

THE ITERATION PROCESS FOR THE NONLINEAR
TWO-DIMENSIONAL OSCILLATION PROBLEM

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Abstract. The initial boundary value problem for an integro-differential Kirchhoff equation is considered in the case of a square domain. To find an approximate solution, step-by-step discretization is performed with respect to spatial variables and a time argument. The obtained cubic system is solved by the iteration method. The method error is estimated.

Key words: Kirchhoff equation, Jacobi nonlinear iteration process, error estimate

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Let us consider the following initial boundary value problem

$$w_{tt} - \left(\lambda + \frac{4}{\pi^2} \int_{\Omega} (w_x^2 + w_y^2) dx dy \right) (w_{xx} + w_{yy}) = 0, \quad (1)$$

$$(x, y) \in \Omega, \quad 0 < t \leq T,$$

$$w(x, y, 0) = w^{(0)}(x, y), \quad w_t(x, y, 0) = w^{(1)}(x, y), \quad (x, y) \in \Omega, \quad (2)$$

$$w(x, y, t) = 0, \quad (x, y) \in \partial\Omega, \quad 0 \leq t \leq T, \quad (3)$$

where $\Omega = \{(x, y) \mid 0 < x < \pi, 0 < y < \pi\}$, $\partial\Omega$ is the boundary of the domain Ω , $w^{(0)}(x, y)$ and $w^{(1)}(x, y)$ are given functions, $\lambda > 0$ and T are given constants.

Equation (1) is a two-dimensional analogue of the well-known Kirchhoff equation [1]

$$w_{tt} - \left(\lambda + \frac{2}{\pi} \int_0^\pi w_x^2 dx \right) w_{xx} = 0 \quad (4)$$

for string vibration. The studies of many researchers are devoted to Kirchhoff type equations (for the bibliography see, e.g., [2], [3]).

We present here one numerical method of the solution of problem (1)–(3). An approximate solution will be sought for as a finite sum

$$w_n(x, y, t) = \sum_{i,j=1}^n w_{nij}(t) \sin ix \sin jy,$$

where the coefficients $w_{nij}(t)$ are calculated from the Galerkin system

$$w''_{nij}(t) + \left[\lambda + \sum_{p,l=1}^n (p^2 + l^2) w_{npl}^2(t) \right] (i^2 + j^2) w_{nij}(t) = 0, \quad (5)$$

$$\begin{aligned}
 & i, j = 1, 2, \dots, n, \\
 w_{nij}(0) &= \frac{4}{\pi^2} \int_{\Omega} w^{(0)}(x, y) \sin ix \sin jy \, dx \, dy, \\
 & (6) \\
 w'_{nij}(0) &= \frac{4}{\pi^2} \int_{\Omega} w^{(1)}(x, y) \sin ix \sin jy \, dx \, dy, \quad i, j = 1, 2, \dots, n.
 \end{aligned}$$

We replace problem (5), (6) by the problem of finding functions $\underline{w}_{nij}(t)$, $i, j = 1, 2, \dots, n$, where the relation between the functions $w_{nij}(t)$ and $\underline{w}_{nij}(t)$ is such that

$$w_{nij}(t) = \frac{1}{\sqrt{i^2 + j^2}} \underline{w}_{nij}(t). \quad (7)$$

We have to solve the following problem

$$\underline{w}''_{nij}(t) + \left[\lambda + \sum_{p,l=1}^n \underline{w}_{npl}^2(t) \right] (i^2 + j^2) \underline{w}_{nij}(t) = 0, \quad (8)$$

$$\begin{aligned}
 & i, j = 1, 2, \dots, n, \\
 \underline{w}_{nij}(0) &= \frac{4}{\pi^2} \sqrt{i^2 + j^2} \int_{\Omega} w^{(0)}(x, y) \sin ix \sin jy \, dx \, dy, \\
 & (9)
 \end{aligned}$$

$$\begin{aligned}
 \underline{w}'_{nij}(0) &= \frac{4}{\pi^2} \sqrt{i^2 + j^2} \int_{\Omega} w^{(1)}(x, y) \sin ix \sin jy \, dx \, dy, \\
 & i, j = 1, 2, \dots, n.
 \end{aligned}$$

To solve the obtained Cauchy problem (8), (9) we will use a difference scheme of symmetrical type. To this end, we introduce the grid $\{t_m \mid 0 = t_0 < t_1 < \dots < t_M = T\}$ with a constant step $\tau = t_m - t_{m-1}$. The approximate values $\underline{w}_{nij}(t_m)$ are denoted by \underline{w}_{nij}^m , $i, j = 1, 2, \dots, n$, $m = 0, 1, \dots, M$.

