

## Ordinary differential equations

first order ODEs

$$\dot{y} = f(t, \underline{y})$$

$t = \text{time}$

$$\frac{dy}{dt}$$

continuous

$$f: I \times D \rightarrow \mathbb{R}^d; \quad D \subset \mathbb{R}^d \ni \underline{y}(t)$$

$d > 1$  system of ODEs

$$\underline{y} = \begin{bmatrix} y_1 \\ \vdots \\ y_d \end{bmatrix}$$

Examples

1)

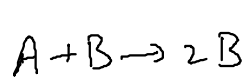
$$\dot{y} = -\lambda y \quad \text{with } \lambda > 0 \text{ constant}$$

$$y(t) = y_0 e^{-\lambda t}, \quad t \in \mathbb{R}$$

2)

$$\dot{y} = (\alpha - \beta y) y$$

logistic differential equation ( $\alpha, \beta \in \mathbb{R}$  constants)



with rate

$$r = k c_A c_B$$

concentrations of A, B

$$\begin{cases} \dot{c}_A = -r \\ \dot{c}_B = r \end{cases}$$

$$c_A + c_B = c_A(0) + c_B(0) = D \text{ constant}$$

$\Rightarrow$  decoupled

$$\begin{cases} \dot{c}_A = -k(D - c_A)c_A \\ \dot{c}_B = k(D - c_B)c_B \end{cases}$$

$$y(0) = y_0 > 0 \Rightarrow y(t) = \frac{\alpha y_0}{\beta y_0 + (\alpha - \beta y_0) e^{-\alpha t}}$$

$$y_0 = 0 \Rightarrow y(t) = 0$$

$$y_0 = 1 \Rightarrow y(t) = 1$$

$(\alpha, \beta = 1)$

$f(y^*) = 0$  (here at  $0, \frac{\alpha}{\beta}$ )  $\rightarrow$  call  $y^* = \text{stationary points}$

if we start at a stationary point

$y(0) = y^*$ , then  $y(t) = y^*$  for all times  $t$

$$[\dot{y}(0) = f(y^*) = 0 = \dot{y} \text{ constant!}]$$

3) Lotka-Volterra ODE

$$\begin{cases} \dot{u} = (\alpha - \beta v) u \\ \dot{v} = (\delta u - \gamma) v \end{cases}$$

$$\underline{y} = \begin{bmatrix} u \\ v \end{bmatrix} \text{ and } f(\underline{y}) = \begin{bmatrix} (\alpha - \beta v) u \\ (\delta u - \gamma) v \end{bmatrix}$$

$u(t) =$  no of prey (rabbits)

$v(t) =$  no of predators (fox)

Initial value problem (IVP): ODE + IV

$$\begin{cases} \dot{y} = f(t, y) \\ y(t_0) = y_0 \end{cases}$$

Remark We may consider only autonomous ODEs, i.e.  $f$  does not depend explicitly on time

$$\dot{y} = f(y)$$

Why? Trick: define a new unknown  $t_1 = t$

unknown  $\underline{v} = \begin{bmatrix} y \\ t_1 \end{bmatrix}$  and write the equation

$$\dot{y} = f(t, y) \text{ as } \underline{\dot{v}} = \begin{bmatrix} \frac{d}{dt} y \\ \frac{d}{dt} t_1 \end{bmatrix} = \begin{bmatrix} f(t, y) \\ 1 \end{bmatrix} = g(\underline{v})$$

$\underline{\dot{v}} = f(\underline{v})$  autonomous!

Hence: every ODE can be re-written as an autonomous ODE

Note Some trick: a higher order ODE (several derivatives of  $y$  are involved)

$$\dot{y}, y'', y^{(3)}, \dots, y^{(k)}$$

can be re-written as a system of ODEs of first order

Note stick mostly on autonomous ODEs of first order

Note If  $f$  is smooth enough, we know that solutions of ODEs

exist (and are unique) for given IV  $y_0$ .

