, \infty); R)$ such that for some $\beta \ge 1$ and for some $H \in \omega$,

$$\lim_{t \to \infty} \sup \frac{1}{H(t, t_0)} \int_{t_0}^t \left[H(t, s) \left(\overline{\psi}(s) - \frac{p^2(s) \overline{v}(s)}{2\mu r(s)} \right) - \frac{\mu \beta}{2} \overline{v}(s) r(s) h^2(t, s) \right] ds = \infty, \quad (2.37)$$

where

$$\overline{v}(t) = \exp\left(-2\mu \int_{-\infty}^{t} y(s)ds\right),$$

$$\overline{\psi}(t) = \overline{v}(t)\left(q(t) + \mu r(t)y^{2}(t) - p(t)y(t) - (r(t)y(t))'\right).$$
(2.38)

Then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Proof. As in Theorem 2.1, without loss of generality, we assume that a nonoscillatory solution u(x,t) of the problem (1.1), (1.5) satisfies u(x,t) > 0, $u(x,t-\rho_k(t)) > 0$ and $u(x,t-\sigma_j) > 0$ in $\Omega \times [t_1,\infty)$, $t_1 \ge t_0$, $k \in I_s$, $j \in I_m$. Define a generalized Riccati transformation

$$\overline{w}(t) = \overline{v}(t)r(t)\left[\frac{U'(t)}{f(U(t))} + y(t)\right], \quad t \ge t_1, \tag{2.39}$$

where $\overline{v}(t)$ is given by (2.38). Then

$$\overline{w}'(t) \leq -2\mu y(t)\overline{w}(t) + \overline{v}(t) \left\{ -q(t) + (r(t)y(t))' - p(t) \left[\frac{\overline{w}(t)}{\overline{v}(t)r(t)} - y(t) \right] - \mu r(t) \left[\frac{\overline{w}(t)}{\overline{v}(t)r(t)} - y(t) \right]^{2} \right\}$$

$$= -\overline{\psi}(t) - \frac{p(t)}{r(t)}\overline{w}(t) - \mu \frac{1}{r(t)\overline{v}(t)}\overline{w}^{2}(t), \quad t \geq t_{1}.$$

$$(2.40)$$

Using an elementary inequality

$$-ax^2 + bx \le -\frac{a}{2}x^2 + \frac{b^2}{2a'} \tag{2.41}$$

for all a > 0 and for all $b, x \in R$, we conclude from (2.40) that

$$\overline{\psi}(t) - \frac{p^2(t)\overline{v}(t)}{2\mu r(t)} \le -\overline{w}'(t) - \frac{\mu}{2\overline{v}(t)r(t)}\overline{w}^2(t), \quad t \ge t_1. \tag{2.42}$$

Multiplying (2.42) by H(t, s) and integrating from T < t, we obtain, for some $\beta \ge 1$ and for all $t \ge T \ge t_1$,

$$\int_{T}^{t} H(t,s) \left(\overline{\psi}(s) - \frac{p^{2}(s)\overline{v}(s)}{2\mu r(s)} \right) ds \leq H(t,T)\overline{w}(T) - \int_{T}^{t} h(t,s)\sqrt{H(t,s)}\overline{w}(s)ds - \int_{T}^{t} \frac{\mu\overline{w}^{2}(s)}{2\overline{v}(s)r(s)} ds$$

$$\leq H(t,T)\overline{w}(T) + \frac{\mu\beta}{2} \int_{T}^{t} \overline{v}(s)r(s)h^{2}(t,s)ds$$

$$- \mu \int_{T}^{t} \frac{(\beta - 1)H(t,s)}{2\beta\overline{v}(s)r(s)} \overline{w}^{2}(s)ds$$

$$- \frac{\mu}{2} \int_{T}^{t} \left(\sqrt{\frac{H(t,s)}{\beta\overline{v}(s)r(s)}} \overline{w}(s) + \sqrt{\beta\overline{v}(s)r(s)}h(t,s) \right)^{2} ds.$$
(2.43)