The scheme has the form

$$\begin{aligned}
 & \frac{\underline{w}_{nij}^m - 2\underline{w}_{nij}^{m-1} + \underline{w}_{nij}^{m-2}}{\tau^2} + \frac{1}{2} \sum_{u=-1,0}^n \left\{ \left[\lambda + \sum_{p,l=1}^n \frac{(\underline{w}_{npl}^{m+u})^2 + (\underline{w}_{npl}^{m+u-1})^2}{2} \right] \right. \\
 & \quad \left. \times (i^2 + j^2) \frac{\underline{w}_{nij}^{m+u} + \underline{w}_{nij}^{m+u-1}}{2} \right\} = 0, \\
 & \quad i, j = 1, 2, \dots, n, \quad m = 2, 3, \dots, M,
 \end{aligned} \quad (10)$$

$$\begin{aligned}
 \underline{w}_{nij}^0 &= \underline{w}_{nij}(0), \\
 \underline{w}_{nij}^1 &= \underline{w}_{nij}(0) + \frac{\tau}{2} \underline{w}'_{nij}(0) \\
 & + \frac{2\tau^2}{\pi^2} \left[\lambda + \frac{4}{\pi^2} \int_{\Omega} ((w_x^{(0)})^2 + (w_y^{(0)})^2) \, dx \, dy \right] \times \\
 & \times \int_{\Omega} (w_{xx}^{(0)} + w_{yy}^{(0)}) \sin ix \sin jy \, dx \, dy, \quad i, j = 1, 2, \dots, n.
 \end{aligned} \quad (11)$$

From system (10) we write a subsystem for fixed m , $2 \leq m \leq M$. Then

$$\frac{8}{\tau^2(i^2 + j^2)} \underline{w}_{nij}^m + \left\{ 2\lambda + \sum_{p,l=1}^n [(\underline{w}_{npl}^m)^2 + (\underline{w}_{npl}^{m-1})^2] \right\} \times (\underline{w}_{nij}^m + \underline{w}_{nij}^{m-1}) = \frac{8}{\tau^2(i^2 + j^2)} \underline{f}_{nij}^m, \quad i, j = 1, 2, \dots, n, \quad (12)$$

where

$$\underline{f}_{nij}^m = 2\underline{w}_{nij}^{m-1} - \underline{w}_{nij}^{m-2} - \frac{\tau^2}{2} \left(\lambda + \sum_{p,l=1}^n \frac{(\underline{w}_{npl}^{m-1})^2 + (\underline{w}_{npl}^{m-2})^2}{2} \right) \times (i^2 + j^2) \frac{\underline{w}_{nij}^{m-1} + \underline{w}_{nij}^{m-2}}{2}, \quad i, j = 1, 2, \dots, n.$$

Let us agree that for the solution of problem (10), (11) we will use iteration layerwise, more exactly, knowing the approximate solutions $\underline{w}_{nij}^{m-2}$ and $\underline{w}_{nij}^{m-1}$, $i, j = 1, 2, \dots, n$, from (12) we find \underline{w}_{nij}^m , $i, j = 1, 2, \dots, n$, by iteration. As to the initial values, i.e. values at the zero and the first layer, formulas (11) make it possible to find \underline{w}_{nij}^0 and \underline{w}_{nij}^1 , $i, j = 1, 2, \dots, n$.

Denote the k -th iteration approximation \underline{w}_{nij}^m by $\underline{w}_{nij,k}^m$, $i, j = 1, 2, \dots, n$, $k = 0, 1, \dots$

