

**THE SET OF PERIODIC SOLUTIONS OF A NEUTRAL  
DIFFERENTIAL EQUATION WITH CONSTANT DELAY  
AND PIECEWISE CONSTANT ARGUMENT**

GEN-QIANG WANG and SUI SUN CHENG

*Recommended by L. Sanchez*

**Abstract:** The complete set of  $\omega$ -periodic solutions is found for a neutral differential equation with constant delay, piecewise constant argument and  $\omega$ -periodic coefficients.

Considerable attention has been given to delay differential equations with piecewise constant arguments by several authors including Cooke and Wiener [1], Shah and Wiener [2], Aftabizaded et al. [3], and others. This class of differential equations has useful applications in biomedical models of disease that has been developed by Busenerg and Cooke [4] and in stabilization of hybrid control systems with feedback delay, where a hybrid system is one with a continuous plant with a discrete (sampled) controller.

In [3], the set of all periodic solutions of the following linear differential equation with constant coefficient and piecewise constant deviating argument

$$y'(t) + ay(t) + by([t - 1]) = 0, \quad t \geq 0,$$

are found, where  $[\cdot]$  is the greatest integer function. In [4], the set of all periodic solutions of a more general equation

$$(1) \quad y'(t) + a(t)y(t) + b(t)y([t - 1]) = 0, \quad t \geq 0,$$

---

*Received:* March 15, 2004; *Revised:* April 21, 2004.

*AMS Subject Classification:* 34K13.

*Keywords:* neutral differential equation; delay; piecewise constant argument; set of periodic solutions.

is identified with the set of solutions of a linear system. It then follows that necessary and sufficient conditions for the existence of nontrivial periodic solutions of (1) can be found by considering a circulant matrix.

In this note, we proceed one step further and consider the following nonhomogeneous equation

$$(2) \quad \left( y(t) + cy(t-1) \right)' + a(t) \left( y(t) + cy(t-1) \right) + b(t) y([t-1]) = p(t),$$

where  $c$  is a real constant such that  $|c| \neq 1$ , and  $a(t)$ ,  $b(t)$  as well as  $p(t)$  are real, continuous functions with positive integer period  $\omega$ . Note that we allow  $c = 0$ , in which case, (2) reduces to the following nonhomogeneous equation with piecewise constant argument.

$$y'(t) + a(t)y(t) + b(t)y([t-1]) = p(t).$$

We will be interested in finding all the  $\omega$ -periodic solutions of (2). To this end, let  $R$  be the set of real numbers and  $Z$  the set of all integers. By a solution  $y = y(t)$  of (2), we mean a continuous function on  $R$  such that  $(y(t) + cy(t-1))'$  exists at each point  $t \in R$ , with the possible exception of the points  $[t] \in R$  where one-sided derivatives exist, and equation (2) is satisfied on each interval  $[n, n+1) \subset R$  with integral endpoints.

**Lemma 1.** *Let  $c$  be a real constant such that  $|c| \neq 1$ . If  $u(t)$  is a real, continuous and  $\omega$ -periodic function on  $R$ , then there is unique  $\omega$ -periodic continuous function  $x(t)$  which is defined on  $R$  and*

$$(3) \quad u(t) = x(t) + cx(t-1), \quad t \in R.$$

Furthermore, if  $|c| < 1$ , then

$$(4) \quad x(t) = \sum_{i=0}^{\infty} (-1)^i c^i u(t-i), \quad t \in R,$$

while if  $|c| > 1$ , then

$$(5) \quad x(t) = \sum_{i=0}^{\infty} (-1)^i \left( \frac{1}{c} \right)^{i+1} u(t+i+1), \quad t \in R.$$

**Proof:** In case  $|c| < 1$ , it is easy to see that  $\sum_{i=0}^{\infty} (-1)^i c^i u(t-i)$  is uniform convergent on compact intervals of  $R$ . If we define  $x(t)$  by (4), then  $x(t)$  is a

real, continuous and  $\omega$ -periodic function on  $R$ . Furthermore, it is not difficult to check that  $x(t)$  satisfies (3). Similarly, in case  $|c| > 1$ , if we definite  $x(t)$  as (5), then  $x(t)$  is real, continuous and  $\omega$ -periodic, and (3) holds.