Therefore, for all $t \ge T \ge t_1$, we have

$$\int_{T}^{t} \left[H(t,s)\overline{\psi}(s) - H(t,s) \frac{p^{2}(s)\overline{v}(s)}{2\mu r(s)} - \frac{\mu\beta}{2}\overline{v}(s)r(s)h^{2}(t,s) \right] ds$$

$$\leq H(t,T)\overline{w}(T) - \mu \int_{T}^{t} \frac{(\beta-1)H(t,s)}{2\beta\overline{v}(s)r(s)}\overline{w}^{2}(s)ds$$

$$- \frac{\mu}{2} \int_{T}^{t} \left(\sqrt{\frac{H(t,s)}{\beta\overline{v}(s)r(s)}}\overline{w}(s) + \sqrt{\beta\overline{v}(s)r(s)}h(t,s) \right)^{2} ds.$$
(2.44)

Following the same lines as in the proof of Theorem 2.1, we have

$$\lim_{t \to \infty} \sup \frac{1}{H(t,t_0)} \int_{t_0}^t \left(H(t,s)\overline{\psi}(s) - H(t,s) \frac{p^2(s)\overline{v}(s)}{2\mu r(s)} - \frac{\mu\beta}{2}\overline{v}(s)r(s)h^2(t,s) \right) ds$$

$$\leq \int_{t_0}^{t_1} |\overline{\psi}(s)| ds + |\overline{w}(t_1)| < \infty$$
(2.45)

which contradicts the assumption (2.19).

This completes the proof.

Theorem 2.6. Let (2.19) holds. Assume that there exist functions $f \in C^1([t_0, \infty); R)$ and $\phi \in C([t_0, \infty), R)$ such that, for all $t \ge t_0$, any $T \ge t_0$, and for some $\beta > 1$,

$$\lim_{t \to \infty} \sup \frac{1}{H(t,T)} \int_{T}^{t} \left[H(t,s) \left(\overline{\psi}(s) - \frac{p^{2}(s)\overline{v}(s)}{2\mu r(s)} \right) - \frac{\mu \beta}{2} \overline{v}(s) r(s) h^{2}(t,s) \right] ds \ge \phi(T)$$
 (2.46)

and (2.21) holds, where $\psi(t)$, $\overline{v}(t)$ are defined as in Theorem 2.6 and $\phi_+(t) = \max(\phi(t), 0)$, then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Theorem 2.7. *Let all assumptions of Theorem 2.6 be satisfied except that condition* (2.46) *be replaced by*

$$\lim_{t \to \infty} \inf \frac{1}{H(t,T)} \int_{T}^{t} \left[H(t,s) \left(\overline{\psi}(s) - \frac{p^{2}(s)\overline{v}(s)}{2\mu r(s)} \right) - \frac{\mu \beta}{2} \overline{v}(s) r(s) h^{2}(t,s) \right] ds \ge \phi(T). \tag{2.47}$$

Then every solution u(x,t) of the problem (1.1), (1.5) is oscillatory in G.

Remark 2.8. By introducing the parameter β in Theorem 2.3, we derive new oscillation criteria of the problem (1.1), (1.5) which are simpler than that in Theorem A; furthermore, modifications of the proofs through the refinement of the standard integral averaging method allowed us to shorten significantly the proofs of Theorem 2.3. We can also derive a number of oscillation criteria with the appropriate choice of the function H and ρ , here, we omit the details.

3. Examples

Now, we consider these following examples.

Example 3.1. Consider the partial differential equation

$$\frac{\partial}{\partial t} \left[\frac{1}{t} \frac{\partial}{\partial t} \left(u(x,t) + \frac{1}{2} u(x,t-\pi) \right) \right] + 2\cos t \frac{\partial u(x,t)}{\partial t}$$

$$= \Delta u(x,t) + \frac{1}{t^2} \Delta u \left(x, t - \frac{3}{2} \pi \right)$$

$$- \left(\frac{2}{t^3} + t \cos^2 t + \sin t \right) f(u(x,t)) - \frac{1}{t^2} f_1(u(x,t-\pi)), \quad (x,t) \in (0,\pi) \times (0,\infty),$$
(3.1)

with the boundary condition

$$u_x(0,t) = u_x(\pi,t) = 0, \quad t > 0,$$
 (3.2)

where $f(u) = u^3 + u$, $f_1(u) = ue^u + u$.

Here, N = 1, l = 1, s = 1, m = 1, $\mu = 1$, r(t) = 1/t, $p(t) = 2\cos t$, $q(x,t) = q(t) = (2/t^3 + t\cos^2 t + \sin t)$, $q_1(x,t) = 1/t^2$, $f(u) = f_j(u) = u$, a(t) = 1, $a_1(t) = 1/t^2$, $\rho_1(t) = (3/2)\pi$, $\sigma_1 = \pi$, $\tau_1 = \pi$.