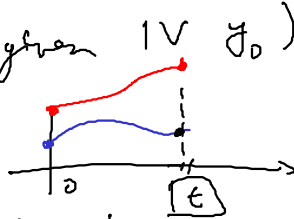
Given

$$\dot{y} = f(y)$$

$$y(0) = y_0$$

$y_0 \mapsto y(t)$  at a later time  $t > 0$

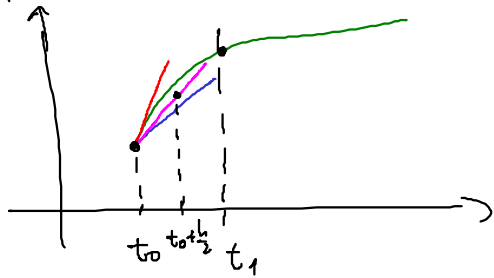
$$\Phi^t y_0 := y(t) \quad \Phi^t : \mathbb{D} \rightarrow \mathbb{D}$$



$\Phi$  = flux associated with the ODE  $\dot{y} = f(y)$   
(continuous flux)

Numerics: construct approximations / discrete versions of  $\Phi$   
discrete flux  $\Phi$

$$\dot{y} = f(y)$$



$h = \text{small}$ ,  $t_1 = t_0 + h$

$$y(t_1) \approx y_1$$

$$y(t) \approx \frac{dy(t)}{dt} = \lim_{h \rightarrow 0} \frac{y(t+h) - y(t)}{h}$$

$$\frac{y_1 - y_0}{h} \approx \begin{cases} f(y_0) \Rightarrow \text{explicit Euler (eE)} \\ f(y(t_0 + \frac{h}{2})) \Rightarrow \text{(iMP)} \\ f(y_1) \Rightarrow \text{implicit Euler (iE)} \end{cases}$$

(eE)  $y_1 = y_0 + h f(y_0)$

(iE)  $y_1 = y_0 + h f(y_1)$  implicit: still have to solve an algebraic eq. for  $y_1$

numerically (later)  
solve

(iMP) implicit midpoint rule

$$y_1 = y_0 + h f(y(t_0 + \frac{h}{2}))$$

$$y_1 = y_0 + h f(\frac{1}{2}(y_0 + y_1))$$

$$f(t, y)$$

↓

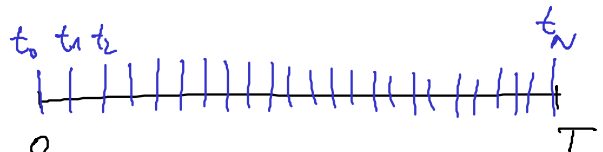
$$f(t_0 + \frac{h}{2}, \frac{1}{2}(y_0 + y_1))$$

another simplification

(iTR)  $y_1 = y_0 + h \frac{1}{2}(f(y_0) + f(y_1))$  implicit trapezoidal rule

Error? Look at error at a given final time T

$$\|y(T) - y_N\|$$



N steps,  $t_0 = 0$

$$h = \frac{T}{N}$$

$$y_0, y_1 \approx y(t_1), y_2 \approx y(t_2), \dots, y_N \approx y(t_N) = y(T)$$

(eE)  $\|y(T) - y_N\| \leq ch^1$

$O(h)$  error

(iE)  $\|y(T) - y_N\| \leq ch^1$

$O(h)$

(iMP)  $\|y(T) - y_N\| \leq ch^2$

$O(h^2)$

(iTR)  $\|y(T) - y_N\| \leq ch^2$

$O(h^2)$

only if solution  $y$  is smooth enough!

Runge-Kutta methods improve it to  $c \cdot h^4, c \cdot h^5, \dots$

ode45 = combination of RK of order  $O(h^4)$  and  $O(h^5)$

Note Convergence order is sometimes not so important...