Suppose  $y(t)$  is a real, continuous and  $\omega$ -periodic function defined on  $R$  which satisfies

$$(6) \quad u(t) = y(t) + cy(t-1), \quad t \in R.$$

From (3), (6) and the fact that  $x(t)$  and  $y(t)$  are  $\omega$ -periodic, we see that for any  $t \in R$ ,

$$(7) \quad |x(t) - y(t)| = |c| |x(t-1) - y(t-1)|.$$

By (7) and the fact that  $x(t)$  and  $y(t)$  are  $\omega$ -periodic, we have

$$(8) \quad \begin{aligned} \max_{0 \leq t \leq \omega} |x(t) - y(t)| &= \sup_{t \in R} |x(t) - y(t)| \\ &= |c| \sup_{t \in R} |x(t-1) - y(t-1)| \\ &= |c| \max_{0 \leq t \leq \omega} |x(t) - y(t)|. \end{aligned}$$

Since  $|c| \neq 1$ , (8) implies  $x(t) = y(t)$  for  $t \in R$ . The proof is complete. ■

The set of all  $\omega$ -periodic solutions of (2) will be denoted by  $\Omega_\omega$ . Note that when  $p(t) = 0$  for  $t \in R$ ,  $\Omega_\omega$ , endowed with the usual addition and (real) scalar multiplication, is a linear space. In order to determine  $\Omega_\omega$ , we set

$$(9) \quad \alpha_n = \exp\left(-\int_{n-1}^n a(u) du\right) - c, \quad n \in Z,$$

$$(10) \quad \beta_n = c \exp\left(-\int_{n-1}^n a(u) du\right) - \int_{n-1}^n b(s) \exp\left(-\int_s^n a(u) du\right) ds, \quad n \in Z,$$

and

$$(11) \quad \gamma_n = -\int_{n-1}^n p(s) \exp\left(-\int_s^n a(u) du\right) ds, \quad n \in Z.$$

Since  $a(t)$ ,  $b(t)$  and  $c(t)$  are  $\omega$ -periodic, it is easy to see that  $\{\alpha_n\}_{n=-\infty}^\infty$ ,  $\{\beta_n\}_{n=-\infty}^\infty$  and  $\{\gamma_n\}_{n=-\infty}^\infty$  are  $\omega$ -periodic sequences.

Let  $\Psi_\omega$  be the set of all solutions of the following system of  $\omega$  linear equations

$$(12) \quad \begin{cases} \beta_1 z_\omega + \alpha_1 z_1 - z_2 = \gamma_1, \\ \beta_2 z_1 + \alpha_2 z_2 - z_3 = \gamma_2, \\ \dots = \dots \\ \beta_\omega z_{\omega-1} + \alpha_\omega z_\omega - z_1 = \gamma_\omega. \end{cases}$$

In case  $\gamma_1 = \dots = \gamma_\omega = 0$ , note that  $\Psi_\omega$  is a linear subspace of  $R^\omega$ .

Note that the system (12) can be written as

$$A_\omega z = \gamma ,$$

where  $z = (z_1, z_2, \dots, z_\omega)^\dagger$ ,  $\gamma = (\gamma_1, \dots, \gamma_\omega)^\dagger$ , and

$$(13) \quad A_\omega = \begin{pmatrix} \alpha_1 & -1 & 0 & 0 & \dots & \dots & \dots & \beta_1 \\ \beta_2 & \alpha_2 & -1 & 0 & \dots & \dots & \dots & 0 \\ 0 & \beta_3 & \alpha_3 & -1 & \dots & \dots & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & 0 & \dots & \beta_{\omega-1} & \alpha_{\omega-1} & -1 \\ -1 & 0 & 0 & 0 & \dots & 0 & \beta_\omega & \alpha_\omega \end{pmatrix}$$

when  $\omega \geq 3$ ,

$$(14) \quad A_\omega = \begin{pmatrix} \alpha_1 & \beta_1 - 1 \\ \beta_2 - 1 & \alpha_2 \end{pmatrix}$$

when  $\omega = 2$  and

$$(15) \quad A_\omega = (\alpha_1 + \beta_1 - 1)$$

when  $\omega = 1$ .

