

# ON A PROBLEM OF W. J. FIREY IN CONNECTION WITH THE CHARAC- TERIZATION OF SPHERES

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**Abstract:** Let  $F$  be an arbitrary  $C_\infty$ -smooth closed strictly convex hyper-  
surface in the euclidean space  $\mathbb{R}_d$ . We describe characterizations for  $F$  to be  
a sphere or an ellipsoid in the form that the support function of  $F$  is a pre-  
scribed function of special curvatures of  $F$ . The method of proof consists in  
the application of the theory of evolutions.

## 1. Introduction

In 1967 U. Simon ([9], Satz 6.1.) proved that a  $C_\infty$ -smooth, strictly convex closed hypersurface (ovaloid) in the euclidean space  $\mathbb{R}_d$  is a sphere if its  $k$ -th normalized elementary symmetric function  $H_k$  of the principal curvatures  $\kappa_k$  of  $F$  and its support function  $h > 0$  with respect to the inner point  $o$  of  $F$  are related by

$$(1) \quad H_k = G(h) \quad (1 \leq k \leq d-1)$$

for a  $C_1$ -function  $G$  with

$$(2) \quad \frac{dG}{dh} \leq 0.$$

Especially  $F$  must be a sphere for  $k = 1$  with  $H_1 =: H =$  mean curva-  
ture and  $k = d - 1$  with  $H_{d-1} =: K =$  Gauss Kronecker curvature of

*F.* 1953 K. P. Grotemeyer ([5], 3.) showed that hereby assumption (2) cannot be dropped in general by indication of a rotational ellipsoid in  $\mathbb{R}_3$  with

$$(3) \quad 2H = G(h) := \left(\frac{\alpha}{\beta}\right)^{\frac{1}{4}} h + \beta \left(\frac{\alpha}{\beta}\right)^{\frac{3}{4}} h^3 \quad (\alpha, \beta = \text{const.} > 0)$$

where obviously

$$(4) \quad \frac{dG}{dh} > 0.$$

Now the question arises whether a sphere also may be characterized by  $H_k = G(h)$  ( $k = 1, \dots, d-1$ ) for special functions  $G$  with  $\frac{dG}{dh} > 0$  as

$$(5) \quad H = c^{-2} \cdot h$$

or

$$(6) \quad K = c^{-d} \cdot h$$

( $c = \text{const.} > 0$ ). We want to emphasize that the question of a characterization of a sphere by (6) just occurred in a paper of W. J. Firey (1974 [3]). There Firey investigated the evolution procedure  $\{F_t\}_{t \geq 0}$  of the surface of a worn stone in  $\mathbb{R}_3$ , initially being smooth and strictly convex, which is controlled by the evolution equation

$$(7) \quad \frac{\partial h_t}{\partial t} = -\alpha v_t K_t \quad (\alpha = \text{const.} > 0)$$

( $h_t =$  support function,  $v_t =$  volume and  $K_t =$  Gauss curvature of  $F_t$ ). Under the additional assumption that  $F_0$  (and therefore all the  $F_t$  for  $t \geq 0$ ) are centrally symmetric with respect to the origin  $o$  he was able to prove that the  $F_t$  for  $t \rightarrow \infty$  contract to the "round point"  $o$ . In other words this means that the rescaled surfaces

$$(8) \quad \tilde{F}_t := \left(\frac{v_0}{v_t}\right)^{\frac{1}{3}} \cdot F_t \quad (t \geq 0)$$

with constant volume  $v_0$  converge to a sphere  $F_\infty$  of radius  $\left(\frac{3v_0}{4\pi}\right)^{\frac{1}{3}}$  about  $o$  for  $t \rightarrow \infty$  in the Hausdorff topology.

A basic tool of Firey's was the following

**Lemma 1.** *If  $F$  is a smooth ovaloid in the euclidean space  $\mathbb{R}_d$  ( $d \geq 3$ ), centrally symmetric with respect to the origin  $o$  of  $\mathbb{R}_d$ , whose support function  $h$  and Gauss curvature  $K$  are related by*

$$(9) \quad h = c^d \cdot K \quad (c = \text{const.} > 0)$$

(compare (6)) then  $F$  must be the sphere  $S_c(o)$  about  $o$  with radius  $c > 0$  (see [3], Th. 3 in the case  $d = 3$ ).