Ex  $y' = -\lambda y$ ,  $y(t) = y_0 e^{-\lambda t} \xrightarrow{t \rightarrow \infty} 0$

What happens for a large  $T$

(eE):  $y_1 = y_0 + hf(y_0) = y_0 - \lambda h y_0 = \underline{\underline{(1 - \lambda h)y_0}}$  (one step)

$y_2 = y_1 + hf(y_1) = \underline{\underline{(1 - \lambda h)^2 y_0}}$

$y_N = (1 - \lambda h)^N y_0 \xrightarrow{N \rightarrow \infty} \begin{cases} \infty & \text{if } |1 - \lambda h| > 1 \\ 1 & \text{if } |1 - \lambda h| = 1 \\ 0 & \text{if } |1 - \lambda h| < 1 \end{cases}$

Stability condition:  $|1 - \lambda h| < 1 \Leftrightarrow h < \frac{2}{\lambda}$  ( $\lambda > 0$ )

Large  $\lambda \Rightarrow$  need small  $h$  in order to have reasonable results.

(E)  $y_1 = y_0 + hf(y_1) = y_0 - \lambda h y_1 \Rightarrow y_1 + \lambda h y_1 = y_0 \Rightarrow y_1 = \frac{1}{1 + \lambda h} y_0$

$y_2 = \left(\frac{1}{1 + \lambda h}\right)^2 y_0$

...  
 $y_N = \left(\frac{1}{1 + \lambda h}\right)^N y_0 \rightarrow 0$  because  $\lambda, h > 0 \Rightarrow 1 + \lambda h > 1 \Rightarrow 0 < \frac{1}{1 + \lambda h} < 1$

note: no restriction on the time step  $h$

typical for implicit methods (stiff problems!)

but has its price: solving an algebraic equation at each time-step!

note ode45 is explicit!  $\rightarrow$  fast (but sometimes bad!!)  
stiff equation

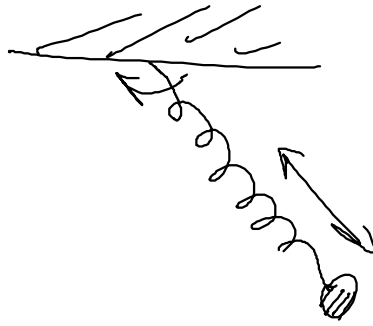
(IMP)  $\rightarrow$  no need of such time step restrictions.

$\hookrightarrow$  "Stability"

note More important than precision might be conservation of important quantities

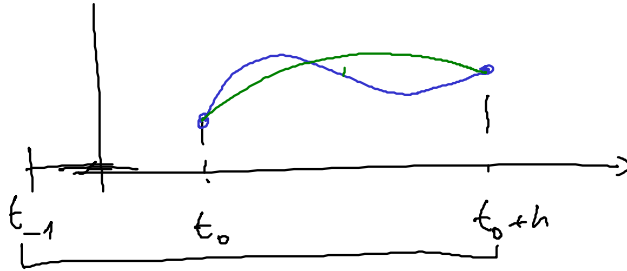
"good" methods

- symplectic integrators
- geometric integrators
- (iMP)



Störmer-Verlet method

$$\begin{cases} \ddot{y} = f(y) \\ y(t_0) = y_0 \\ \dot{y}(t_0) = v_0 \end{cases}$$



given  $y_{k-1} \approx y(t_{k-1})$

$y_k \approx y(t_k)$

approximate  $y(t)$  on  $[t_{k-1}, t_k]$  by a parabola  $p(t)$

going through  $(t_{k-1}, y_{k-1})$  and  $(t_k, y_k)$  such that

$$p''(t_k) = f(y_k)$$

then:  $p(t_{k+1}) \approx y(t_{k+1})$

$t_k - t_{k-1} = t_{k+1} - t_k = h$  (time steps of equal length!)  $\Rightarrow$

$$y_{k+1} = -y_{k-1} + 2y_k + h^2 f(y_k)$$

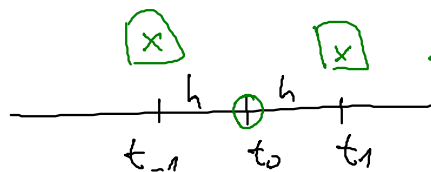
→ explicit method!  $O(h^2)$

→ two-step method, needs a second starting point!