**Theorem 1.** *There is a one to one and onto mapping from  $\Omega_\omega$  to  $\Psi_\omega$ . Furthermore, if  $p(t) = 0$  for  $t \in R$ , then  $\Omega_\omega$  and  $\Psi_\omega$  are isomorphic.*

**Proof:** Let  $y(t)$  be an  $\omega$ -periodic solution of (2). Then for  $n \in Z$ ,

$$\left( (y(t) + cy(t-1))' \right) + a(t) (y(t) + cy(t-1)) + b(t) y(n-1) = p(t), \quad n \leq t < n+1,$$

so that

$$\begin{aligned} & \frac{d}{dt} \left( (y(t) + cy(t-1)) \exp \left( \int_n^t a(u) du \right) \right) + b(t) \exp \left( \int_n^t a(u) du \right) y(n-1) \\ &= p(t) \exp \left( \int_n^t a(u) du \right), \end{aligned}$$

and

$$\begin{aligned} & (y(t) + cy(t-1)) \exp \left( \int_n^t a(u) du \right) - (y(n) + cy(n-1)) \\ &+ y(n-1) \int_n^t b(s) \exp \left( \int_n^s a(u) du \right) ds \\ &= \int_n^t p(s) \exp \left( \int_n^s a(u) du \right) ds, \end{aligned}$$

for  $n \leq t < n + 1$ . Thus

$$\begin{aligned}
 y(t) + cy(t-1) &= \left( y(n) + cy(n-1) \right) \exp\left(-\int_n^t a(u) du\right) \\
 &\quad - y(n-1) \int_n^t b(s) \exp\left(-\int_s^t a(u) du\right) ds \\
 &\quad + \int_n^t p(s) \exp\left(-\int_s^t a(u) du\right) ds
 \end{aligned}
 \tag{16}$$

for  $n \leq t < n + 1$ . Since  $\lim_{t \rightarrow (n+1)^-} y(t) = y(n+1)$ , we see further that for  $n \in Z$ ,

$$\begin{aligned}
 y(n+1) + cy(n) &= \left( y(n) + cy(n-1) \right) \exp\left(-\int_n^{n+1} a(u) du\right) \\
 &\quad - y(n-1) \int_n^{n+1} b(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds \\
 &\quad + \int_n^{n+1} p(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds ,
 \end{aligned}$$

so that

$$\begin{aligned}
 y(n+1) &= \left( \exp\left(-\int_n^{n+1} a(u) du\right) - c \right) y(n) \\
 &\quad + \left\{ c \exp\left(-\int_n^{n+1} a(u) du\right) - \int_n^{n+1} b(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds \right\} y(n-1) \\
 &\quad + \int_n^{n+1} p(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds .
 \end{aligned}$$

In terms of  $\alpha_n, \beta_n$  and  $\gamma_n$  defined by (9), (10) and (11),

$$y(n+1) = \alpha_{n+1}y(n) + \beta_{n+1}y(n-1) - \gamma_{n+1} .
 \tag{17}$$

If we now let  $z_k = y(k-1)$  for  $k \in Z$ , then  $\{z_k\}_{k=-\infty}^\infty$  is a periodic sequence and from (17) we see that the column vector  $(z_1, z_2, \dots, z_\omega)^\dagger$  is a solution of (12), that is,  $(z_1, z_2, \dots, z_\omega)^\dagger \in \Psi_\omega$ .

Conversely, let  $(z_1, z_2, \dots, z_\omega)^\dagger \in \Psi_\omega$ . Define  $z_0 = z_\omega$  and extend the finite sequence  $\{z_0, z_1, \dots, z_\omega\}$  to the unique  $\omega$ -periodic sequence  $\{z_n\}_{n=-\infty}^\infty$ . Let  $y_n = z_{n+1}$  for  $n \in Z$ , and let the function  $u(t)$  on each interval  $[n, n+1) \subset R$  be defined by

$$\begin{aligned}
 u(t) &= (y_n + cy_{n-1}) \exp\left(-\int_n^t a(u) du\right) - y_{n-1} \int_n^t b(s) \exp\left(-\int_s^t a(u) du\right) ds \\
 &\quad + \int_n^t p(s) \exp\left(-\int_s^t a(u) du\right) ds .
 \end{aligned}
 \tag{18}$$

and

$$\begin{aligned}
 (19) \quad u(n+1) &= (y_n + cy_{n-1}) \exp\left(-\int_n^{n+1} a(u) du\right) \\
 &\quad - y_{n-1} \int_n^{n+1} b(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds \\
 &\quad + \int_n^{n+1} p(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds.
 \end{aligned}$$