**Proof.** If one denotes the unit sphere about  $o$  by  $\Omega$  and its volume resp. surface area element by  $\omega_d$  resp.  $d\omega$  then the volumes  $v$  and  $v^*$  of  $F$  and its polar surface  $F^*$  with respect to  $o$  may be computed by

$$(10) \quad v = \frac{1}{d} \int_{\Omega} h \frac{d\omega}{K} = \frac{1}{d} \int_{\Omega} c^d d\omega = c^d \omega_d$$

(see (9)) and

$$(11) \quad v^* = \frac{1}{d} \int_{\Omega} (r^*)^d d\omega = \frac{1}{d} \int_{\Omega} \frac{d\omega}{h^d}$$

if  $r^*$  is the radius function of  $F^*$  with respect to  $o$ . But trivially

$$(12) \quad r \geq h$$

for the radial function  $r$  of  $F$ , and combining (10), (11) and (12) with the formula

$$(13) \quad \omega_d = \frac{1}{d} \int_{\Omega} \frac{h d\omega}{K r^d} = \frac{c^d}{d} \int_{\Omega} \frac{d\omega}{r^d}$$

(see (9)) for the volume  $\omega_d$  of  $\Omega$  we get

$$(14) \quad v \cdot v^* \geq \omega_d \cdot \frac{c^d}{d} \int_{\Omega} \frac{d\omega}{r^d} = \omega_d \cdot \omega_d.$$

But on the other hand the Blaschke–Santaló inequality yields for the minimal value of  $v \cdot v^*$ , attained at the Santaló point of  $F$  which is the origin because of the central symmetry of  $F$ , the estimate

$$(15) \quad v \cdot v^* \leq \omega_d \cdot \omega_d.$$

Now (14) and (15) imply equality in (12) from which the assertion of Lemma 1 immediately follows.  $\diamond$

In this proof the assumption that  $F$  is centrally symmetric with respect to  $o$  is only used to guarantee that  $o$  is the Santaló point of  $F$ , needed for the validity of (15). Therefore Firey in his paper [3], p.10 settles the conjecture that *his Lemma 1 holds true without this assumption*. This problem has not yet been solved up to now but we are able to prove similar sphere characterizations as well as, by the same method, a local version of Firey's problem expressed in the following two theorems:

## 2. Characterizations of spheres

**Theorem 1.** *If  $F$  is an arbitrary smooth ovaloid in  $\mathbb{R}_d$  with*

$$(16) \quad h = c^2 \cdot H$$

or

$$(17) \quad h = c^2 \cdot K^{\frac{1}{d-1}}$$

( $c = \text{const.} > 0$ ) then  $F$  must be a sphere about  $o$  of radius  $c$ :

$$(18) \quad F = S_c(o).^*$$

**Theorem 2.** *Let  $F$  be an arbitrary smooth ovaloid in  $\mathbb{R}_d$  ( $d \geq 3$ ) the shape of which is sufficiently close to a sphere in the sense of the validity of the inequalities*

$$(19) \quad \kappa_k \geq C(\beta)(\kappa_1 + \dots + \kappa_{d-1}) \quad (k = 1, \dots, d-1)$$

for its principal curvatures where  $C(\beta)$  is a suitable constant depending only on  $\beta > \frac{1}{d-1}$  with

$$(20) \quad C(\beta) < \frac{1}{d-1}$$

and

$$(21) \quad \lim_{\beta \rightarrow +\infty} C(\beta) = \frac{1}{d-1}.$$

If moreover  $F$  fulfils the condition

$$(22) \quad h = c^{(d-1)\beta+1} \cdot K^\beta$$

(generalizing (17)) then  $F$  must be a sphere about  $o$  with radius  $c$ :

$$(23) \quad F = S_c(o).$$

**Remark 1.** *In the special case  $\beta = 1$  Th. 2 provides the fact that the sphere locally is the only solution of Firey's problem in the case of no symmetry assumptions for  $F$ .*

**Proof of Th. 1.** We consider the array

$$(24) \quad F_\tau := \gamma(t) \cdot F$$

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\*Addendum after submission: the first part of the theorem has just been proven more generally for hypersurfaces with nonnegative mean curvature  $H$  by G. Huisken (Asymptotic behaviour for singularities of the mean curvature flow, *J. Diff. Geom.* 31 (1990), 285-299, Th. 4.1).

of hypersurfaces homothetic to  $F$  with respect to the origin  $o$  where  $\gamma(t)$  is a suitable  $C_\infty$ -function of the parameter  $\tau \geq 0$  which fulfils the initial condition

$$(25) \quad \gamma(0) = 1.$$

Then we easily compute

$$(26) \quad h_\tau = \gamma(\tau) \cdot h$$

as well as

$$(27) \quad H_\tau = (\gamma(\tau))^{-1} \cdot H$$

and

$$(28) \quad K_\tau = (\gamma(\tau))^{-(d-1)} \cdot K$$

for the support function  $h_\tau$ , the mean curvature  $H_\tau$  and the Gauss curvature  $K_\tau$  of  $F_\tau$ .