From our algorithmic approach it follows that (12) is a system of equations with respect to \underline{w}_{nij}^m , $i, j = 1, 2, \dots, n$. To solve this system we use the non-linear Jacobi iteration process. From (12) it follows that the process has the form

$$\underline{w}_{nij,k+1}^3 + a_{ij} \underline{w}_{nij,k+1}^2 + b_{ij} \underline{w}_{nij,k+1} + c_{ij} = 0, \quad (13)$$

where

$$\begin{aligned} a_{ij} &= \underline{w}_{nij}^{m-1}, \quad b_{ij} = d_{ij} + (\underline{w}_{nij}^{m-1})^2 + \frac{8}{\tau^2(i^2 + j^2)}, \\ c_{ij} &= (d_{ij} + (\underline{w}_{nij}^{m-1})^2) \underline{w}_{nij}^{m-1} - \frac{8}{\tau^2(i^2 + j^2)} \underline{f}_{nij}^m, \\ d_{ij} &= 2\lambda + \sum_{\substack{p,l=1 \\ p \neq i \\ l \neq j}}^n [(\underline{w}_{npl,k}^m)^2 + (\underline{w}_{npl}^{m-1})^2], \quad k = 0, 1, \dots \end{aligned} \quad (14)$$

In addition to (14) we need the values

$$\begin{aligned} r_{ij} &= d_{ij} + \frac{2}{3} (\underline{w}_{nij}^{m-1})^2 + \frac{8}{\tau^2(i^2 + j^2)}, \\ s_{ij} &= \frac{2}{3} \left(d_{ij} + \frac{10}{9} (\underline{w}_{nij}^{m-1})^2 \right) \underline{w}_{nij}^{m-1} - \frac{8}{\tau^2(i^2 + j^2)} \left(\frac{1}{3} \underline{w}_{nij}^{m-1} + \underline{f}_{nij}^m \right), \quad i, j = 1, 2, \dots, n. \end{aligned} \quad (15)$$

Using the Cardano formulas [4], by (13) we write $\underline{w}_{nij,k+1}^m$ in the explicit form

$$\underline{w}_{nij,k+1}^m = \psi_{ij,k}, \quad (16)$$

where

$$\psi_{ij,k} = -\frac{a_{ij}}{3} + \sigma_{ij,1} - \sigma_{ij,2}, \quad i, j = 1, 2, \dots, n, \quad (17)$$

$$\sigma_{ij,v} = \left[(-1)^v \frac{s_{ij}}{2} + \left(\frac{s_{ij}^2}{4} + \frac{r_{ij}^3}{27} \right)^{\frac{1}{2}} \right]^{\frac{1}{3}}, \quad v = 1, 2, \quad i, j = 1, 2, \dots, n. \quad (18)$$

To establish the convergence conditions for the iteration process used, we consider the Jacobi matrix

$$J = \left\{ \frac{\partial \psi_{ij,k}}{\partial \underline{w}_{ni_1j_1,k}} \right\}, \quad i, j = 1, 2, \dots, n, \quad i_1, j_1 = 1, 2, \dots, n. \quad (19)$$

Here ij and i_1j_1 are respectively the numbers of a row and a column of matrix (19).

By (14)–(18) on the principal diagonal of the matrix J we find zeros. As to nondiagonal elements, $i \neq i_1, j \neq j_1$, we find

$$\begin{aligned} \frac{\partial \psi_{ij,k}}{\partial \underline{w}_{ni_1j_1,k}^m} &= -\frac{1}{9} \underline{w}_{ni_1j_1,k}^m \sum_{v=1}^2 \frac{1}{\sigma_{ij,v}^2} \left[2\underline{w}_{nij}^{m-1} \right. \\ &\quad \left. + (-1)^v \left(\underline{w}_{nij}^{m-1} s_{ij} + \frac{1}{3} r_{ij}^2 \right) \left(\frac{s_{ij}^2}{4} + \frac{r_{ij}^3}{27} \right)^{-\frac{1}{2}} \right]. \end{aligned} \quad (20)$$

From (18) we obtain the formulas

$$\sigma_{ij,1} \sigma_{ij,2} = \frac{r_{ij}}{3}, \quad \sigma_{ij,2}^3 - \sigma_{ij,1}^3 = s_{ij}, \quad \left(\frac{s_{ij}^2}{4} + \frac{r_{ij}^3}{27} \right)^{\frac{1}{2}} = \frac{\sigma_{ij,1}^3 + \sigma_{ij,2}^3}{2},$$

which together with (20) give

$$\frac{\partial \psi_{ij,k}}{\partial \underline{w}_{ni_1j_1,k}^m} = \psi_{ij i_1 j_1}^{(1)} + \psi_{ij i_1 j_1}^{(2)}, \quad (21)$$

where

$$\begin{aligned} \psi_{ij i_1 j_1}^{(v)} &= \left(\frac{2}{3} \underline{w}_{ni_1j_1,k}^m \right) (s_{ij})^{v-1} \left(-\frac{2}{3} \underline{w}_{nij}^{m-1} \right)^{2-v} \\ &\quad \times \left(\sigma_{ij,1}^{2v} + \left(-\frac{r_{ij}}{3} \right)^v + \sigma_{ij,2}^{2v} \right)^{-1}, \quad v = 1, 2. \end{aligned} \quad (22)$$