Noting that  $\{\alpha_n\}_{n=-\infty}^{\infty}$ ,  $\{\beta_n\}_{n=-\infty}^{\infty}$ ,  $\{\gamma_n\}_{n=-\infty}^{\infty}$  and  $\{z_n\}_{n=-\infty}^{\infty}$  are  $\omega$ -periodic sequences and  $y_n = z_{n+1}$  for  $n \in Z$ , from (12) and (19), we see that

$$\begin{aligned}
 u(n+1) &= \left\{ \exp\left(-\int_n^{n+1} a(u) du\right) - c \right\} y_n \\
 &\quad + \left\{ c \exp\left(-\int_n^{n+1} a(u) du\right) - \int_n^{n+1} b(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds \right\} y_{n-1} \\
 &\quad - \left\{ -\int_{n-1}^{n+1} p(s) \exp\left(-\int_s^{n+1} a(u) du\right) ds \right\} + cy_n \\
 &= \alpha_{n+1} y_n + \beta_{n+1} y_{n-1} - \gamma_{n+1} + cy_n \\
 &= y_{n+1} + cy_n,
 \end{aligned}$$

that is

$$(20) \quad u(n+1) = y_{n+1} + cy_n, \quad n \in Z.$$

By (18), (19) and (20), we see that  $u(t)$  is a real, continuous and  $\omega$ -periodic function on  $R$ , furthermore, from Lemma 1, there is a unique  $\omega$ -periodic continuous function  $y(t)$  which is defined on  $R$  and

$$(21) \quad u(t) = y(t) + cy(t-1), \quad t \in R.$$

By (18) and (21), we have

$$\begin{aligned}
 (22) \quad y(t) + cy(t-1) &= (y_n + cy_{n-1}) \exp\left(-\int_n^t a(u) du\right) \\
 &\quad - y_{n-1} \int_n^t b(s) \exp\left(-\int_s^t a(u) du\right) ds \\
 &\quad + \int_n^t p(s) \exp\left(-\int_s^t a(u) du\right) ds
 \end{aligned}$$

for  $t \in R$ . Now, we prove that  $y(n) = y_n$  for  $n \in Z$ . From (22), we have

$$(23) \quad y(n) + cy(n-1) = y_n + cy_{n-1}, \quad n \in Z.$$

Noting that  $\{y(n)\}_{n=-\infty}^{\infty}$  and  $\{y_n\}_{n=-\infty}^{\infty}$  are  $\omega$ -periodic sequences, (23) implies

$$\begin{aligned}
 \max_{0 \leq n \leq \omega-1} |y(n) - y_n| &= \sup_{n \in Z} |y(n) - y_n| = |c| \sup_{n \in Z} |y(n-1) - y_{n-1}| \\
 (24) \qquad \qquad \qquad &= |c| \sup_{n \in Z} |y(n) - y_n| = |c| \max_{0 \leq n \leq \omega-1} |y(n) - y_n|.
 \end{aligned}$$

Because  $|c| \neq 1$ ,  $y(n) = y_n$  for  $n \in Z$ . Furthermore, in view of (22), we know that  $y(t)$  satisfies (16), it is therefore not difficult to check that the function  $y(t)$  is an  $\omega$ -periodic solution of (2). In other words, we have found a one to one and onto mapping from  $\Omega_\omega$  to  $\Psi_\omega$ .

Note that in case  $p(t) \equiv 0$ , we have  $\gamma_1 = \gamma_2 = \dots = \gamma_\omega = 0$ . Thus the solution sets  $\Omega_\omega$  and  $\Psi_\omega$  are linear spaces. It is easily seen that the mapping found in the proof of Theorem 1 is linear. We may thus conclude that the solution spaces  $\Omega_\omega$  and  $\Psi_\omega$  are isomorphic. The proof is complete. ■

In view of the above identification theorem, we can apply standard results in linear algebra to yield the nature of the solutions of (2).

**Theorem 2.** *Suppose  $p(t) \equiv 0$ . Then the dimension of  $\Omega_\omega$  is  $\omega - \text{Rank}(A_\omega)$ .*

In particular, when  $p(t) \equiv 0$ , (2) has a nontrivial  $\omega$ -periodic solution if, and only if,  $\det A_\omega = 0$ . In case  $\omega = 1$ ,  $\det A_1 = 0$  if, and only if,  $\beta_1 + \alpha_1 = 1$ ; in case  $\omega = 2$ ,  $\det A_2 = 0$  if, and only if,  $(\beta_1 - 1)(\beta_2 - 1) = \alpha_1\alpha_2$ ; and in case  $\omega \geq 3$ ,  $\det A_\omega = 0$  if  $\beta_k + \alpha_k - 1 = 0$  for  $k = 1, \dots, \omega$  (since 0 is an eigenvalue and  $(1, 1, \dots, 1)^\dagger$  is the corresponding eigenvector of  $A_\omega$ ).