Let

$$y(t) = \frac{1}{t} + t\cos t,\tag{3.3}$$

then

$$v(t) = t^2, \qquad \psi(t) = t^{-1}.$$
 (3.4)

Let n = 3, for any $\beta \ge 1$,

$$\lim_{t \to \infty} \sup \frac{1}{t^2} \int_1^t \left[(t - s)^2 s^{-1} - \beta s^2 \frac{1}{s} \right] ds = \lim_{t \to \infty} \sup \frac{1}{t^2} \int_1^t \left[(t - s)^2 s^{-1} - \beta s \right] ds = \infty.$$
 (3.5)

Therefore, Corollary 2.2 holds, then every solution u(x,t) of the problem (3.1), (3.2) oscillates in $(0,\pi)\times(0,\infty)$.

Example 3.2. Consider the partial differential equation

$$\frac{\partial}{\partial t} \left[\left(1 + \frac{1}{2t^3} \right) (2 + \sin t) \frac{\partial u(x, t)}{\partial t} \right] + \frac{3}{t} \left(1 + \frac{1}{2t^3} \right) (2 + \sin t) \frac{\partial u(x, t)}{\partial t}$$

$$= 3\Delta u(x, t) + (2 - \cos t) \Delta u \left(x, t - \frac{3}{2}\pi \right) - t^{-3} \left(\left(1 - t^3 + 2t^2 - 6t \right) \sin t + 12t \right) f(u(x, t))$$

$$- (2t + \sin t) f_1(u(x, t - \pi)) - 2f_2 \left(u \left(x, t - \frac{\pi}{2} \right) \right), \quad (x, t) \in (0, \pi) \times (0, \infty)$$
(3.6)

with the boundary condition (3.2), where $f(u) = u^5 + u$, $f_1(u) = u \sin^2 u$, $f_2(u) = u^3 \cos^2 u$. Here N = 1, s = 1, m = 2, $\mu = 1$, $r(t) = (1 + 1/2t^3)(2 + \sin t)$, $p(t) = (3/t)(1 + 1/2t^3)(2 + \sin t)$, $q(t) = t^{-3}[(1 - t^3 + 2t^2 - 6t)\sin t + 12t]$, q(t) = 3, $q(t) = 2 - \cos t$, $q(t) = 2 + \sin t$, q(t) =

Let y(t) = 0, then $v(t) = t^3$ and $\psi(t) = v(t)q(t) = (1 - t^3 + 2t^2 - 6t)\sin t + 12t$. Choose $\beta = 2$, n = 3, a straightforward computation yields

$$\limsup_{t \to \infty} \frac{1}{t^2} \int_{T}^{t} \left[(t - s)^2 \left(\left(1 - s^3 + 2s^2 - 6s \right) \sin s + 12s \right) - \left(2s^3 + 1 \right) (2 + \sin s) \right] ds$$

$$= 16 - T^3 \cos T + T^2 (2 \cos T - 6 + 3 \sin T) - 4T \sin T - 3 \cos T = \phi(T).$$
(3.7)

Let $\phi_+(t) = \max(\phi(t), 0)$. It is not difficult to see that

$$\limsup_{t \to \infty} \int_{1}^{t} \frac{\phi_{+}^{2}(s)}{(s^{3} + (1/2))(2 + \sin s)} ds \ge \limsup_{t \to \infty} \int_{1}^{t} \frac{\phi_{+}^{2}(s)}{3(s^{3} + (1/2))} ds = \infty.$$
 (3.8)

By Corollary 2.4, we obtain that every solution of problem (3.6), (3.2) oscillates in $(0, \pi) \times (0, \infty)$.

Note that in this example,

$$\limsup_{t \to \infty} \frac{1}{t^2} \int_1^t 4\left(s^3 + \frac{1}{2} + (2 + \sin s)\right) ds = \infty, \tag{3.9}$$

so the condition (C_2) would not have been satisfied with the same choices of v(t).

Acknowledgments

This research was supported by National Science Foundation of China (11171178), the fund of subject for doctor of ministry of education (20103705110003), and the Natural Science Foundations of Shandong Province of China (ZR2009AM011 and ZR2009AL015).

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