$t_{-1} = t_0 - h, y_1 = -y_{-1} + 2y_0 + h^2 f(y_0)$

↳ where?

$\dot{y}(t_0) = v_0 \Rightarrow$  put  $\frac{y_1 - y_{-1}}{2h} = v_0 \Rightarrow y_{-1} = y_1 - 2hv_0$



$$\Rightarrow y_1 = y_0 + hv_0 + \frac{h^2}{2} f(y_0)$$

Note Another possibility:

$$f(t_k, y(t_k)) = \ddot{y}(t_k) \approx \frac{\cancel{\ddot{y}(t_k)} - \cancel{\ddot{y}(t_{k-1})}}{h} \leftarrow \frac{\ddot{y}(t_{k+1}) - \ddot{y}(t_k)}{h} \quad \frac{\ddot{y}(t_k) - \ddot{y}(t_{k-1})}{h}$$

$\approx \ddot{y}(t_k)$        $\approx \ddot{y}(t_{k-1})$

with  $y_k \approx y(t_k) \Rightarrow f(t_k, y_k) = \frac{y_{k+1} - 2y_k + y_{k-1}}{h^2} \Rightarrow \dots$  St-V

Starting point: Taylor - Approximation of order 2

$$y_1 \approx y(t_0 + h) = y(t_0) + h \dot{y}(t_0) + \frac{h^2}{2} \ddot{y}(t_0) + \dots$$

$\uparrow$   
 error  $O(h^3)$

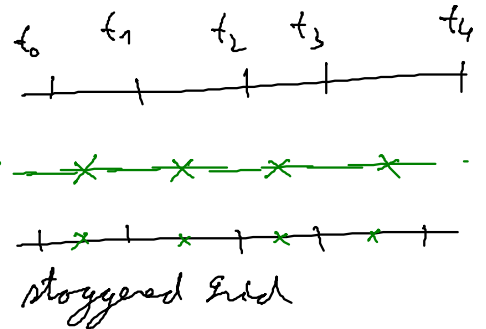
$$\Rightarrow y_1 \approx y_0 + h v_0 + \frac{h^2}{2} f(t_0, y_0)$$

Note The 2-step formulation of StV might be numerically unstable (round-off errors accumulate!)

↓ one-step reformulation

$$\ddot{y} = f(y)$$

$$\leftrightarrow \begin{cases} \dot{y} = v \\ \dot{v} = f(y) \end{cases}$$



$$y_{k+1} - 2y_k + y_{k-1} = h^2 f(y_k)$$

$$\begin{cases} v_{k+\frac{1}{2}} = v_k + \frac{h}{2} f(y_k) \\ y_{k+1} = y_k + h v_{k+\frac{1}{2}} \\ v_{k+1} = v_{k+\frac{1}{2}} + \frac{h}{2} f(y_{k+1}) \end{cases}$$

put together

$$\rightarrow \begin{cases} v_{k+\frac{1}{2}} = v_{k-\frac{1}{2}} + h f(y_k) \\ y_{k+1} = y_k + h v_{k+\frac{1}{2}} \end{cases} \leftarrow \text{one-step formulation of StV}$$

Even better way to implement StV:

velocity-Verlet:

$$\begin{cases} y_{k+1} = y_k + h v_k + \frac{h^2}{2} f(y_k) \\ v_{k+1} = v_k + \frac{h}{2} (f(y_k) + f(y_{k+1})) \end{cases}$$

stable  
 $y_k, v_k$  in  
 one step  
 $O(h^2)$   
 energy cons...