**Theorem 3.** *Equation (2) has an  $\omega$ -periodic solution if, and only if,  $\text{Rank}(A_\omega) = \text{Rank}([A_\omega \ \gamma])$ , and has infinitely many  $\omega$ -periodic solutions if, and only if,  $\text{Rank}(A_\omega) = \text{Rank}([A_\omega \ \gamma]) < \omega$ . Here  $\gamma = (\gamma_1, \dots, \gamma_\omega)^\dagger$  is defined by (11) and  $[A_\omega \ \gamma]$  is the augmented matrix formed from  $A_\omega$  and  $\gamma$ .*

In particular, equation (2) has a unique  $\omega$ -periodic solution if, and only if,  $\det A_\omega \neq 0$ . As a consequence, when  $\omega = 1$ , (2) has a unique 1-periodic solution if and only if  $\beta_1 + \alpha_1 \neq 1$ ; when  $\omega = 2$ , (2) has a unique  $\omega$ -periodic solution if, and only if,  $(\beta_1 - 1)(\beta_2 - 1) \neq \alpha_1\alpha_2$ .

**Example 1.** Consider the following equation

$$\begin{aligned}
 (y(t) - 3y(t-1))' + (\sin \pi t)(y(t) - 3y(t-1)) + \left(\exp\left(\pi^{-1} \cos \pi t\right)\right) y([t-1]) \\
 (25) \qquad \qquad \qquad = (\cos \pi t) \exp\left(\sin\left(2^{-1}\pi t\right)\right).
 \end{aligned}$$

Let  $c = -3$ ,  $a(t) = \sin \pi t$ ,  $b(t) = \exp(\pi^{-1} \cos \pi t)$  and  $p(t) = (\cos \pi t) \exp(\sin(2^{-1}\pi t))$ . It is easy to verify that  $a(t)$ ,  $b(t)$  and  $p(t)$  are continuous real functions with period 2. Then  $\alpha_1 = 3 + \exp(-2\pi^{-1})$ ,  $\alpha_2 = 3 + \exp(2\pi^{-1})$ ,  $\beta_1 = -3 \exp(-2\pi^{-1}) - \exp(-\pi^{-1})$  and  $\beta_2 = -3 \exp(2\pi^{-1}) - \exp(\pi^{-1})$ . Thus  $(\beta_1 - 1)(\beta_2 - 1) \neq \alpha_1 \alpha_2$ . By Theorem 3, our equation (25) has a unique 2-periodic solution. Furthermore, since the trivial function is not a solution of our equation, this 2-periodic solution is not trivial.  $\square$

### REFERENCES

- [1] COOKE, K.L. and WIENER, J. – Retarded differential equations with piecewise constant delays, *J. Math. Anal. Appl.*, 99 (1984), 265–297.
- [2] SHAN, S.M. and WIENER, J. – Advanced differential equations with piecewise constant argument deviations, *Intern. J. Math. & Math. Sci.*, 6 (1983), 671–703.
- [3] AFTABIZADEH, A.R.; WIENER, J. and XU, J. – Oscillatory and periodic solutions of delay differential equations with piecewise constant argument, *Proc. Amer. Math. Soc.*, 99 (1987), 673–679.
- [4] BUSENERG, S. and COOKE, K.L. – *Models of vertically transmitted diseases with sequential-continuous dynamics*, in “Nonlinear Phenomena in Mathematical Sciences” (V. Lakshmikantham, Ed.), Academic Press, New York, 1982, pp. 179–187.
- [5] WANG, G.Q. and CHENG, S.S. – Note on the set of periodic solutions of a delay differential equations with piecewise constant argument, *Intern. J. Pure & Appl. Math.*, 9 (2003), 139–143.

Gen-qiang Wang,  
Department of Computer Science, Guangdong Polytechnic Normal University,  
Guangzhou, Guangdong 510665 – P.R. CHINA

and

Sui Sun Cheng,  
Department of Mathematics, Tsing Hua University,  
Hsinchu, Taiwan 30043, – R.O. CHINA