The idea of the proof of Th. 1 is now to choose the factor  $\gamma(\tau)$  in (24) in such a way that  $F_\tau$  is the solution of the well known evolution equation for  $F$ :

$$(29) \quad \frac{\partial h_\tau}{\partial \tau} = -H_\tau$$

or

$$(30) \quad \frac{\partial h_\tau}{\partial \tau} = -(K_\tau)^{\frac{1}{d-1}}.$$

Indeed, using (16) and (27) or (17) and (28) we find

$$(31) \quad \frac{\partial h_\tau}{\partial \tau} = \frac{d\gamma(\tau)}{d\tau} \cdot h = \frac{d\gamma(\tau)}{d\tau} \cdot c^2 H = \frac{d\gamma(\tau)}{d\tau} c^2 \gamma(\tau) \cdot H_\tau$$

or

$$(32) \quad \frac{\partial h_\tau}{\partial \tau} = \frac{d\gamma(\tau)}{d\tau} \cdot h = \frac{d\gamma(\tau)}{d\tau} \cdot c^2 K^{\frac{1}{d-1}} = \frac{d\gamma(\tau)}{d\tau} c^2 \gamma(\tau) \cdot (K_\tau)^{\frac{1}{d-1}}$$

such that we have to solve the ordinary differential equation

$$(33) \quad \frac{d\gamma(\tau)}{d\tau} c^2 \gamma(\tau) = -1$$

under the initial condition (25) in order to obtain (29) or (30). The solution of (33) and (25) is

$$(34) \quad \gamma(\tau) = (1 - 2c^{-2}\tau)^{\frac{1}{2}} \quad (0 \leq \tau < \frac{1}{2}c^2 =: T)$$

and therefore we have to consider the array

$$(35) \quad F_\tau := (1 - 2c^{-2}\tau)^{\frac{1}{2}} \cdot F$$

during the time interval  $[0, T)$  with

$$(36) \quad T := \frac{1}{2}c^2 > 0 .$$

The array (35) represents the solution of the evolution procedures (29) and (30) which contract to the origin  $o$ . As Firey did in his paper [3] it is important to introduce the “normalized procedures”

$$(37) \quad \hat{F}_\tau := \left( \frac{A_0}{A_\tau} \right)^{\frac{1}{d-1}} \cdot F_\tau$$

( $A_\tau$  = total area of  $F_\tau$ ) respectively

$$(38) \quad \tilde{F}_\tau := \left( \frac{v_0}{v_\tau} \right)^{\frac{1}{d}} \cdot F_\tau$$

( $v_\tau$  = total volume of  $F_\tau$ ). Obviously in our case the rescaled hypersurfaces  $\hat{F}_\tau$  resp.  $\tilde{F}_\tau$  coincide with  $F$ :

$$(39) \quad \hat{F}_\tau = \tilde{F}_\tau = F \quad (0 \leq \tau < T)$$

because of (35).

We can now apply a theorem Gage and Hamilton ([4], p. 70) that says that in the case  $d = 2$  the coinciding conditions (29) and (30) imply (in the  $C_\infty$ -topology)

$$(40) \quad \lim_{\tau \rightarrow T} \tilde{F}_\tau = S_c(o)$$

as a circle obeying (16) or (17) so that (18) holds because of (39). Moreover, after another theorem of G. Huisken ([6], p. 238) we have in the case  $d > 2$  because of (29) (again in the  $C_\infty$ -topology)

$$(41) \quad \lim_{\tau \rightarrow T} \hat{F}_\tau = S_c(o)$$

as a sphere obeying (16) which yields (18) in connection with (39). Last not least condition (30) in the case  $d > 2$  provides

$$(42) \quad \lim_{\tau \rightarrow T} \tilde{F}_\tau = S_c(o)$$

(in the  $C_\infty$ -topology) as a sphere obeying (17) after Th. 1.3 of B. Chow in [2] and so again (18) holds true because of (39) as the convergence of the rescaled hypersurfaces  $\tilde{F}_\tau$  is an obvious consequence of (39)<sup>1)</sup>. All these facts complete the proof of Th. 1.  $\diamond$

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<sup>1)</sup>We need the convergence of the  $\tilde{F}_\tau$  because there is a gap in the proof of Chow of this fact.