# Splitting Methods

$$\begin{cases} \dot{y} = f(y) = f_a(y) + f_b(y) \\ y(0) = y_0 \end{cases}$$

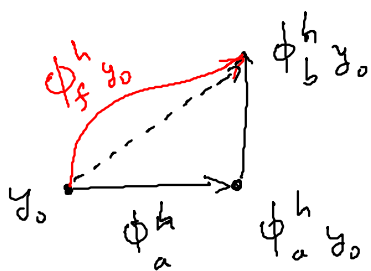
consider 2 ODEs:

$$\dot{u} = f_a(u) \quad \text{and} \quad \dot{v} = f_b(v)$$

$$\text{solution } u(h) = \Phi_a^h u(0)$$

$$v(h) = \Phi_b^h v(0)$$

Idea: recombine the two solutions  $u(h)$  and  $v(h)$  in clever way:

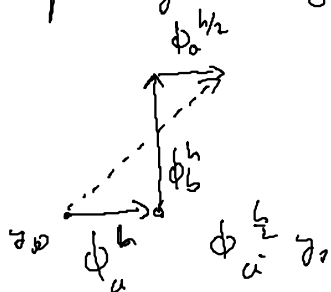


$$\Psi^h y_0 := \Phi_b^h \Phi_a^h y_0 \approx \Phi_f^h y_0$$

Lie-Trotter-splitting  $O(h)$   
not energy conserving

$$\Phi_a^h \Phi_b^h y_0$$

better: put symmetry in method!



$$\Psi^h y_0 := \Phi_a^{h/2} \Phi_b^h \Phi_a^{h/2} y_0$$

Strang splitting  $O(h^2)$   
energy conserving!

Provide we know the exact solutions  $\Phi_a^h u(0)$ ,  $\Phi_b^h v(0)$

then the scheme is simple and explicit!

Newton:  $\dot{r} = a(r) \Leftrightarrow \dot{y} = \begin{bmatrix} \dot{r} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ a(r) \end{bmatrix} = F(y) = \begin{bmatrix} 0 \\ a(r) \end{bmatrix} + \begin{bmatrix} v \\ 0 \end{bmatrix}$

depends only on r!! not on v

$$\begin{cases} \dot{r} = 0 \\ \dot{v} = a(r) \end{cases} \Rightarrow \begin{cases} r(t+h) = r(t) \\ v(t+h) = v(t) + h a(r) \end{cases} \left. \begin{array}{l} r_1 = r_0 \\ v_1 = v_0 + h a(r_0) \end{array} \right\}$$

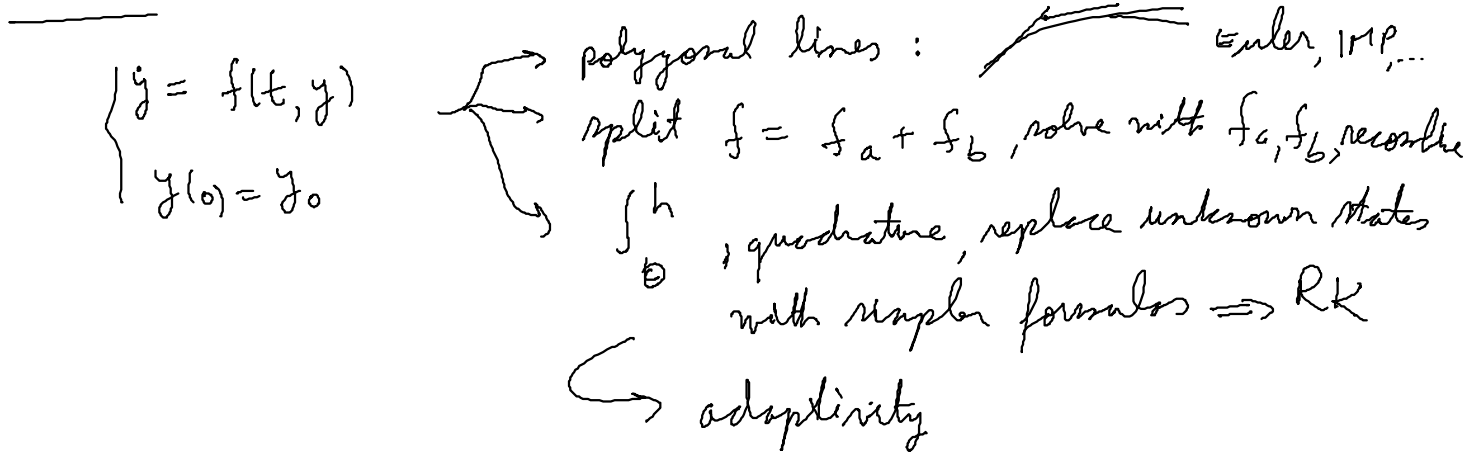
$$\begin{cases} \dot{r} = v \\ \dot{v} = 0 \end{cases} \Rightarrow \begin{cases} r_1 = r_0 + h v_0 \\ v_1 = v_0 \end{cases}$$