Let us estimate $|\psi_{ij i_1 j_1}^{(v)}|$, $v = 1, 2$. To this end, we consider the functions

$$\begin{aligned} \psi^{(u)}(z) &= (-r)^u + \sum_{v=1}^2 \left[z + (-1)^v (z^2 + r^3)^{\frac{1}{2}} \right]^{\frac{2u}{3}}, \\ -\infty &< z < \infty, \quad r = \text{const} > 0, \end{aligned}$$

$u = 1, 2$. Each function $\psi^{(u)}(z)$, $u = 1, 2$, is even and increasing for any $z \geq 0$ and therefore

$$\min_{-\infty < z < \infty} |\psi^{(u)}(z)| = \psi^{(u)}(0) = (2u - 1)r^u, \quad u = 1, 2.$$

The latter relation, (14), (15) and (22) imply

$$|\psi_{ij_1 j_1}^{(u)}| \leq \frac{[\tau^2(i^2 + j^2)]^u}{26u - 20} |s_{ij}|^{u-1} |\underline{w}_{nij}^{m-1}|^{2-u} |\underline{w}_{ni_1 j_1, k}^m|, \quad u = 1, 2. \quad (23)$$

From formulas (21), (23) and (14), (15) it follows that

$$\begin{aligned} \left| \frac{\partial \psi_{ij, k}}{\partial \underline{w}_{ni_1 j_1, k}} \right| &\leq \frac{1}{4} \tau^2(i^2 + j^2) |\underline{w}_{ni_1 j_1, k}^m| \left\{ \frac{1}{6} \tau^2(i^2 + j^2) |\underline{w}_{nij}^{m-1}| \left[\lambda + \right. \right. \\ &+ \left. \sum_{p, l=1}^n \left(\frac{1}{2} (\underline{w}_{npl, k}^m)^2 + \frac{5}{9} (\underline{w}_{npl}^{m-1})^2 \right) \right] + |\underline{w}_{nij}^{m-1}| + |\underline{f}_{nij}^m| \left. \right\}. \quad (24) \end{aligned}$$

Let us introduce the vectors

$$\underline{w}_n^{m-1} = (\underline{w}_{nij}^{m-1})_{i, j=1}^n, \quad \underline{w}_{n, k}^m = (\underline{w}_{ni_1 j_1, k}^m)_{i, j=1}^n, \quad \underline{f}_n^m = (\underline{f}_{nij}^m)_{i, j=1}^n.$$

Besides, we also need vector and matrix norms. For the vector $\mu = (\mu_s)_{s=1}^N$ and the matrix $G = (g_{rs})_{r, s=1}^N$ we define $\|\mu\|_1 = \sum_{s=1}^N |\mu_s|$ and $\|G\|_1 = \max_{1 \leq s \leq N} \sum_{r=1}^N |g_{rs}|$.

Let us consider the sums $\sum_{i, j=1}^n (i^2 + j^2)^u$, $u = 1, 2$. Taking into account that [5]

$$\sum_{l=1}^n l^{2u} = \frac{n(n+1)(2n+1)}{6} \left(\frac{3n^2 + 3n - 1}{5} \right)^{u-1}, \quad u = 1, 2,$$

we obtain

$$\sum_{i, j=1}^n (i^2 + j^2)^u \leq \frac{n^2(n+1)(2n+1)}{3} \left(\frac{6n^2 + 6n - 2}{5} \right)^{u-1}, \quad u = 1, 2. \quad (25)$$

Note that for $u = 1$ in (25) we have the equality.