**Proof of Th. 2.** In the same manner as in the proof of Th. 1 we consider the array (24) for which we conclude, inserting (22) together with (28),

$$(43) \quad \begin{aligned} \frac{\partial h_\tau}{\partial \tau} &= \frac{d\gamma(\tau)}{d\tau} h = \frac{d\gamma(\tau)}{d\tau} c^{(d-1)\beta+1} K^\beta = \\ &= \frac{d\gamma(\tau)}{d\tau} c^{(d-1)\beta+1} (\gamma(\tau))^{(d-1)\beta} (K_\tau)^\beta \end{aligned}$$

( $\beta > \frac{1}{d-1}$ ) instead of (32). For this reason we solve the differential equation

$$(44) \quad \frac{d\gamma(\tau)}{d\tau} c^{(d-1)\beta+1} (\gamma(\tau))^{(d-1)\beta} = -1$$

(compare (33)) together with (25) in order to get the evolution equation

$$(45) \quad \frac{\partial h_\tau}{\partial \tau} = -(K_\tau)^\beta$$

for the hypersurfaces  $F_\tau$ . The solution of (44) has the form

$$(46) \quad \gamma(\tau) = \left(1 - ((d-1)\beta + 1)c^{-((d-1)\beta+1)}\tau\right)^{\frac{1}{(d-1)\beta+1}}$$

( $0 \leq \tau < \frac{c^{(d-1)\beta+1}}{(d-1)\beta+1} =: T$ ). Again we have (39) and a generalization of Th. 1.3 of Chow, namely his Th. 5.1 which is valid because of (45) under the additional assumptions (19), (20) and (21) for  $F$ , now implies

$$(47) \quad \lim_{\tau \rightarrow T} \tilde{F}_\tau = S_c(o)$$

is a sphere obeying (22). Finally the combination of (39) and (47) provides the assertion (23) of Th. 2.  $\diamond$

### 3. Characterizations of ellipsoids

At the end of this paper we shall give a new short proof of a characterization of ellipsoids given first by C. Petty (1985 [8]), Def. 7.3 and Lemma 9.6, although under weaker assumptions for the hypersurface  $F$ . The reason to do so is the fact that our proof works with the same method also used for the proofs of Ths. 1 and 2. This characterization of ellipsoids may be formulated as follows in

**Theorem 3.** *If  $F$  is an arbitrary smooth ovaloid in  $\mathbb{R}_d$  with*

$$(48) \quad h = c^{\frac{2d}{d+1}} \cdot K^{\frac{1}{d+1}} \quad (c = \text{const.} > 0)$$

(compare (22) for  $\beta = \frac{1}{d+1} < \frac{1}{d-1}$  which was excluded in Th. 2) then  $F$  must be an ellipsoid about  $o$  of volume  $c^d \omega_d$ :

$$(49) \quad F = E_c(o).$$

**Remark 2.** Without loss of generality we may refer the support function  $h$  of  $F$  in Th. 3 to the Santalo point  $s$  of  $F$  instead to  $o$  (as Petty did) because (48) and Minkowski's relation

$$(50) \quad \int_{\Omega} n \frac{d\omega(n)}{K(n)} = 0$$

imply

$$(51) \quad \int_{\Omega} n \frac{d\omega(n)}{(h(n))^{d+1}} = 0$$

being characteristic for

$$(52) \quad s = o$$

(compare [8] (3.1)).

**Proof of Th. 3.** As before we see as a consequence of assumption (48) that the array of homothetic hypersurfaces

$$(53) \quad F_{\tau} = \gamma(\tau) \cdot F := \left(1 - \frac{2d}{d+1} c^{-\frac{2d}{d+1}} \tau\right)^{\frac{d+1}{2d}} \cdot F$$

( $0 \leq \tau < \frac{d+1}{2d} c^{\frac{2d}{d+1}} =: T$ ) fulfils the evolution equation

$$(54) \quad \frac{\partial h_{\tau}}{\partial \tau} = -(K_{\tau})^{\frac{1}{d+1}}$$

(compare (46) with  $\beta = \frac{1}{d+1}$ !). But it is well known that the evolution controlled by (54) is equivalent to an affine evolution controlled by

$$(55) \quad \frac{\partial x_{\tau}}{\partial \tau} = y_{\tau} := \frac{1}{d-1} \Delta^a x_{\tau}$$

( $x_{\tau}$  = position vector of a point of  $F_{\tau}$  with the affine normal vector  $y_{\tau}$  and the affine Beltrami operator  $\Delta^a$ ; see [1], (1-1) and (1-2) as well as [7]) because of

$$(56) \quad (K_{\tau})^{\frac{1}{d+1}} = \langle y_{\tau}, n_{\tau} \rangle$$

( $n_{\tau}$  = euclidean inner unit normal vector of  $F_{\tau}$ ). Therefore we have especially for  $\tau = 0$

$$(57) \quad y = \frac{d\gamma(\tau)}{d\tau}(0) \cdot x$$

(see (55) and (53)). But (57) characterizes  $F$  as a smooth strictly convex closed affine hypersphere which must be an ellipsoid about  $o$ , the intersection point of the affine normals of  $F$ . Finally a trivial computation shows that  $F$  must have the volume  $c^d \omega_d$  such that we are sure that (49) holds.  $\diamond$

It is possible that our method of proof also is applicable for other characterization problems where the support function and a special curvature are involved.

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