recombine: Lie-Trotter  $\Rightarrow \Psi^h \begin{pmatrix} r \\ v \end{pmatrix} = \begin{bmatrix} r+h(v+ha(r)) \\ v+ha(r) \end{bmatrix}$

symplectic error  $O(h)$

Störny-Splitting  $\Rightarrow$  velocity-Verlet!!  $O(h^2)$ , energy conserving, etc.

$\rightarrow$  due to computer problems, read how on black-board.





# Stability and stiff ODEs

logistic  $y' = \lambda y(1-y) \rightarrow ?$

Model problem  $y' = \lambda y$

Usually:  $y^*$  stationary point  $f(y^*) = 0$   
around attractive stat. points: here  $y^* = 1$  attr. stat. point

$\rightarrow$  linearisation around attr. stationary point

$$f(y) = \underbrace{f(y^*)}_0 + (y - y^*) \underbrace{f'(y^*)}_{- \lambda} + O(|y - y^*|^2)$$

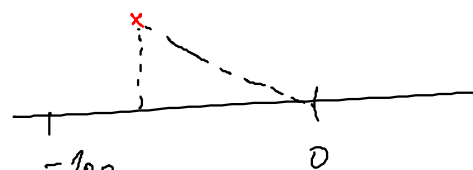
logistic eq.  $f'(y) = \lambda(1-y) + \lambda(-y) = \lambda - 2\lambda y$

at  $y^* = 1 \Rightarrow f'(y^*) = -\lambda$

$\Rightarrow$  linearised eq. is:  $y' = (y-1)(-\lambda)$

Denote  $u = y-1 \Rightarrow \dot{u} = -\lambda u \quad \left\{ \begin{array}{l} \Rightarrow \dot{u} = \omega u \\ \omega = -\lambda \end{array} \right.$

Hence: investigate the stability/stiffness for the model problem  $y' = \lambda y$  in order to conclude of stability/stiffness near attr. stat. points of general non-linear problems.



e. Euler:  $y' = \lambda y \Rightarrow h < \frac{2}{|\lambda|}$

$$y(t) = e^{\lambda t} y(0)$$

$$h < \frac{2}{|-100|}$$

IE: no time-step restriction!

Example explicit trapezoidal rule:

$$\left. \begin{aligned} k_1 &= \lambda y_0 \\ k_2 &= \lambda (y_0 + k_1 h) \end{aligned} \right\} \Rightarrow y_1 = \underbrace{\left(1 + \lambda h + \frac{1}{2}(\lambda h)^2\right)}_{S(\lambda h)} y_0$$

$$y_2 = S(\lambda h) y_1 = S(\lambda h)^2 y_0 \Rightarrow \dots$$

$$y_n = \boxed{S(\lambda h)}^n y_0 \rightarrow 0 \text{ for } n \rightarrow \infty, \text{ when } e^{\lambda t} y_0 \rightarrow 0$$