Let us estimate the norm of the Jacobi matrix (19). By virtue of (24) and (25)

$$\begin{aligned} \|J\|_1 &\leq \frac{\tau^2 n^2 (n+1)(2n+1)}{12} \\ &\times \max_{i_1, j_1} |\underline{w}_{ni_1 j_1, k}^m| \left\{ \frac{\tau^2 (3n(n+1) - 1)}{15} \|\underline{w}_n^{m-1}\|_1 \times \right. \\ &\times \left[\lambda + \sum_{i, j=1}^n \left(\frac{1}{2} (\underline{w}_{nij, k}^m)^2 + \frac{5}{9} (\underline{w}_{nij}^{m-1})^2 \right) \right] + \|\underline{w}_n^{m-1}\|_1 + \|\underline{f}_n^m\|_1 \left. \right\}, \quad (26) \\ & \quad i_1, j_1 = 1, 2, \dots, n. \end{aligned}$$

Applying the principle of compressed mapping [6], we assume that

$$\begin{aligned} \|J\|_1 &\leq q, \quad 0 < q < 1, \\ \|\underline{w}_{n,k}^m - \underline{w}_{n,0}^m\|_1 &\leq \frac{1}{1-q} \|\underline{w}_{n,1}^m - \underline{w}_{n,0}^m\|_1, \quad k = 1, 2, \dots, \end{aligned} \quad (27)$$

is fulfilled. From (26) it follows that for (27) to be hold, it is sufficient that the following biquadratic inequality

$$\alpha \tau^4 + \beta \tau^2 - \gamma \leq 0 \quad (28)$$

be fulfilled with respect to the step τ . Here

$$\begin{aligned} \alpha &= \frac{3n(n+1)-1}{15} \|\underline{w}_n^{m-1}\|_1 \left[\lambda + \frac{1}{2} \left(\|\underline{w}_{n,0}^m\|_1 + \right. \right. \\ &\quad \left. \left. + \frac{1}{1-q} \|\underline{w}_{n,1}^m - \underline{w}_{n,0}^m\|_1 \right)^2 + \frac{5}{9} \sum_{i,j=1}^n |\underline{w}_{nij}^{m-1}|^2 \right], \\ \beta &= \|\underline{w}_n^{m-1}\|_1 + \|\underline{f}_n^m\|_1, \\ \gamma &= \frac{12q}{n^2(n+1)(2n+1)} \left(\|\underline{w}_{n,0}^m\|_1 + \frac{1}{1-q} \|\underline{w}_{n,1}^m - \underline{w}_{n,0}^m\|_1 \right)^{-1}. \end{aligned}$$

Thus we come to a conclusion that if the parameter q and the step τ satisfy inequality (28), then system (12) has a unique solution with respect to the unknowns \underline{w}_{nij}^m , $i, j = 1, 2, \dots, n$. The vector $\underline{w}_n^m = (\underline{w}_{nij}^m)_{i,j=1}^n$ consisting of the components of solutions \underline{w}_{nij}^m is the limit of a sequence of vectors $\underline{w}_{n,k}^m$, as $k \rightarrow \infty$. Moreover, the estimate

$$\|\underline{w}_{n,k}^m - \underline{w}_n^m\|_1 \leq \frac{q^k}{1-q} \|\underline{w}_{n,1}^m - \underline{w}_{n,0}^m\|_1, \quad k = 0, 1, \dots,$$

is true.

Formula (7) enables us to construct with the aid of $\underline{w}_{nij,k}^m$ approximate solutions of the function $w_{nij}(t)$ at the grid points.

The problem considered here for the one-dimensional equation (4) is studied in [7].

References

- [1] Kirchhoff, G., Vorlesungen über Mechanik, Teubner, Leipzig, 1883.
- [2] Medeiros, L., Limaco, I., Menezes, S., Vibrations of an elastic string, Mathematical Aspects, I and II., J. Comput. Anal. Appl., 4, 2(2002), 211–263.
- [3] Peradze, J., A numerical algorithm for the nonlinear Kirchhoff string equation, Numer. Math., 102, 2(2005), 311–342.
- [4] Kurosh, A., A course on higher algebra, Nauka, Moscow, 1975 (in Russian).
- [5] Gradstein, I., Rizhik, I., Tables of integrals, sums, series and products, Nauka, Moscow, 1971 (in Russian).
- [6] Trenogin, V. Functional analysis, Nauka, Moscow, 1980 (in Russian).
- [7] Peradze, J., The Jacobi nonlinear iteration method for a discrete Kirchhoff system, AMIM, I. Vekua Appl. Math. Inst. Tbilisi State Univ., 6, 1(2001), 81–89.

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