$$|S(\lambda h)| < 1 \quad \leftarrow \text{must be full filled!}$$

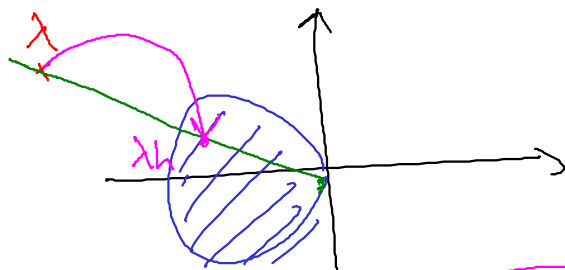
Def  $S(\lambda h) = S(z)$ ,  $z = \lambda h$  is called stability function of the numerical method.

T // stability function of RK:

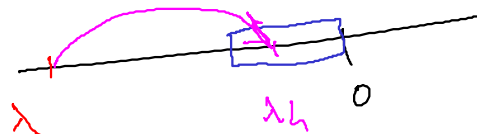
$$\begin{array}{c|c} & A \\ \hline c & b^T \end{array}$$

$$S(z) = \frac{\det(\lambda I - zA + z \mathbf{1} b^T)}{\det(\lambda I - zA)}$$

Ben expl. RK:  $S(z)$  is a polynomial (of degree  $s$ )  
 and  $\{z; |S(z)| < 1\}$  is bounded domain



$$\text{Stability: } |S(\lambda h)| < 1$$



Here: explicit methods always have a bound on  $h$ !!

implicit RK:  $S(z) = \frac{P(z)}{Q(z)}$ , stability domain is unbounded

Typically, implicit methods do not have a bound on  $h$ !

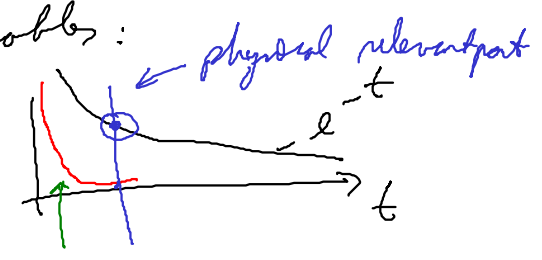
⇒ they allow for large time steps  $h$ !

Def An ODE is called stiff if its stability condition needs a smaller time-step  $h$  than the accuracy of the solution requires.

Ex 
$$\frac{d}{dt} \begin{bmatrix} z_1 \\ z_2 \end{bmatrix} = \begin{bmatrix} -50 & 49 \\ 49 & -50 \end{bmatrix} \begin{bmatrix} z_1 \\ z_2 \end{bmatrix}, z(0) = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

diagonalisation, change of variables:

$$\begin{cases} z_1(t) = e^{-t} + e^{-99t} \\ z_2(t) = e^{-t} - e^{-99t} \end{cases}$$

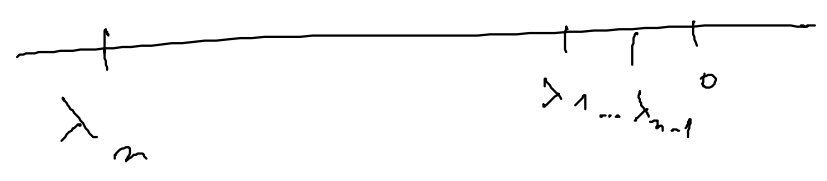


explicit method will be governed by  $e^{-99t}$

e.g. (E):  $h < \frac{2}{99}$

In practical applications it happens

$$\lambda_1, \lambda_2, \dots, \lambda_n$$



odes (explicit methods, time adaptive) will employ very small time steps  $\lambda \Rightarrow \text{small } h$   
 see: time integrator advances very slowly  $\rightarrow$

Then: use implicit methods

ode 23s, ode 15s  
 ↑                      ↑ stiff

price to pay: solving nonlinear equations

→ expensive

→ complicated

→ no guarantees!